

Stars and Nuclei, from Cosmology to Stellar Evolution

Various phases of nucleosynthesis processes
Consequences on nuclei abundances,
on energy production
Nuclear physics characteristics

A. Lefebvre-Schuhl
CSNSM Orsay, France

July 2nd 2009

The Secrets of the Atomic Nucleus, Strasbourg

Periodic table of the elements

Mendeleiev

GROUP 1	IA																18	VIIIA							
1	1	2																	10	2					
	H																	He							
2	Li	Be																	B	C	N	O	F	Ne	
3	Na	Mg																	Al	Si	P	S	Cl	Ar	
4	K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ce	As	Se	Br	Kr							
5	Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe							
6	Cs	Ba																	Hg	Tl	Pb	Bi	Po	At	Rn
7	Fr	Ra																							



Lanthanids

57	58	59	60	61	62	63	64	65	66	67	68	69	70	71
La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu
LANTHANUM	CERURIUM	PRASEODYMUM	NEODYMIUM	PROMETHIUM	SAMARIUM	EUROPIUM	GADOLINIUM	TERBIUM	DYSPROSIUM	HOLEMIUM	ERBIUM	THULIUM	Ytterbium	LUTETIUM

Actinides

89	90	91	92	93	94	95	96	97	98	99	100	101	102	103
Ac	Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr
ACTINIUM	THORIUM	Protactinium	URANIUM	NEPTUNIUM	PLUTONIUM	AMERICIUM	CURVIUM	BERKELIUM	CALIFORNIUM	EINSTEINIUM	FERMURIUM	Mendelevium	Nobelium	LAWRENCIUM

Chemistry → A few number of different elements

First spectroscopic observation of stars → The same elements

How to understand origin and abundance of elements in Universe ?

Where does the star energy production come from?

Sky fascination

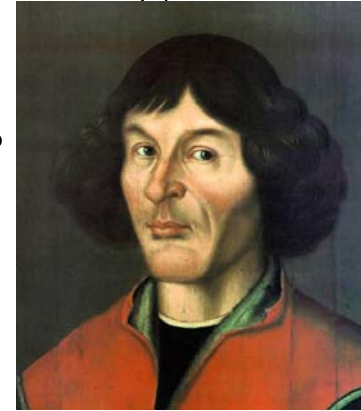
Sun during the day; Stars and other celestial objects during the night

Observation ; Use

Technical advances → always more advanced observations

→ advances in physics knowledge

(Kepler laws ; velocity limitation...)



Now knowledge still advances from Sun and star study

Solar neutrino puzzle → Neutrino mass

Lens, mirror surface advances

New detectors

Numerous star observations

→ Researches in astrophysics ;

→ In nuclear physics ;

→ In particle physics ;

From Galilei's lens,

*To the Hubble Space
Telescope
and ...*



*→ Some stars have similar
characteristics*

→ Star classification

→ Expansion of the Universe



Newton's telescope

But, a remaining haunting question !

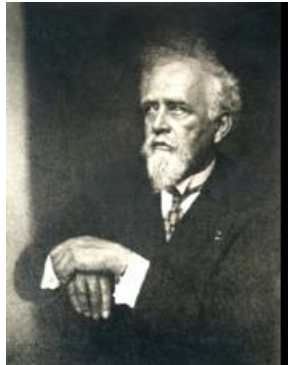
Will the Sun die ?

19th century : Thermodynamics laws

→ The Sun is shining! It's impossible

H. N. Russell in 1919

Too much energy during such a long time !



J. Perrin in 1920

Only in the 1920s : ~ 80 years later

Solar energy is produced inside the Sun
by nuclear reactions

Consequences : → Solar energy may be understood

→ A lot of various results may be connected

→ A new coherent theory appears

The Standard model \equiv Standard big-bang

SBB theory

One basic assumption : Laws of physics as presently known were valid at each time of cosmological evolution

In this framework:

Two fundamental SBB underlying assumptions :

- 1) General Relativity offers a valid description of gravity,
- 2) The Universe got hotter than $\sim 10^{11}$ K, so that statistical equilibrium was established between all of its components
(i. e. Weinberg 1972, Kolb and Turner 1990)

Going back in the time \rightarrow *origin of the Universe : 10^{10} years earlier*

within the SBB model \rightarrow *back to $\sim 10^{-12}$ s after the big-bang*

Then, « Beyond the standard model » until the Planck time :
 $\sim 10^{-43}$ s after the big-bang

Supersymmetry ; String theories ;

Variation of the fundamental physical constants ; ...

Totally outside this lecture

Preliminary remarks

Need to compare theory and observational data

Mainly electromagnetic radiation detections

in any wave-length range

from radiative sites as many and varied as possible

*i.e. galaxy (ours and others) ; various kinds of stars ;
interstellar/intergalactic medium ; ...*

Nuclear reaction importance

They change nuclear species in other nuclear species

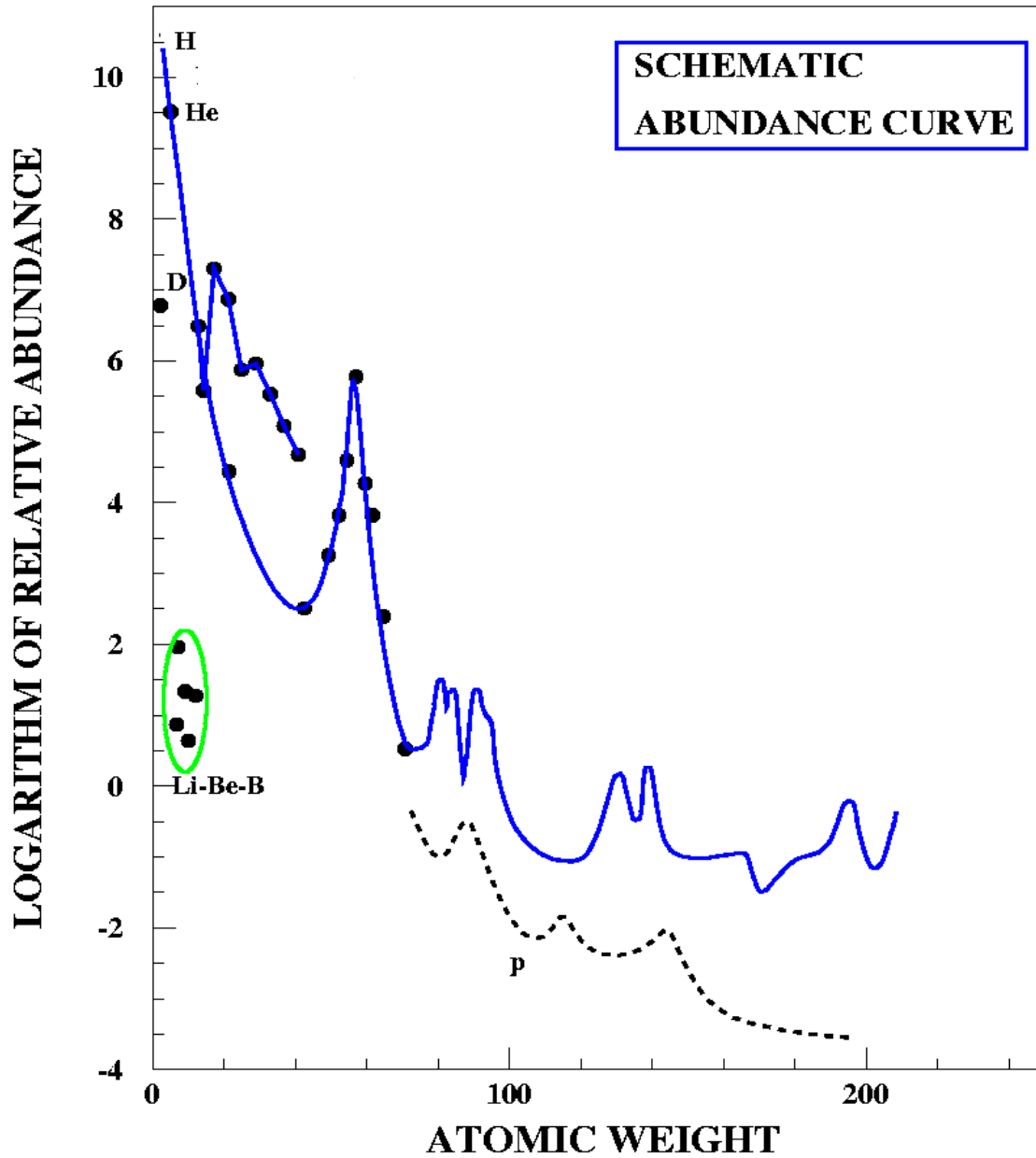
Source of the production of the various chemical elements

They produce or absorb energy

Source of star energy production

...

Some characteristics of the abundances



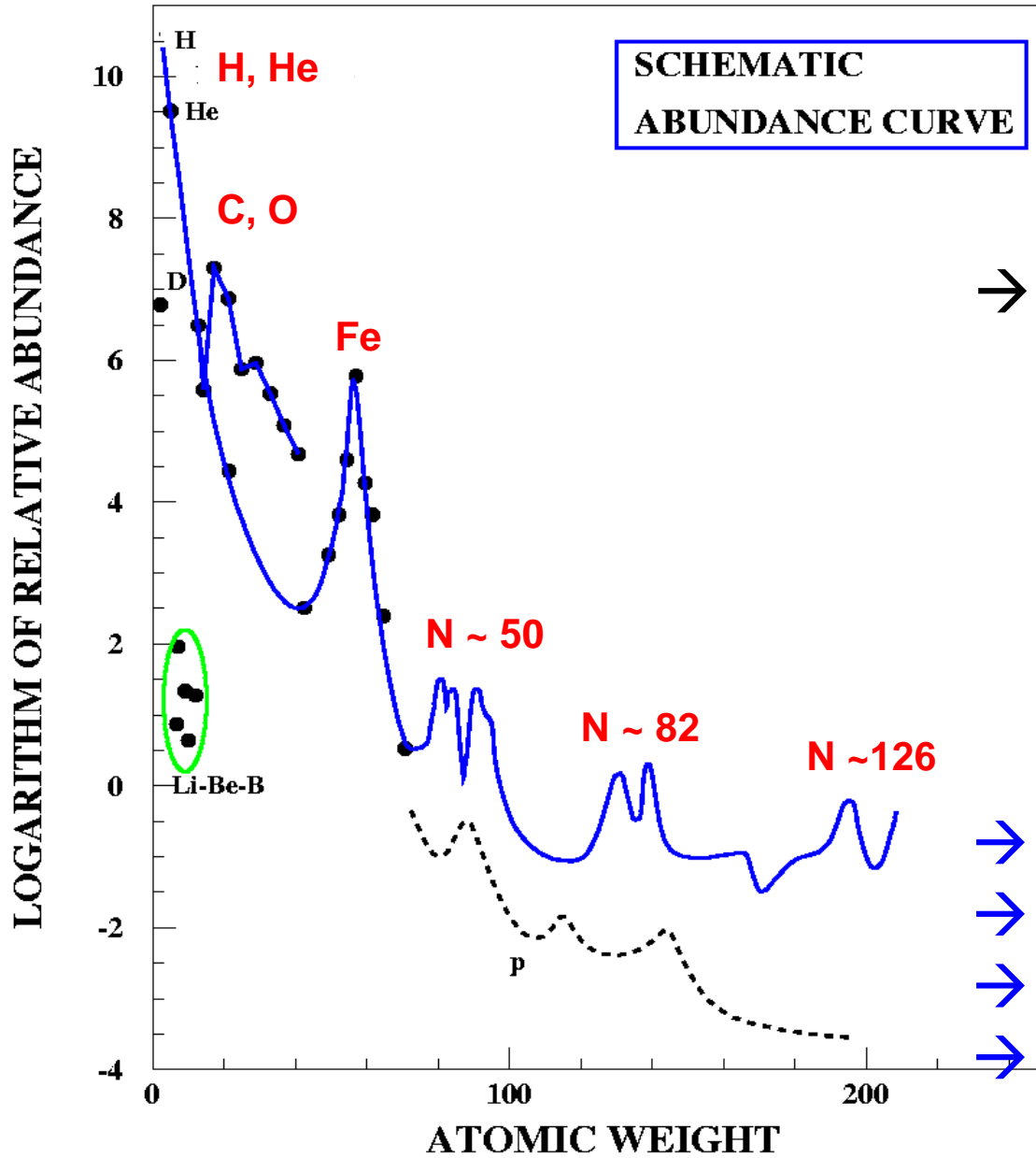
[H] and [He] far above the others

Li, Be and B extremely underabundant

Some peaks are superimposed to a curve decreasing with A

Double peaks appear around $A = 80, 126, 200$

Some characteristics of the abundances (II)



Abundant elements for element species with $N, Z \sim 2, 8, 28, 50, 82, 126$

→ magic numbers correlation

→ Nucleosynthesis is strongly correlated to the nuclear stability, nuclear properties

→ primordial nucleosynthesis

→ stellar nucleosynthesis

→ explosive nucleosynthesis

→ spallative nucleosynthesis

Thermonuclear reaction rate (I)

Reaction : $1 + 2 \rightarrow 3 + 4 + \dots$

$$Q = (M_3 + M_4 + \dots) - (M_1 + M_2)$$

Number of reactions per time unit \leftrightarrow cross section

$$\frac{dN_{12}}{dt} = IxN_2\sigma$$

I = beam intensity

x = target thickness

N₂ = volumic density of the particle in the target

σ = cross section

Stars \neq targets :

Number of interactions per time and volume unit

$$\frac{d\mathcal{N}_{12}}{dt} = N_1N_2\sigma v$$

v = relative velocity of type 1 and 2 particles

N_i = volumic density of type i particles

ℳ₁₂ = number of interactions per volume unit

Thermonuclear reaction rate (II)

Reaction : $1 + 2 \rightarrow 3 + 4 + \dots$

Star : gas mixing with \neq velocities

\rightarrow reaction rate per volume unit

$$r_{12} = \frac{d\mathcal{N}_{12}}{dt} = N_1 N_2 \int_0^\infty \sigma(v) \varphi(v) v dv$$

$\varphi(v)$ = velocity distribution

N_i = volumic densities of type i particles

Thermonuclear reaction rate (per mole of interacting particles) :

$$\mathcal{N}_A \langle \sigma v \rangle = \mathcal{N}_A \int_0^\infty \sigma(v) v \varphi(v) dv$$

where :

\mathcal{N}_A = Avogadro constant

Should be multiplied by $(1 + \delta_{12})^{-1}$: true for \neq or \equiv particles

Reaction rate unit : $\text{cm}^3 \text{s}^{-1} \text{mole}^{-1}$

Thermonuclear reaction rate (II)

In a non degenerate gas mixing composed of non relativistic particles:

$\varphi(v)$ = velocity distribution = Maxwell-Boltzmann distribution

$$\begin{aligned} r_{12} &= \mathcal{N}_A \langle \sigma v \rangle = \mathcal{N}_A \int_0^\infty \sigma(v) v \varphi(v) dv \\ &= \sqrt{8/\pi\mu} \left(\mathcal{N}_A / (kT)^{(3/2)} \right) \int_0^\infty E \sigma(E) e^{(-E/kT)} dE \\ &\quad \text{with } k_B = 8.62 \cdot 10^{-8} \text{ keV/K} \end{aligned}$$

To integrate:

Need of $\sigma(E)$ value for any energy E value

→ measurements, approximations, models ...

Depends on the interacting particles (charged, neutrals, photons, ...)

Astrophysical S factor: $S(E)$

To remove clearly, in $\sigma(E)$:

- *the barrier penetrability*
- *the de Broglie wave length*

$$S(E) = \sigma(E) \cdot E \cdot \exp(2\pi\eta) = \sigma(E) \cdot E \cdot \exp\left(\sqrt{\frac{E_G}{E}}\right)$$

Contributes as : $\sigma(E) \sim \hat{\lambda}^2$
where $\hat{\lambda} \propto E^{-1/2}$ is the de
Broglie wave length

Penetrability
correction ($L=0$)

Gamow energy: $E_G = \frac{1}{2}(2\pi\alpha Z_1 Z_2)^2 \mu c^2 = (0.989 \cdot Z_1 Z_2 A^{1/2})^2$ [MeV]

$\Rightarrow S(E)$ varies more slowly with E compared to $\sigma(E)$

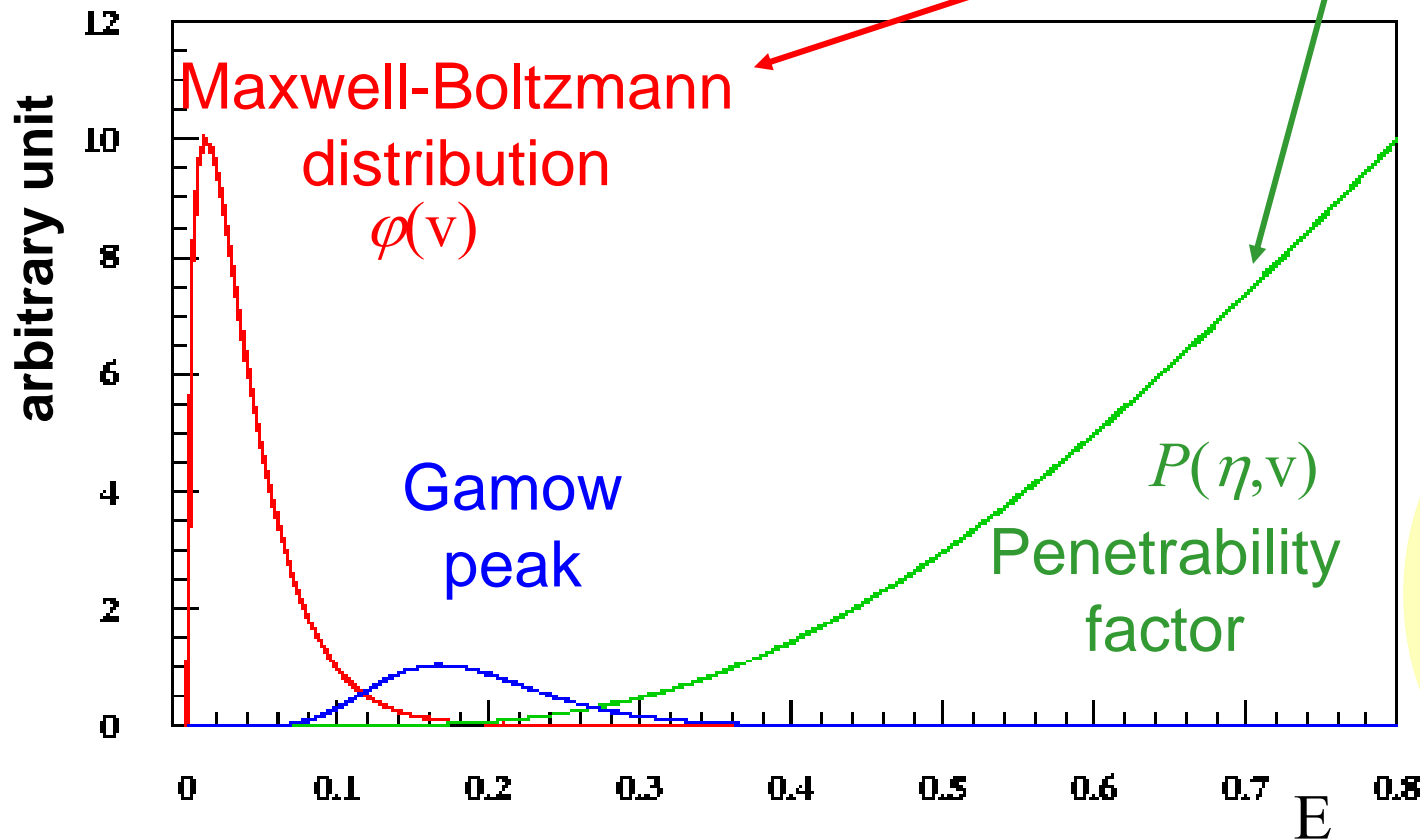
Easier for use when extrapolations are needed

Charged particle reactions: Gamow peak

$$r_{12} = \mathcal{N}_A \langle \sigma v \rangle = \sqrt{8/\pi\mu} \left(\mathcal{N}_A / (kT)^{(3/2)} \right) \int_0^\infty E \sigma(E) e^{(-E/kT)} dE$$

Astrophysical S factor: $S(E) = E \sigma(E) e^{2\pi\eta} = E \sigma(E) e^{(E_G/E)^{1/2}}$

$$r_{12} = \mathcal{N}_A \langle \sigma v \rangle = \sqrt{8/\pi\mu} \left(\mathcal{N}_A / (kT)^{(3/2)} \right) \int_0^\infty S(E) e^{(-E/kT - \sqrt{E_G/E})} dE$$



If $S(E) \approx \text{cst}$:
 Mainly the
 Gamow peak
 energy range
 will contribute

p + p
 1st stellar reaction

$T \approx 15 \cdot 10^6 \text{K}$

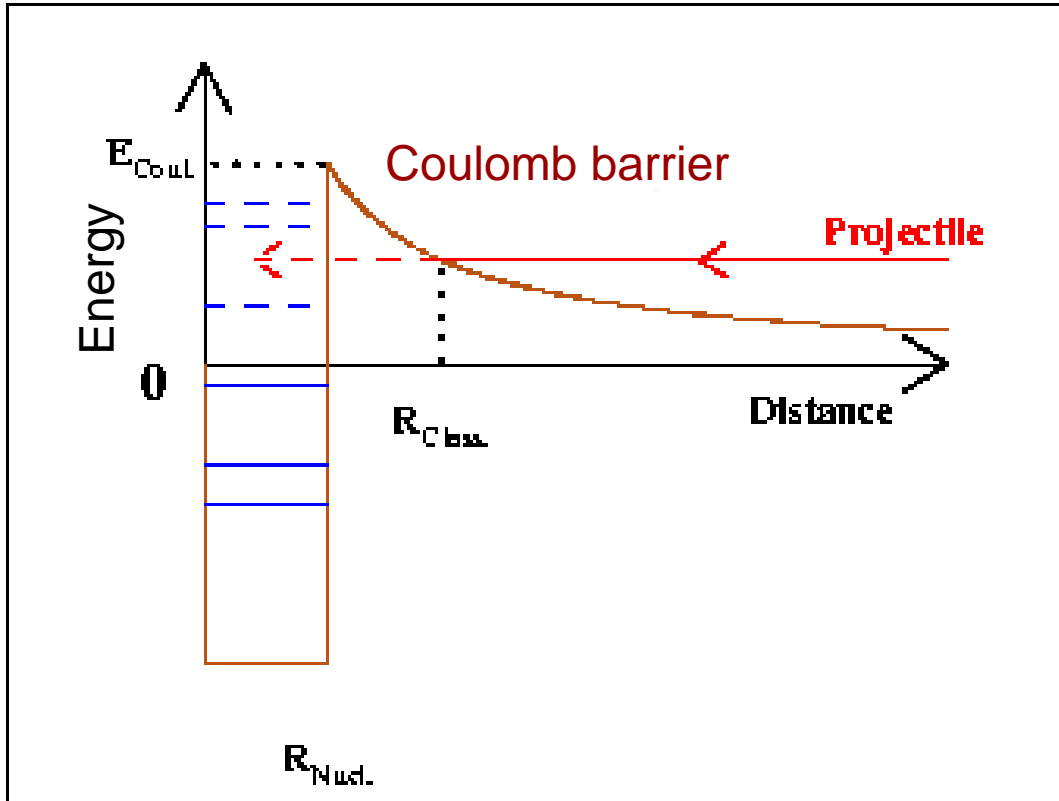
$E_0 \approx 6 \text{ keV}$

$\Delta \approx 6.4 \text{ keV}$

Charged particle reactions

In stars: energy is always smaller than the Coulomb barrier !

Very small cross sections



$$V_c = \frac{Z_1 Z_2 e^2}{4\pi\epsilon_0 R}$$

p+p example :

$$V_c \approx 550 \text{ keV}$$

(for $R = 1.3(A_1^{1/3} + A_2^{1/3})$ in fm)

should be compared to E_{int}

In the core of the Sun : $T \sim 15 \cdot 10^6 \text{ K}$

$$kT \sim 1.3 \text{ keV} \ll V_c \Leftrightarrow d_{\text{min,Coul}}(1.3 \text{ keV}) = 1.1 \cdot 10^3 \text{ fm} \gg R \approx 2.6 \text{ fm}$$

Very small interaction probability for p + p

Astrophysical S factor : $S(E)$

Cross section variation

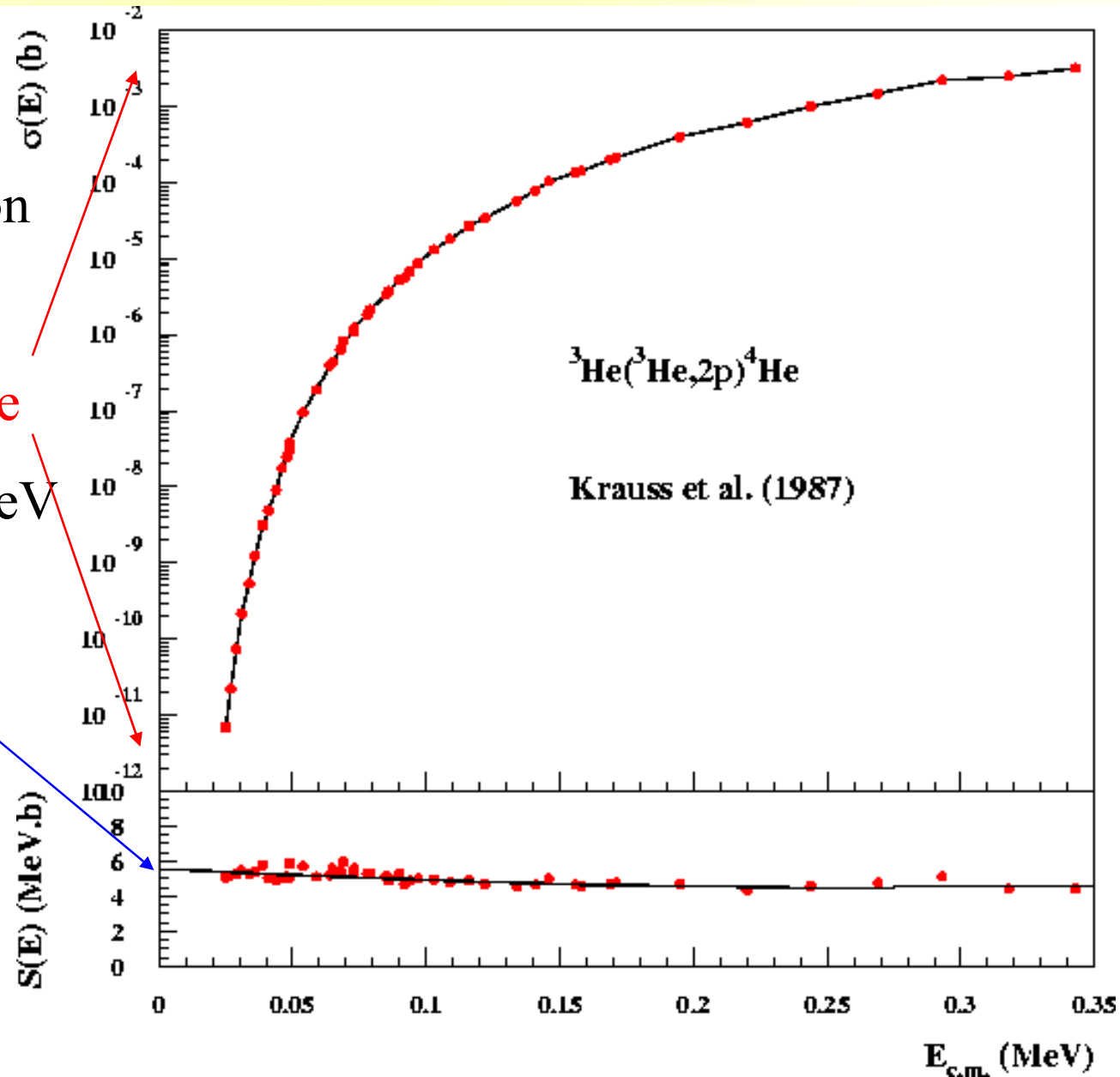
$$\sigma(E) :$$

9 orders of magnitude

between 25 and 340 keV

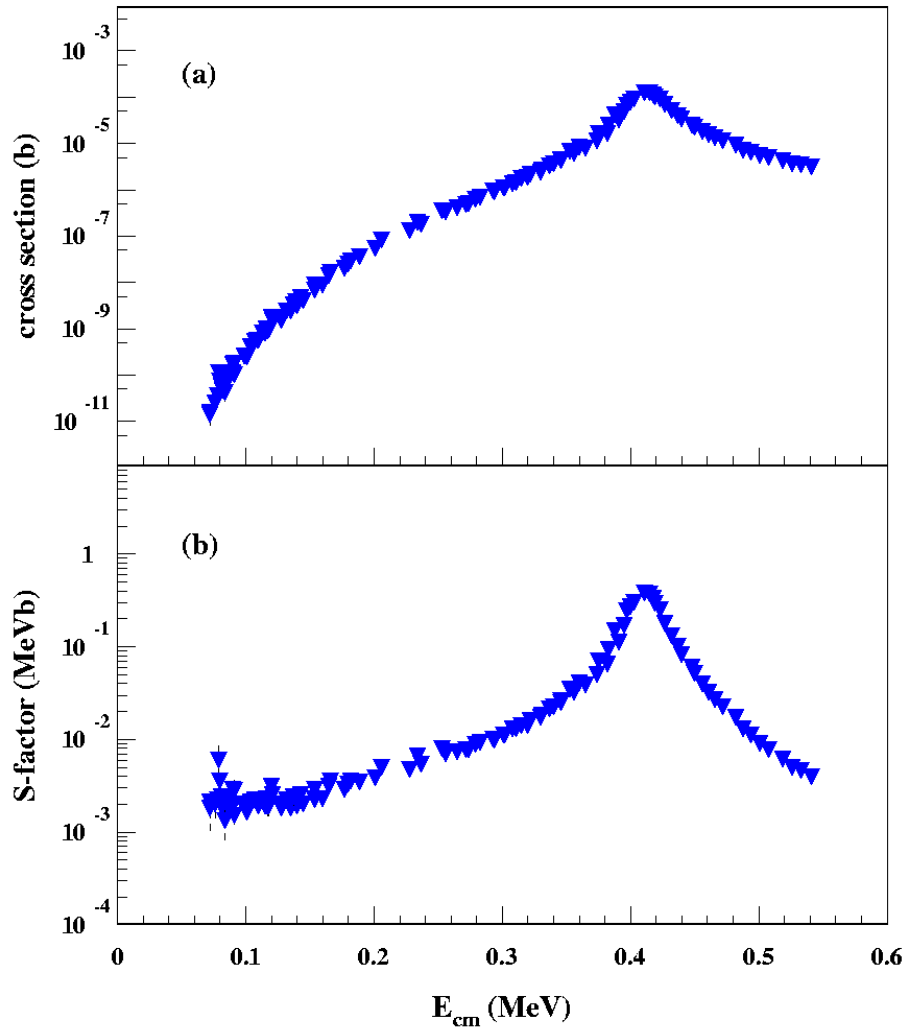
$S(E)$ variation : 50%

⇒ Extrapolation
at low energy

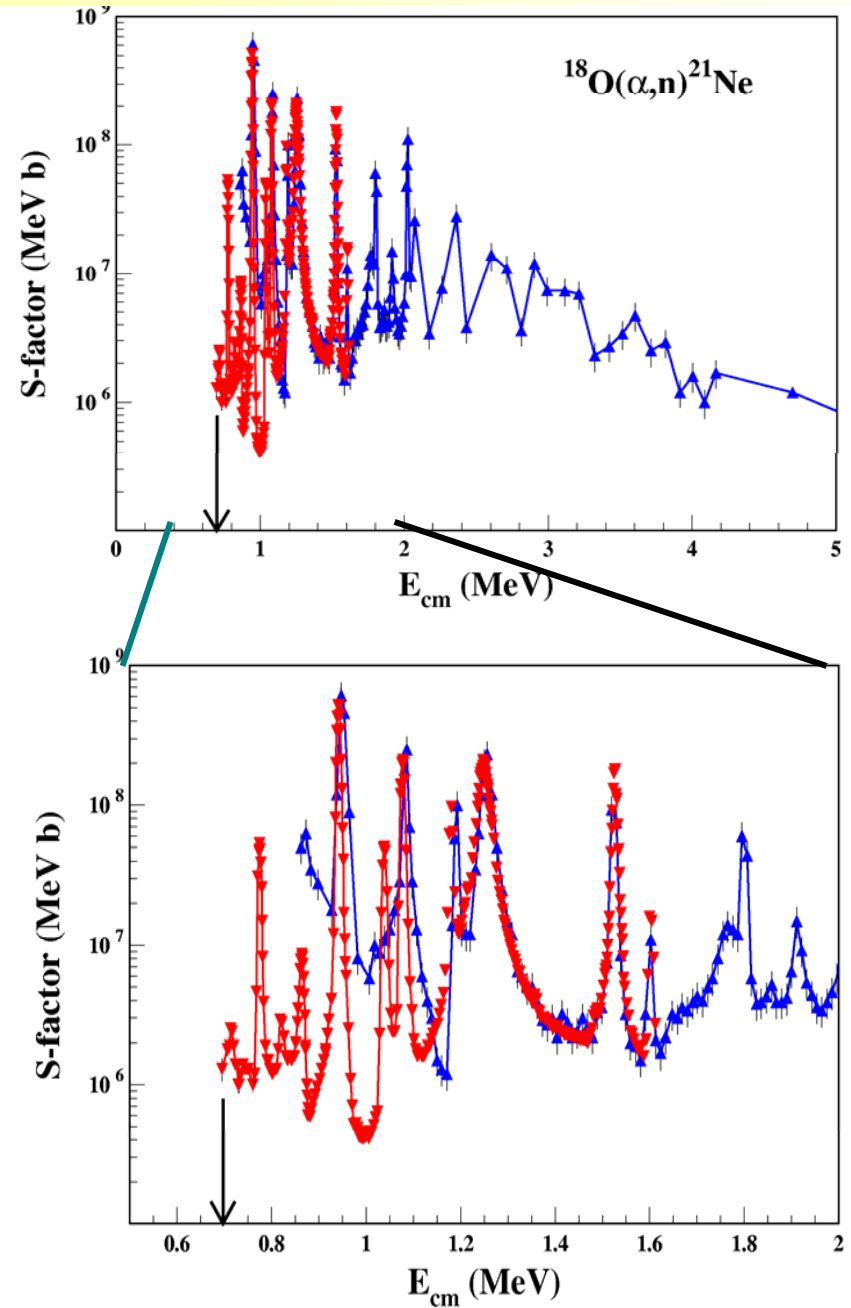


Simple or multiple resonances

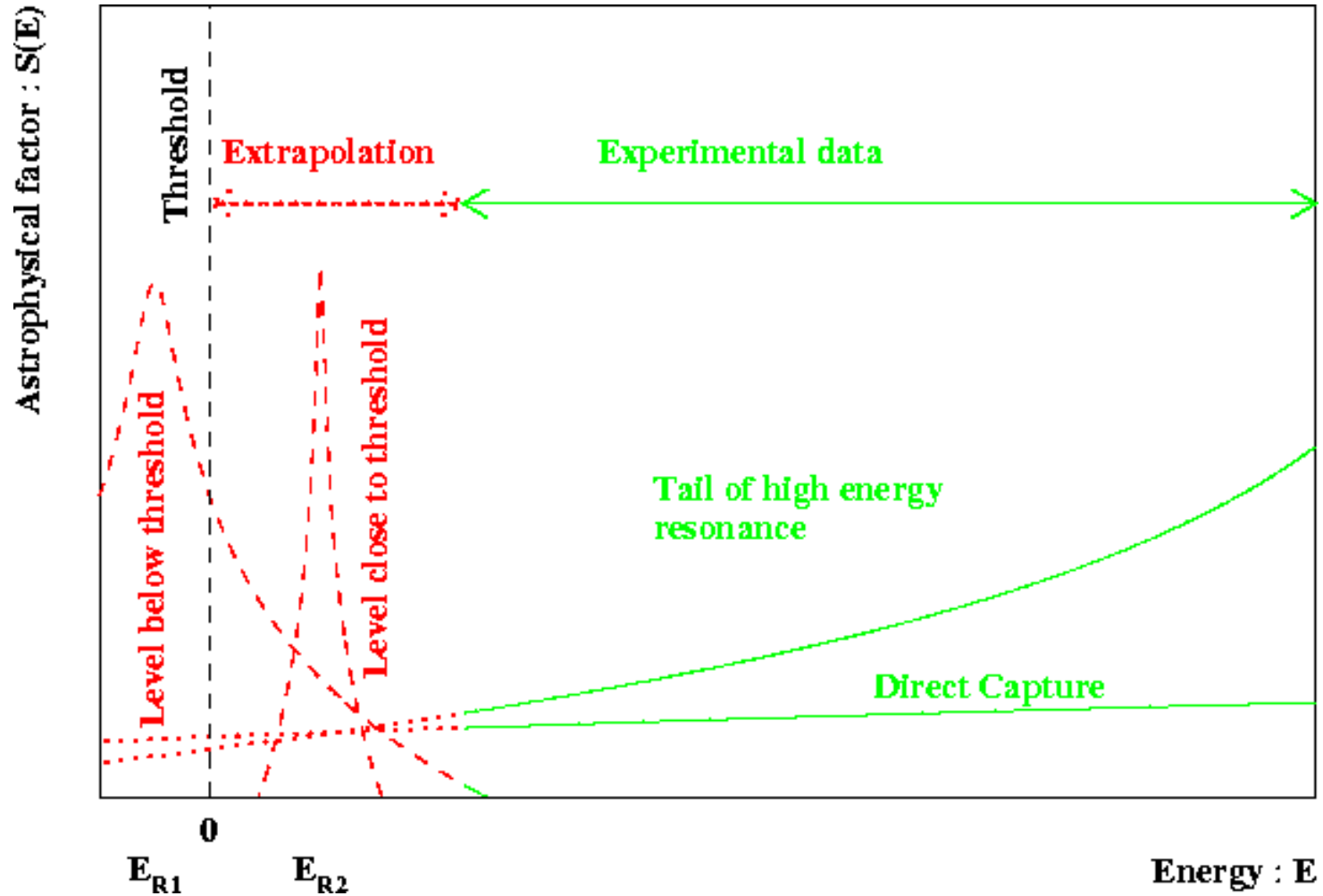
$^{12}\text{C}(p,\gamma)^{13}\text{N}$



$^{18}\text{O}(\alpha,n)^{21}\text{Ne}$



Low energy extrapolation dangers



Nuclear models and experimental data are needed

SBBN framework

Quite well known : Universe is a hot and dense gas

Thermodynamics → Statistical equilibrium between all the particles

Expansion → quick temperature decrease

→ particle evolution

Age of the Universe in the range 1 s to 3 min ;

Temperature cooling down from 10^{10} to 10^9 K

Energy decrease from 1 to 0.1 MeV

Thermodynamics →

Most probable state = state of minimum energy

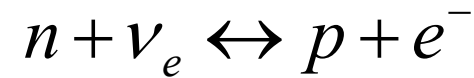
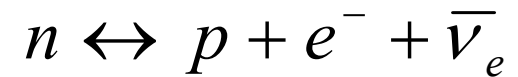
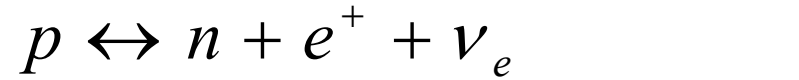
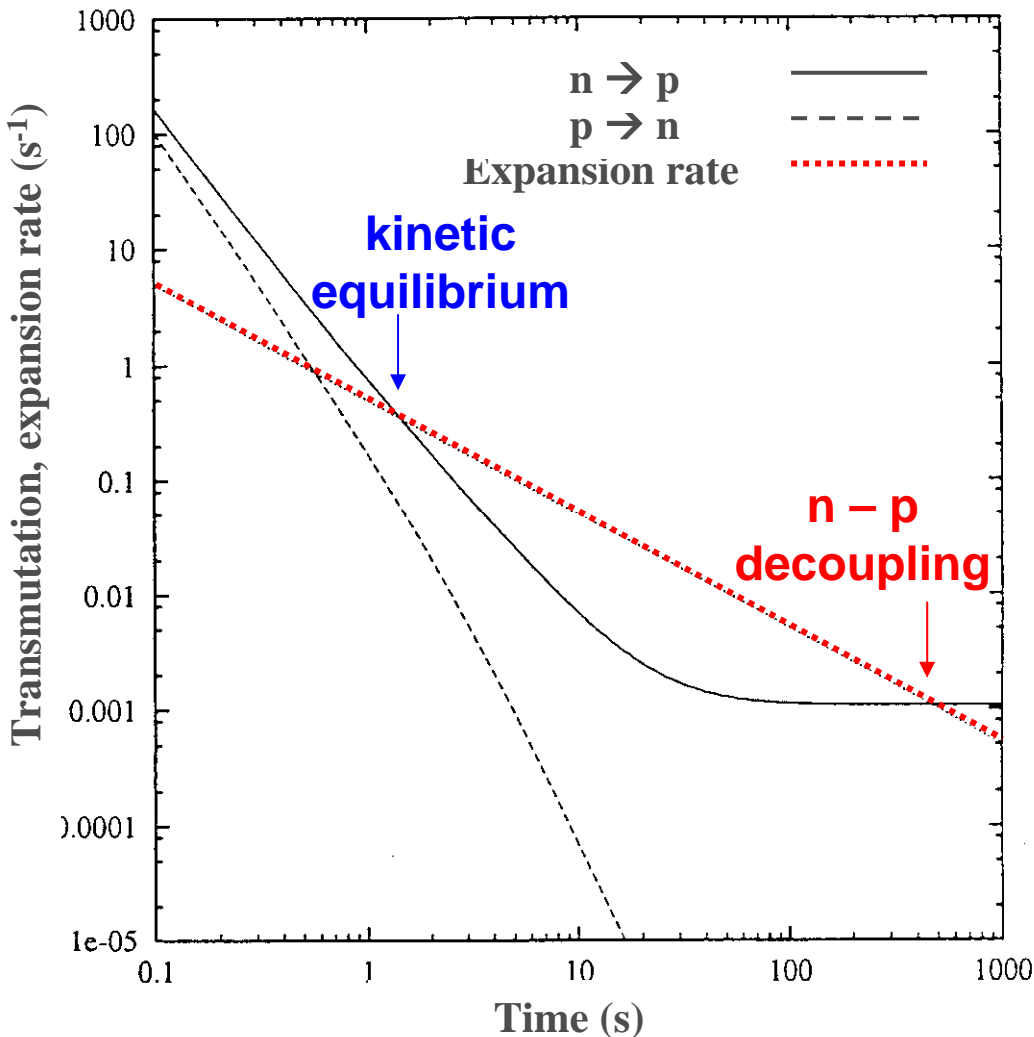
Governs the primordial plasma all along its evolution

Universe is mainly made of photons, ν , $\bar{\nu}$, e^+ , e^- , p and n

Nucleosynthesis (-I)

$T \sim 10^{10} \text{K} \Leftrightarrow 1 \text{ MeV} : \text{équilibrum } p \leftrightarrow n$

Nuclei heavier than $A = 1$ cannot exist : only p and n



Equilibrium $p \leftrightarrow n$ as long as
reaction rate $>$ expansion rate

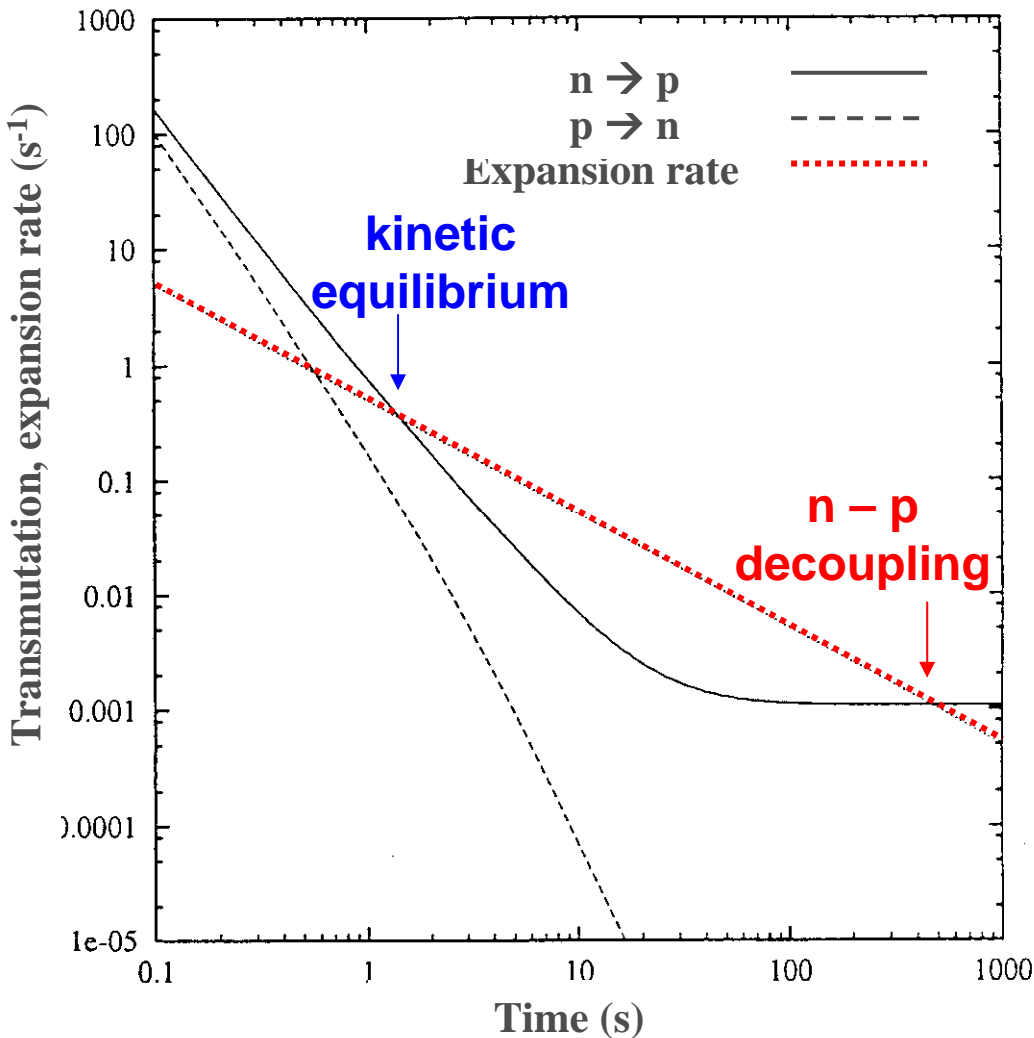
$$N_n/N_p = \exp(-Q/kT)$$

$$Q = M_n - M_p = 1.29 \text{ MeV}$$

n, p = 2 quantum states of the
same particle: the nucleus

Transmutation rate $n \rightarrow p$

Expansion rate $>$ reaction rate
 \rightarrow decoupling $\rightarrow n, p = 2$ different particles
 \rightarrow different evolutions



Decrease of the
neutron mass fraction

$t \sim 1000$ s

Reaction rate $>$ expansion rate
but p is stable
and $n \rightarrow p$
with $\tau_{1/2} = n$ lifetime in vacuum

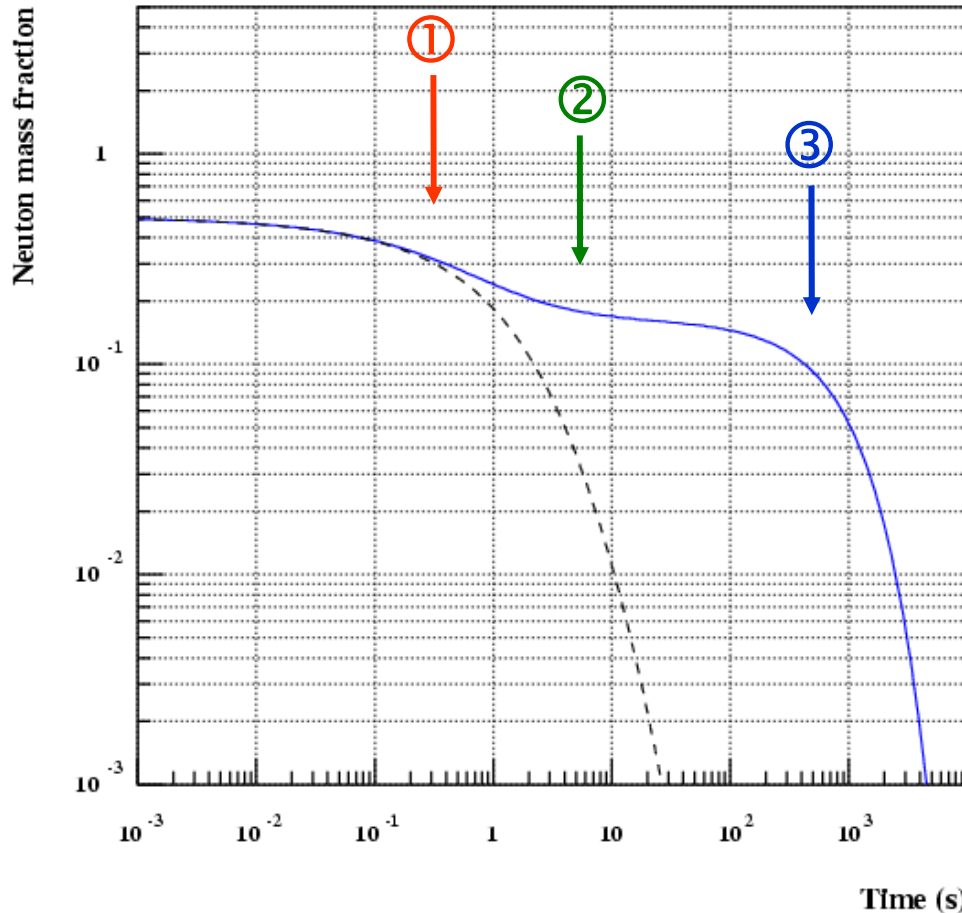
As soon as the more stable state
becomes heavier nuclei

\rightarrow D formation

...

First steps of the BBN

Neutron mass fraction evolution



1. $n \leftrightarrow p$ decoupling
2. free decrease of n
3. D formation

----- $n \leftrightarrow p$ equilibrium
————— Exact calculation

High T , minimal free energy states
 \rightarrow n and p

Low T , minimal free energy states
 \rightarrow heavier nuclei

Because of the low nucleon density
 \rightarrow only 2 particle reactions

Primordial nucleosynthesis = succession of 2 particle reactions

The first one = D formation

n are preserved because trapped in the D nuclei

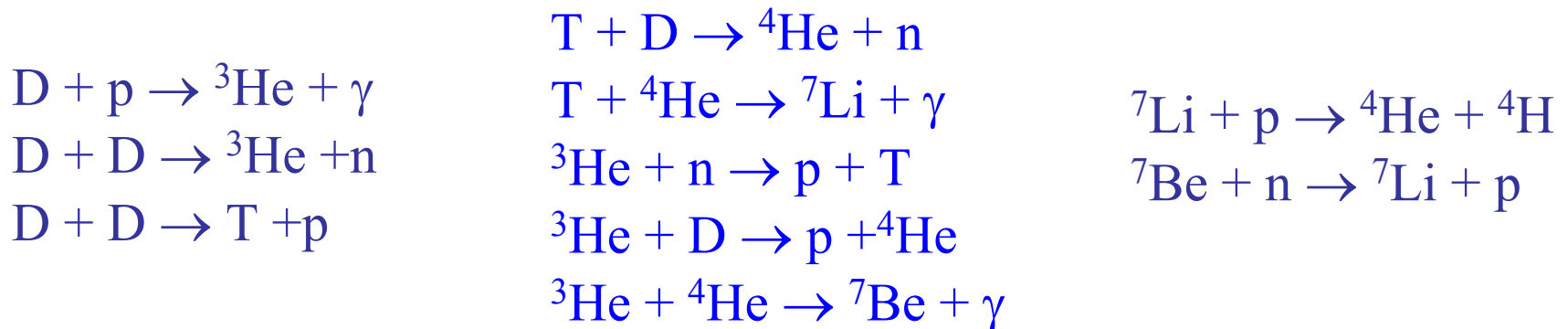
Primordial nucleosynthesis (I)

Neutrons decay until T is low enough: the $n+p \rightarrow D+\gamma$ reaction becomes faster than the deuterium photodisintegration : $D+\gamma \rightarrow n+p$ ($Q = -2.2 \text{ MeV}$)
D energy is smaller than n + p energy

Nuclei are forming : nucleosynthesis begins

Then, $t \approx 3 \text{ mn}$, $T \approx 10^9 \text{ K}$ and N_n has decreased to $N_n/N_p \approx 0.1$

After enough D abundance increase, ^4He is quickly formed via T and ^3He



**No stable nuclei with $A=5$ & $A=8 \rightarrow$ primordial nucleosynthesis limits
only some traces of ^7Li , ^7Be are formed**

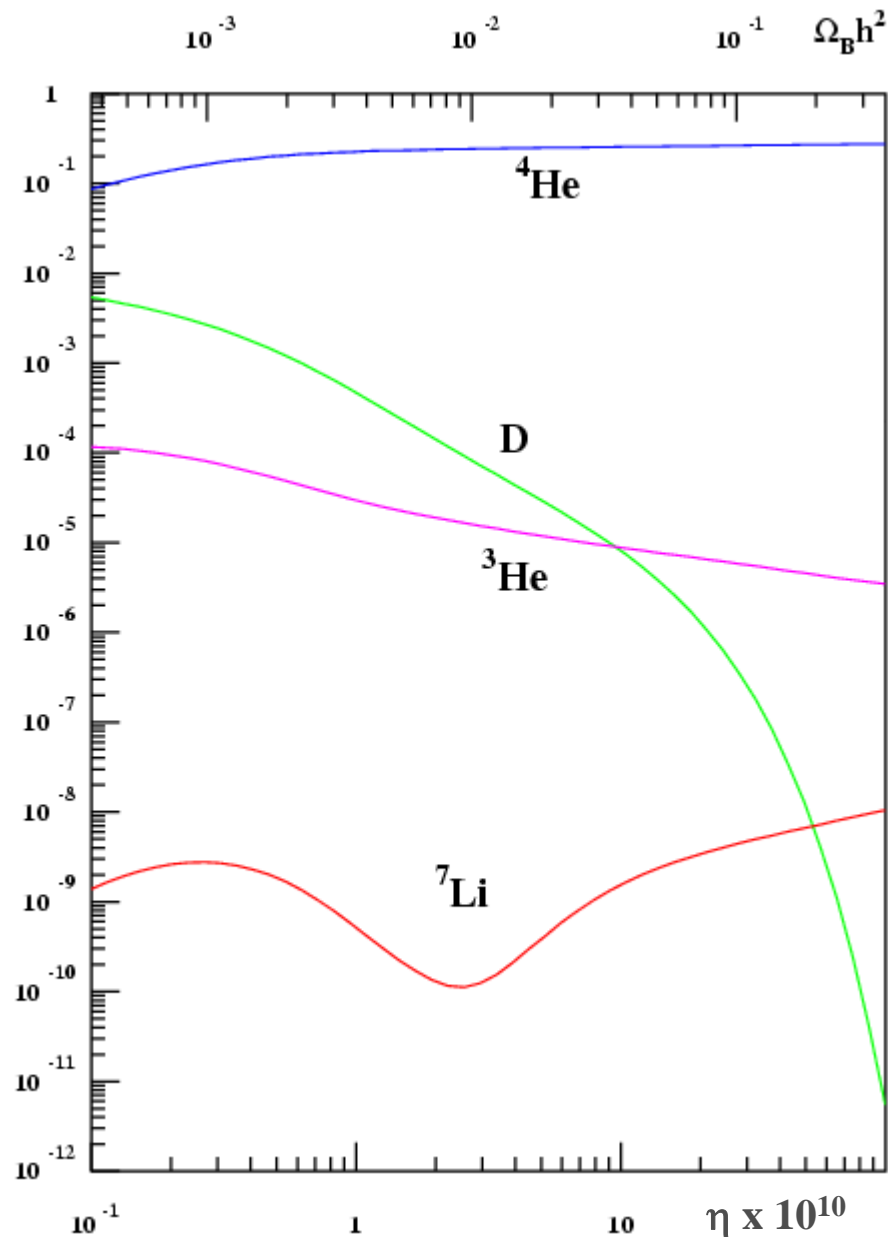
Primordial nucleosynthesis results

Succession of nuclear reactions out of kinetic equilibrium conditions

Light nuclei formation D, ^3He , ^4He and some traces of ^7Li , ^7Be

One important parameter of SBBN : the baryonic density η

Thanks to deuterium, η can be safely restricted to the 10^{-10} to 10^{-9} range



Baryonic density determination

□ Standard primordial Nucleosynthesis (SBBN)

Comparison between calculated and observed primordial abundances of ${}^4\text{He}$, D , ${}^3\text{He}$, ${}^7\text{Li} \Rightarrow \eta$ (ou $\Omega_{\text{B}}h^2$)

Needs:

✓ *precise observational data (primordials from extrapolated observations via chemical evolution models of the galaxies!)*

✓ *precise nuclear data*

□ Anisotropies of the cosmological microwave background (CMB)

BOOMERanG, CBI, DASI, MAXIMA, ARCHEOPS, WMAP, Planck.....

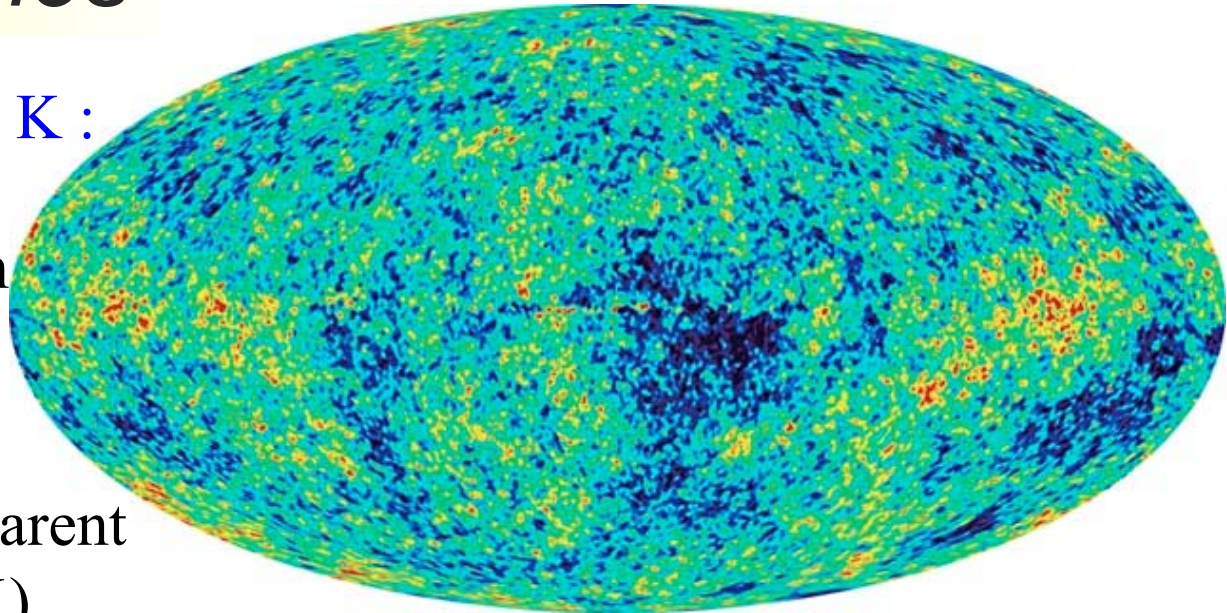
□ Primordial abundance observations : **Lyman- α** HI and HeII absorption lines

(looking in the quasars direction)

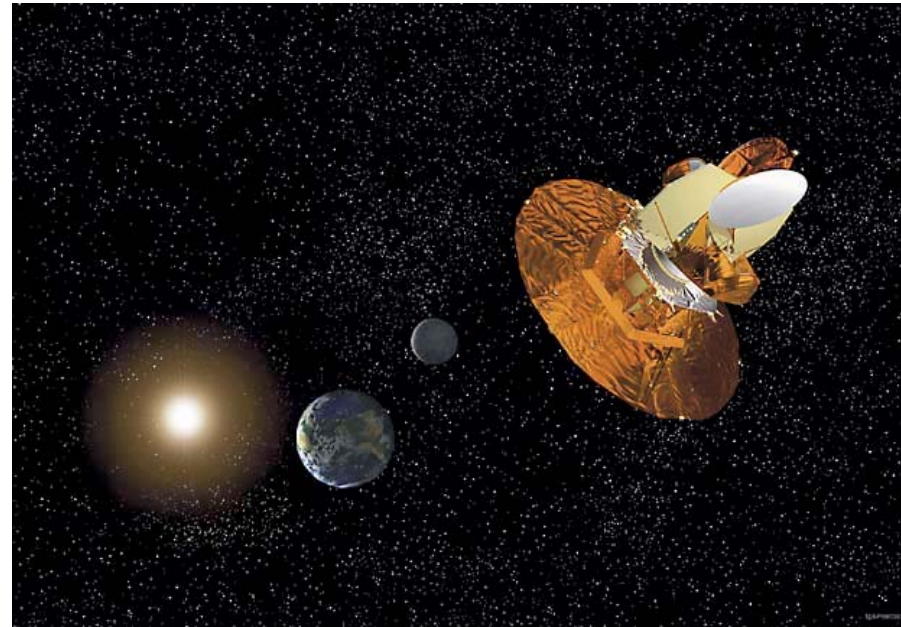
CMB anisotropies

At $t \approx 0.3$ My and $T \approx 3000$ K :
recombination:
 e^- and nuclei combination
to nuclei to form neutral
atoms.

Universe becomes transparent
(now $T = 2.725$ K)
(microwave energy range)



*After COBE, new generations of instruments (BOOMERanG, CBI, DASI, MAXIMA, ARCHEOPS)
Now **WMAP** and **Planck/HFI** dedicated to anisotropy studies of the CMB.*



CMB anisotropies

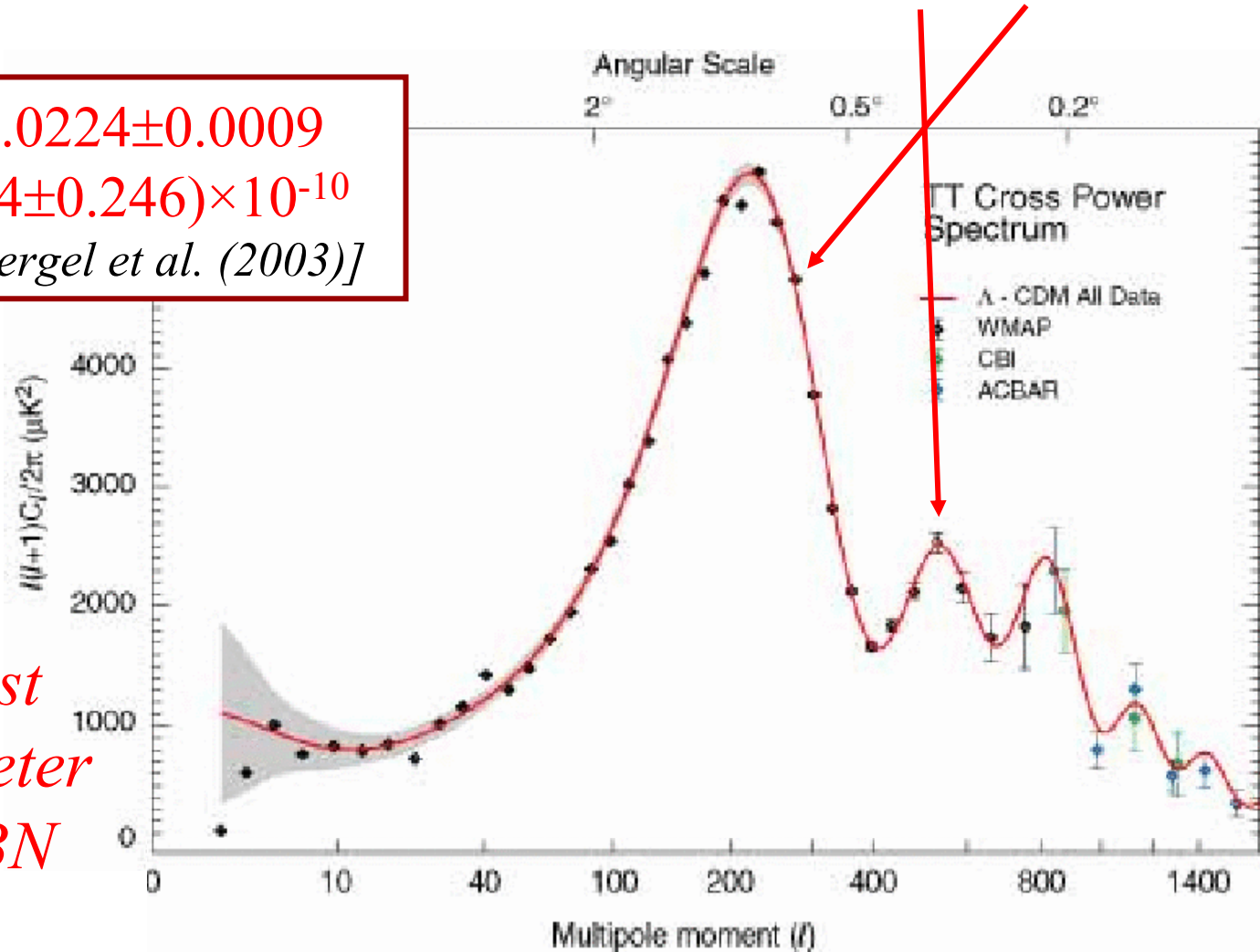
Spatial fluctuations spectrum of CMB

- Geometry ($\Omega_T \approx 1$), 1st peak
- Ω_b (2nd/1st peaks)

$$\Omega_b h^2 = 0.0224 \pm 0.0009$$
$$\eta = (6.134 \pm 0.246) \times 10^{-10}$$

[WMAP: Spergel et al. (2003)]

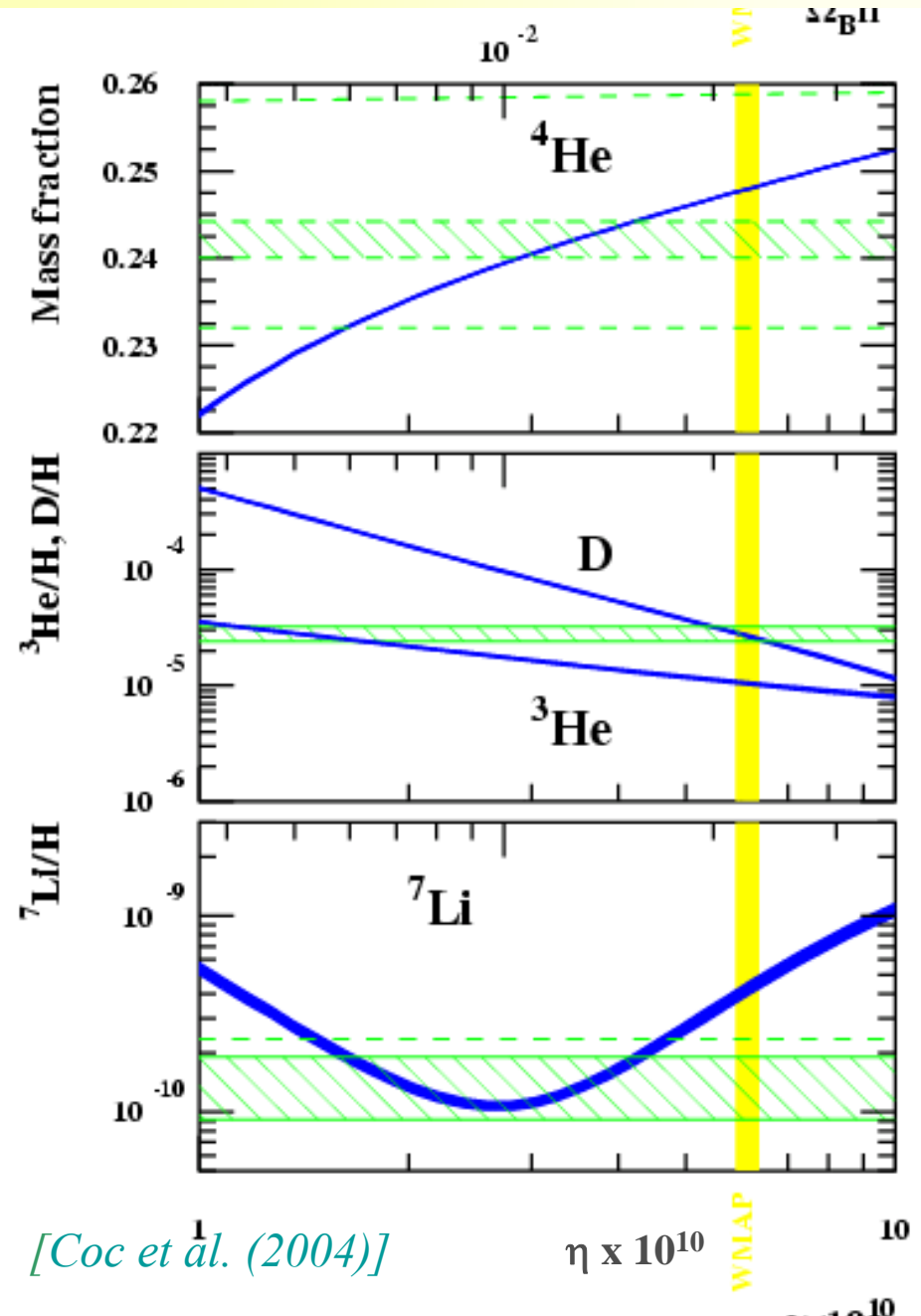
Fix the last free parameter of the SBBN



Observed vs calculated abundances

Primordial abundances

- H : more than 75% in mass
- ^4He : blue compact galaxy observations \rightarrow 24% in mass
- D : local interstellar medium; protosolar clouds; cosmological clouds
- ^7Li : low metallicity stars
- $\Omega_B h^2 = 0.0224 \pm 0.0009$
[WMAP: Spergel et al. (2003)]



[Coc et al. (2004)]

Primordial nucleosynthesis

To find a model in agreement with observation data :
a new theory is needed, « beyond the standard model »

Then?

Universe is around 3 minutes old ;

It is mainly a plasma made of photons, ν , $\bar{\nu}$, e^+ , e^- , p , ${}^4\text{He}$ and traces

Temperature and density are too low to allow other reactions but too high : neutral atoms, molecules, cannot be formed

(Coulomb forces are too small compared to thermal energy)

Universe will expand during $\sim 10^9$ years before the gas may condensate because of tiny density variations, and form proto stars; before nuclear reactions may be initiated (beginning of a star life);

and before the stellar nucleosynthesis will begin

Metallicity and stars

Equilibrium condition: during the whole life of the star

In astrophysics:

“metals” = everything heavier than helium

Metallicity = “metal” mass fraction (Z)

Solar metallicity : $Z_{\odot} \approx 0.013$

A **first generation star** was formed with primordial gas
metallicity ~ 0 : $\sim 3/4$ of H and $1/4$ of ^4He in mass

Stars formed with metal enriched gas : **second generation star**
metallicity as high as the gas was metal enriched

Always a small amount : $< 3\%$

Important observational data

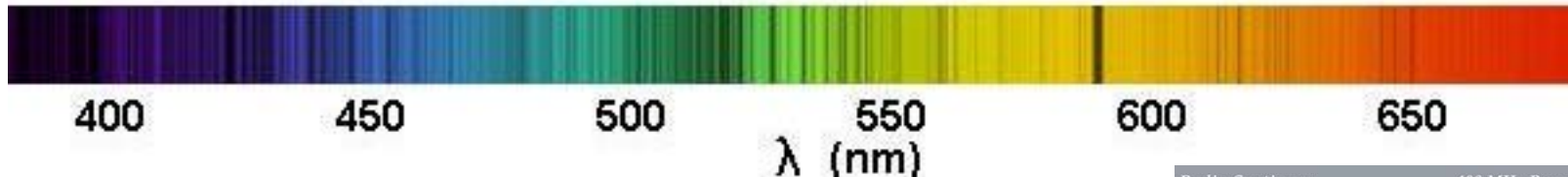
Determination of elemental and isotopic abundances

*from reachable material (only a tiny amount of the universe) ;
solar photosphere observations ; chondrites collection ;
cosmic radiation observations*

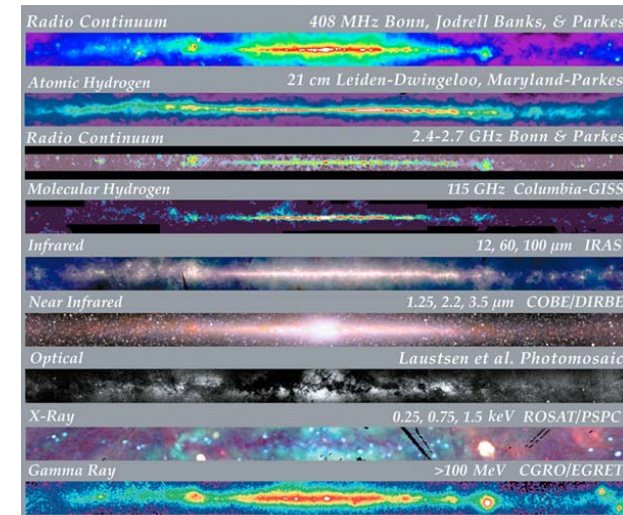
Variety of sources with quite uniform composition

Electromagnetic observables

Traditionally : radio, visible wave lengths



*in addition now : IR, UV, X, γ
(until PeV and beyond)*



→ Elemental abundances relative to H

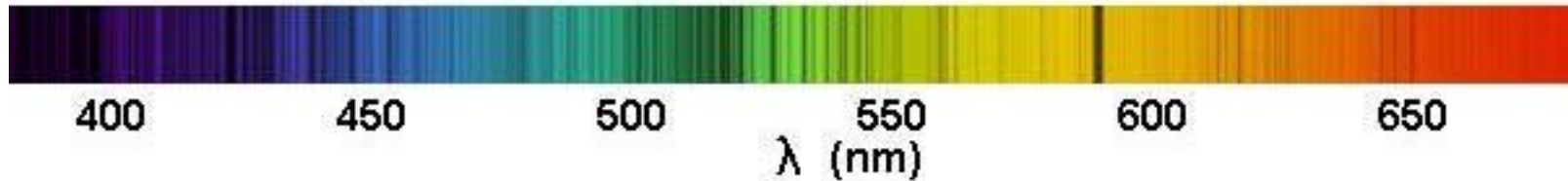
→ Isotopic ratio of abundances

→ Temperature and Luminosity

A new observable : neutrinos

Important observational data

Stellar electromagnetic spectra



Continuous spectrum + (dark or intensive) lines

Continuous spectrum : *black body distribution* → **Temperature (T)**

dark lines : *absorption* (→ *some material lies between the emitting star and the detector*)

intensive lines : *emission* (→ *some processes excite atoms and/or molecules on the stellar surface, they deexcite by radiation emission*)

→ **Elemental abundances relative to H**

→ **Isotopic abundance ratios**

Luminosity (L) → *amount of energy radiated per time unit : deduced from the magnitude (amount of energy received per time and surface unit on top of the atmosphere) and the distance of the star*

→ **Star classification : luminosity vs temperature**

Hertzprung-Russell Diagram

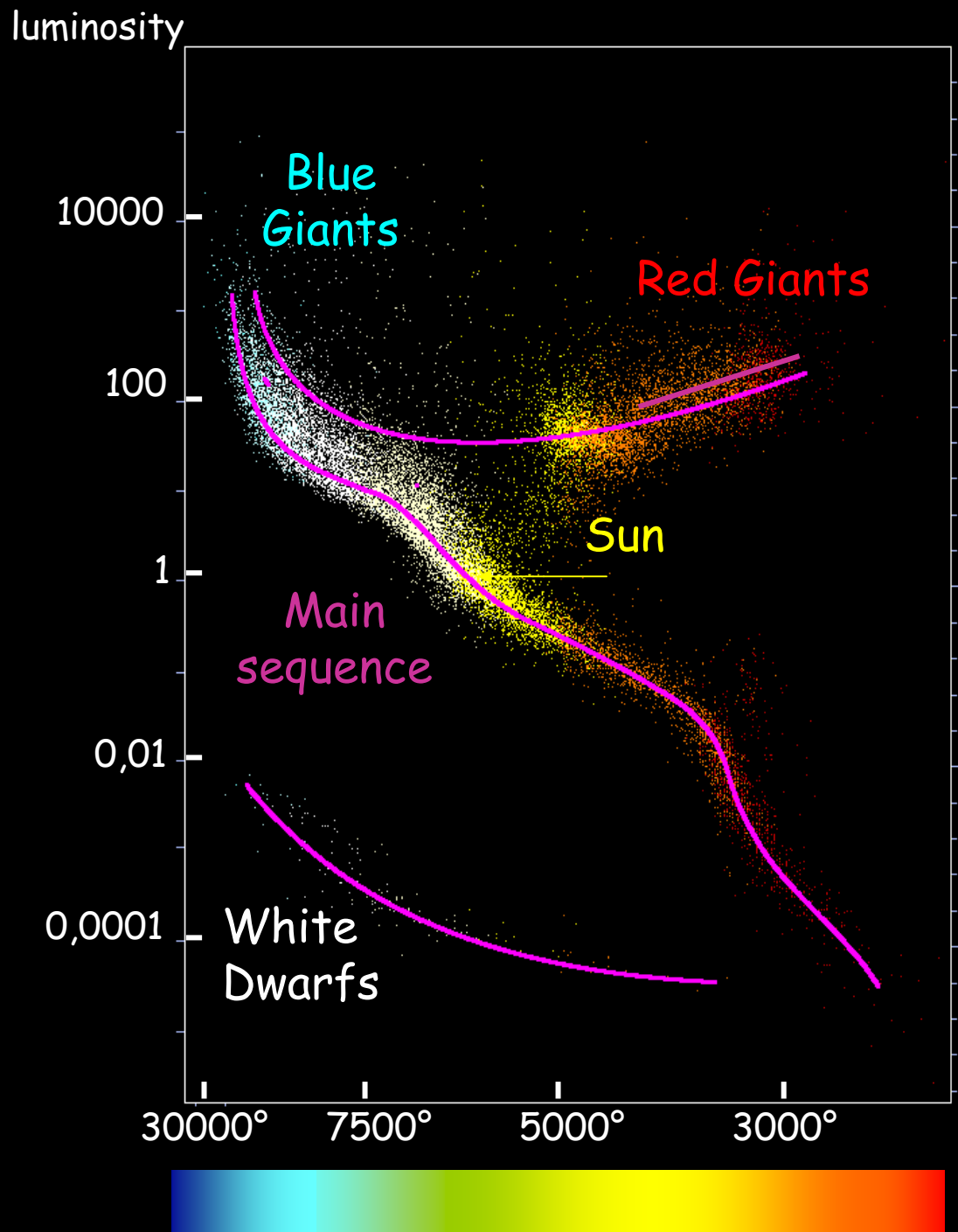
temperature -
luminosity

Correlation lines :
→ star classification

$$T \Leftrightarrow E = k_B T$$

with $k_B = 8.62 \cdot 10^{-8} \text{ keV/K}$

Surface temperature
Color



HR Diagram

*temperature -
luminosity*

At a given time

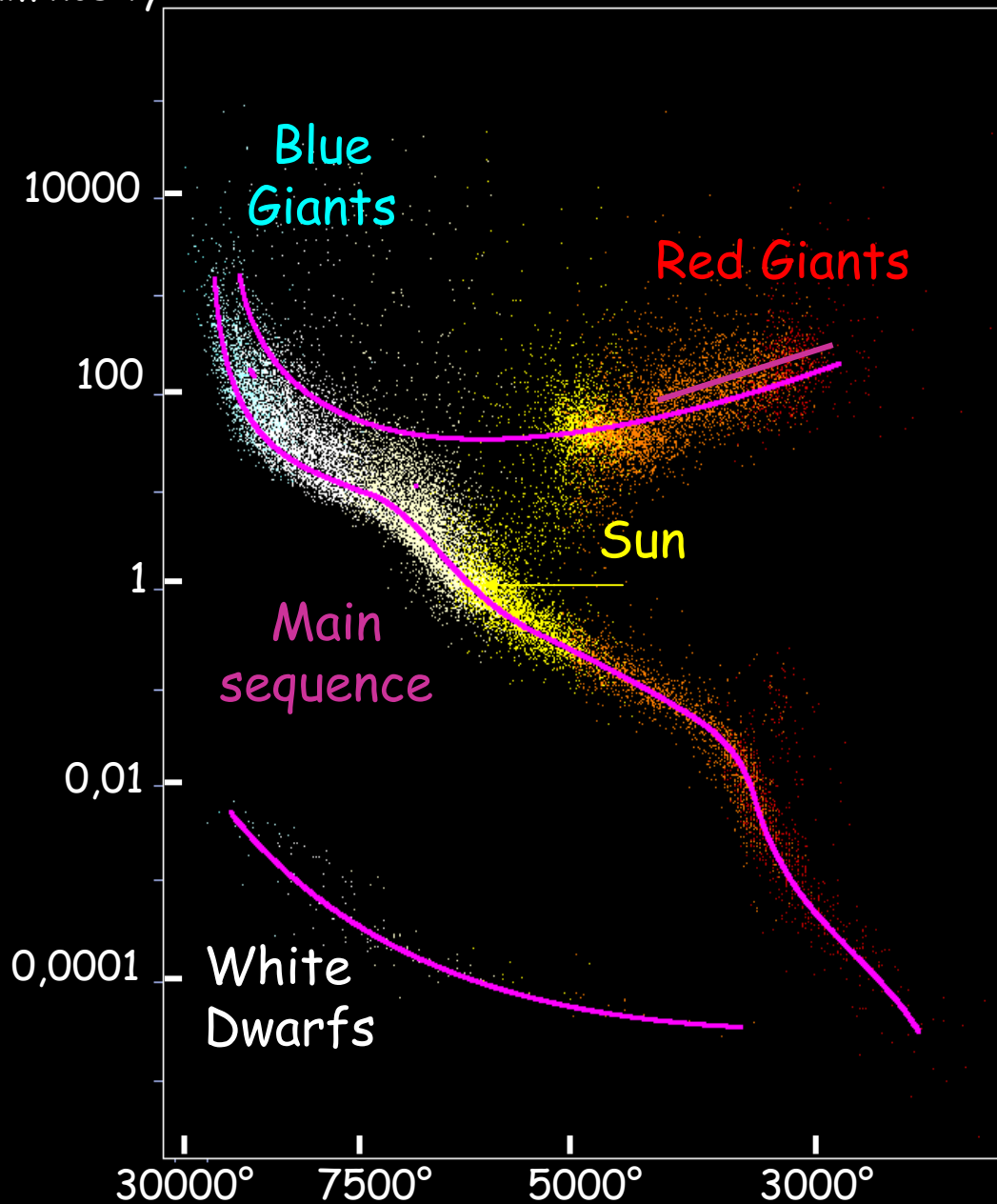
Interpretation in terms
of stellar population :

*Relative proportion of
the various kinds of star*

Surface temperature

Color

luminosity



HR Diagram

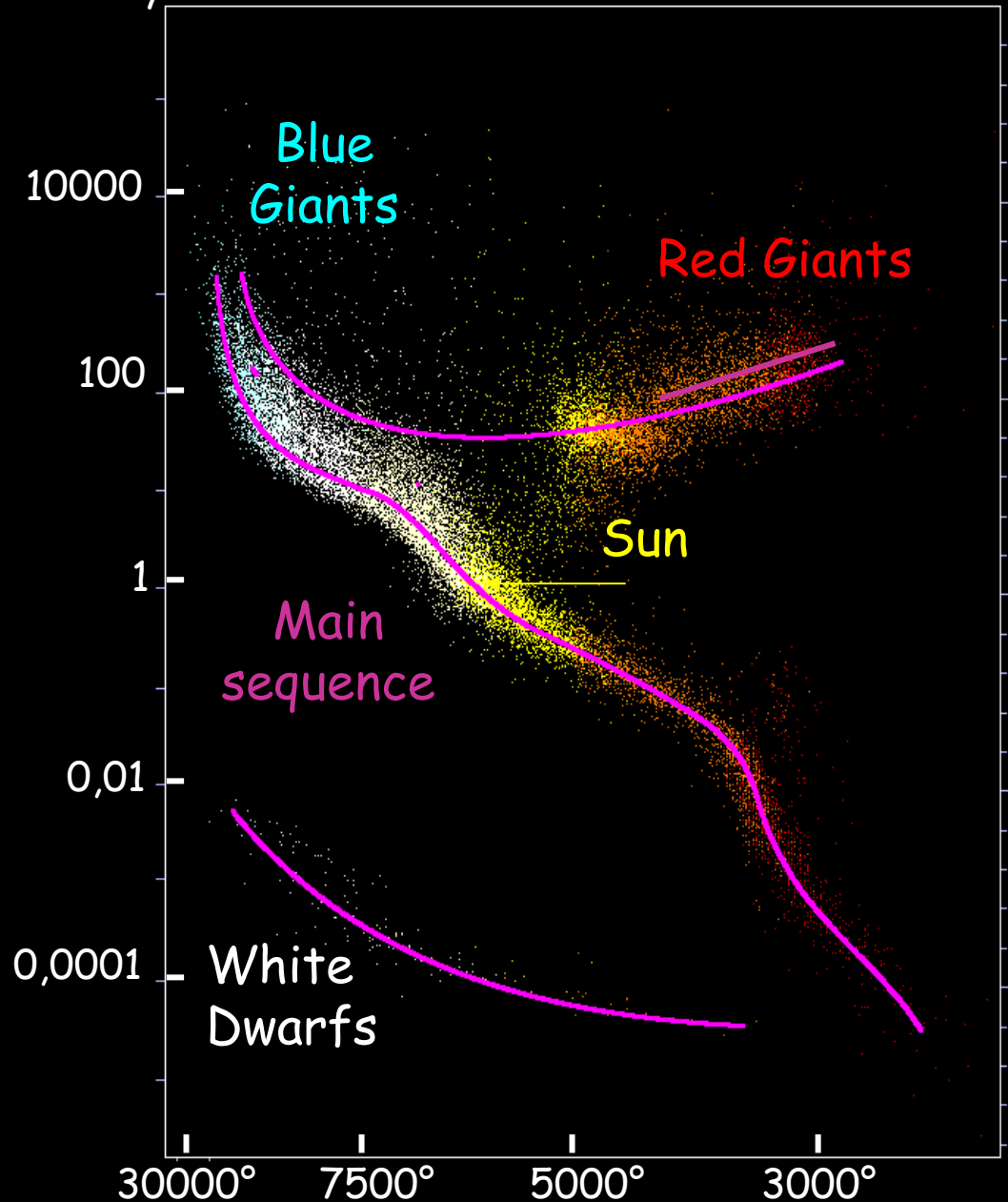
temperature -
luminosity

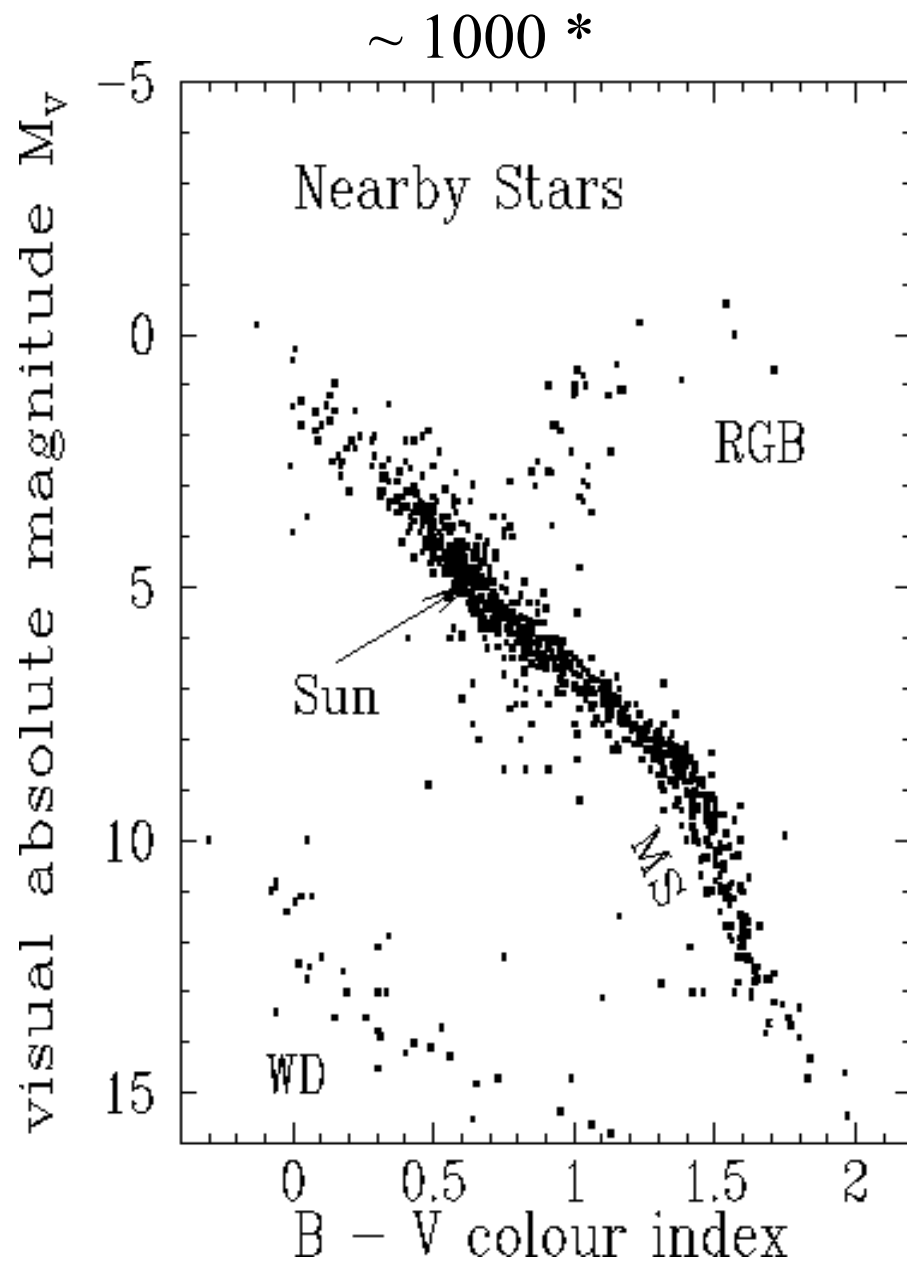
In terms of star evolution

Star classification :
correlation lines

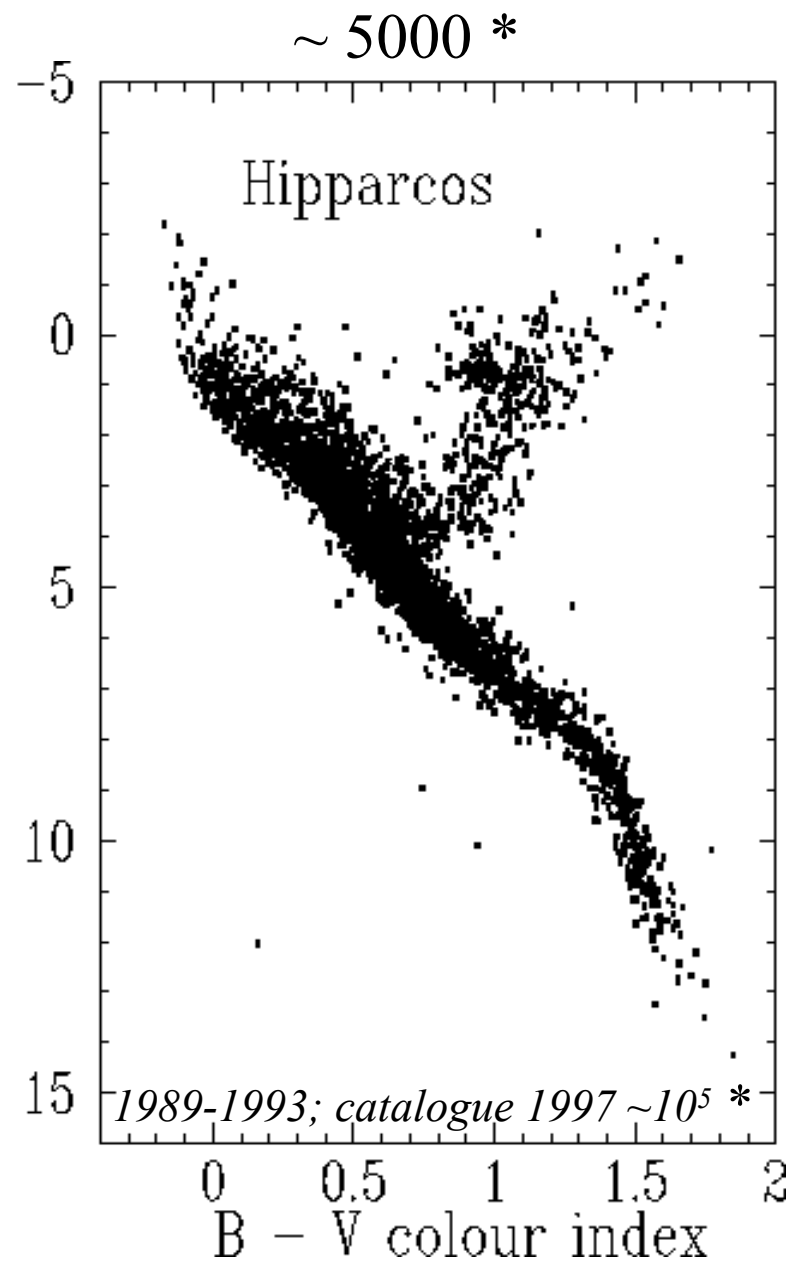
Surface temperature
Color

luminosity

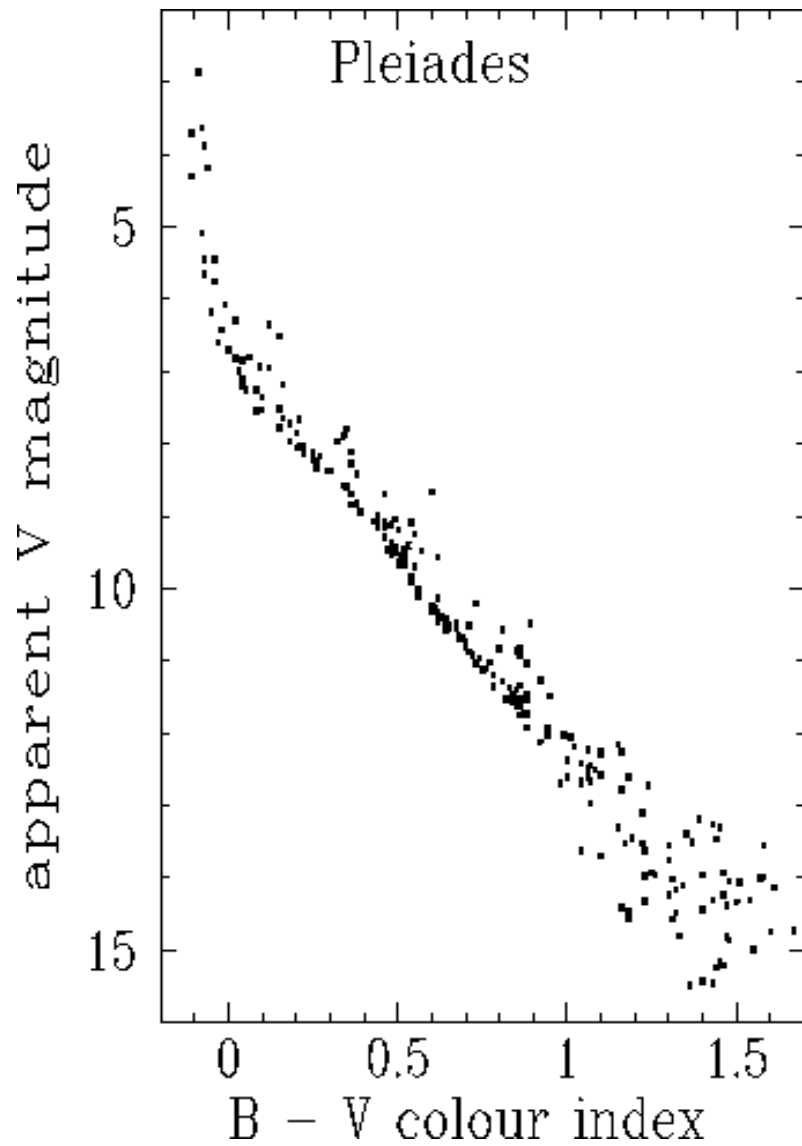




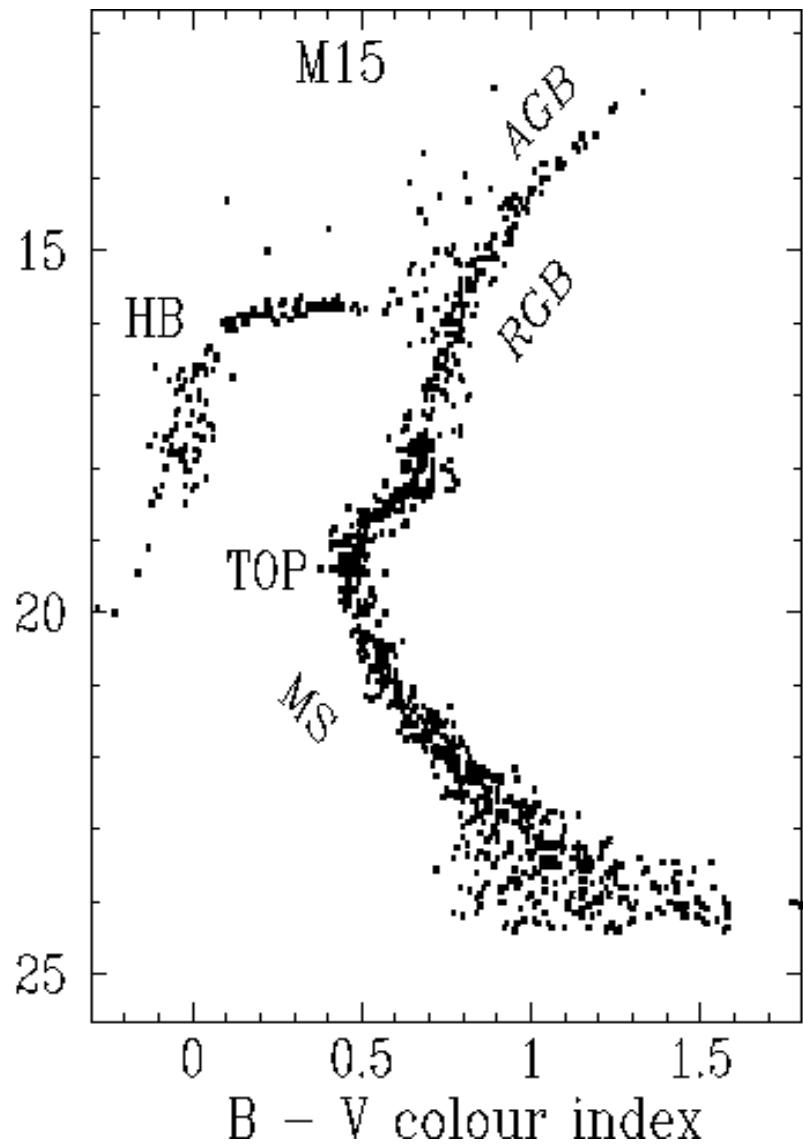
$D < 20pc$



$60 pc < D < 70pc$



Young galactic cluster



Old globular cluster

Star formation (-1)

Orion nebula



After BBN, the primordial gas contains mainly, ν , $\bar{\nu}$, e^+ , e^- , p and light nuclei

Tiny density variations allow gas concentrations : massive, dense and cold molecular clouds

Stars are formed by gas condensation because of the gravitation

Initially : cold gas

gravitation induces the free-fall collapse of gas fragment

($\tau_{\text{ff}} \sim$ hour for the \odot)

→ proto stellar gas fragment

Star formation(0)

Free-fall collapse → progressive increase of the internal pressure
→ particle ionization
→ temperature increase due to opacity increase
→ gas collapse slows down

Until the proto-star reaches a hydrostatic quasi-equilibrium state :
a large amount of gas is totally ionized in the core

Then new contraction with new temperature increase
until density and temperature high enough to initiate nuclear reactions

Only if the **mass** of the gas fragment is **high enough**
otherwise, it will never radiate enough energy :
the gas fragment remains in a quasi equilibrium state
never a star

Star formation

When the gas fragment is massive enough

T increases up to $\sim 15 \times 10^6$ ($T_6 = 15$)

Nuclear reactions may occur \rightarrow **beginning of the star life**

Mechanical equilibrium : stationary state

gravitation + internal pressure = 0

stationary evolution \rightarrow scaling time \gg free-fall time

Energetical equilibrium : stationary state

Nuclear energy compensates energy losses in each point of the star

radiation + convection + ν + thermal conduction

+ mass losses (stellar winds)

stationary \rightarrow long scaling time \gg Helmholtz - Kelvin time

(time needed to release all the emitted power without nuclear energy source :

$\tau_{HK} \sim 10^{10}$ years for the $\odot \gg \tau_{ff}$)

D burning stage ...

Mainly H and He ; but some D traces

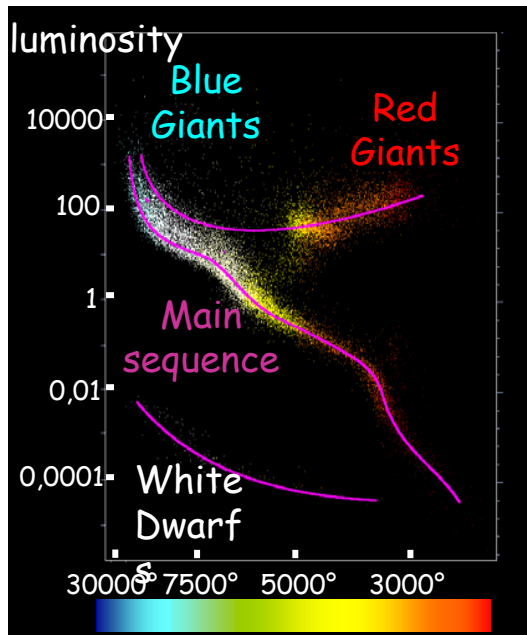
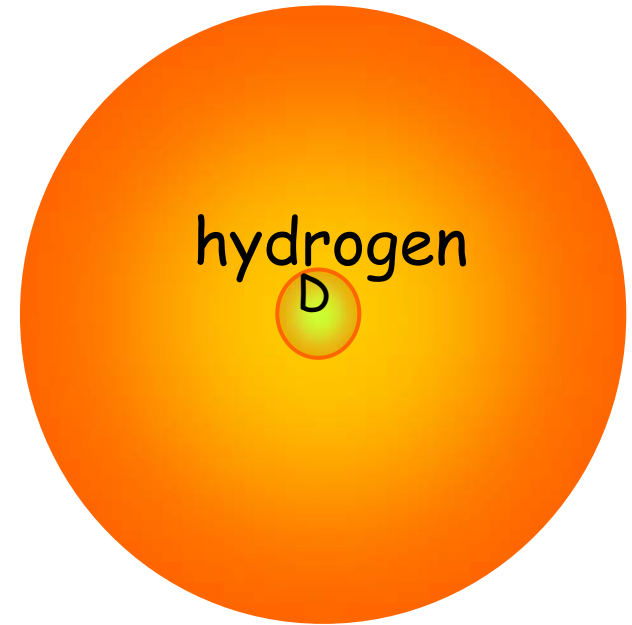
(a few 10^{-4} in mass compared to H)

D is not very stable

→ D photodissociation

1st possible nuclear reaction

1st signals of the star



Then gravitationnal contraction

The star is getting closer to the main sequence

At the beginning of its life,
the star is on the main sequence
in the HR diagram

Stellar evolution (0)

Reference model framework
(approximation)

Chemical homogeneity (when formed)

Stationnary state

No magnetic field

No rotation



Allows , in first approximation,
a description of stellar evolution

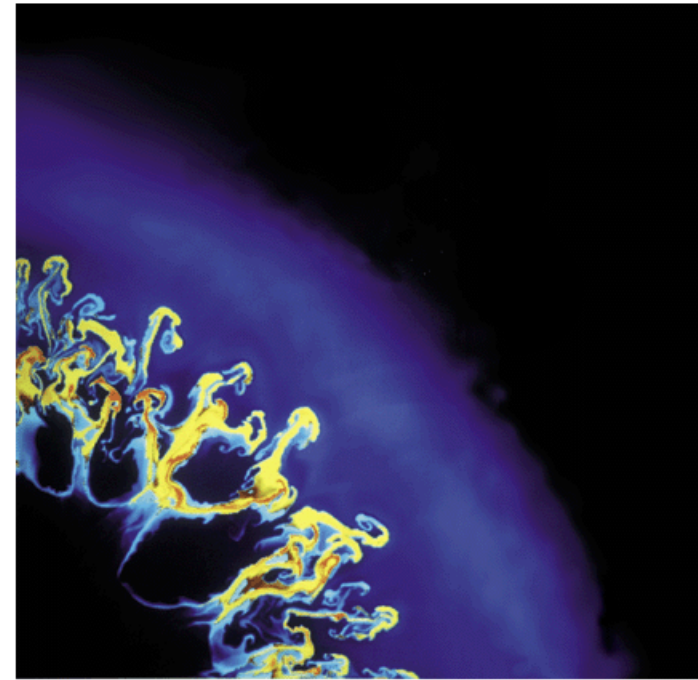
$\sim 3/4$ of H and $1/4$ of ^4He

Ionized particles \rightarrow Electromagnetic field \rightarrow Coulomb barrier

First allowed reactions between H particles

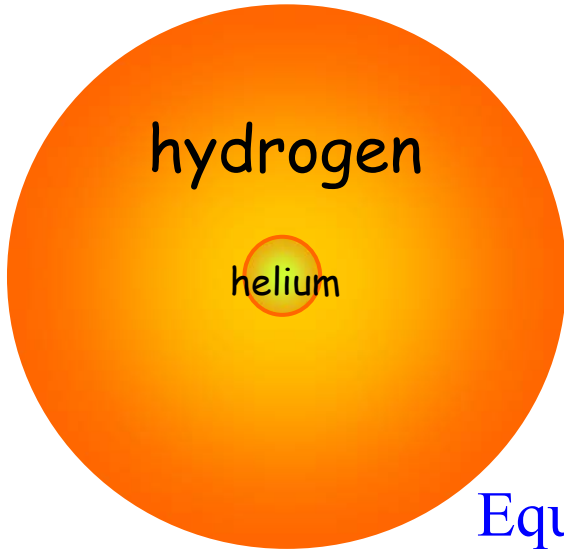
Hydrogen burning phase

*Stars on the main sequence
are in the hydrogen burning stage*



Supernova 1987A provided strong evidence of turbulence emanating from the core of the exploded star because core materials were observed well before they were predicted. The turbulence caused mixing among the layers and greatly complicated the tidy "onion" model of dying stars. [Image reproduced from Muller, Fryxell, and Arnett, *Astronomy & Astrophysics* 251, 505 (1991).]

H burning (Main sequence)



Mainly H and He
First enabled reactions :
Hydrogen fusion
→ He formation

Star is on the Main Sequence

Equilibrium : gravitation vs radiation

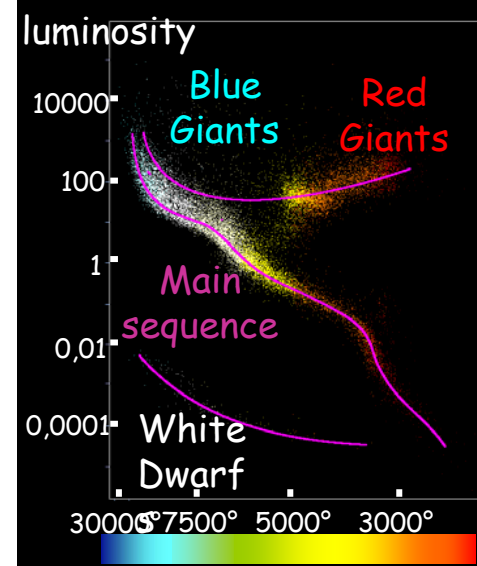
Increase of the central He concentration

≠ kind of burning processes depending on the chemical composition of the star gas (generation) and on its mass

First generation (tiny traces of metal –Li,Be-) and light stars → pp chains
(*equilibrium at $T_6 \sim 15$; no CNO nuclei*)

Second generation and massive stars → CNO cycles
(*equilibrium at $T_6 \gtrsim 20$; \exists CNO nuclei*)

1st generation star observations: mainly H lines



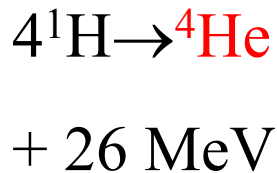
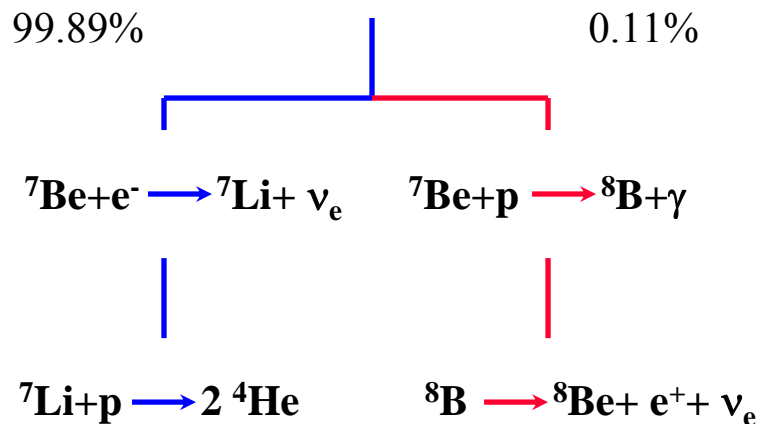
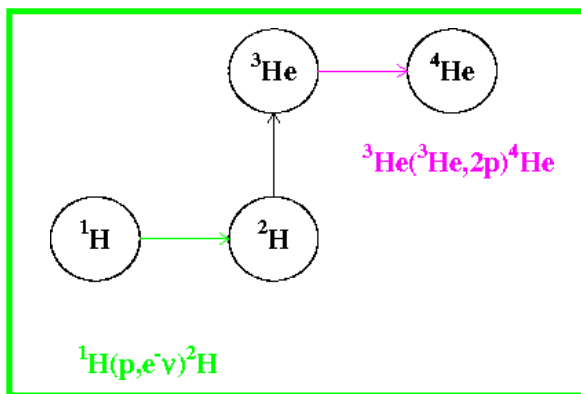
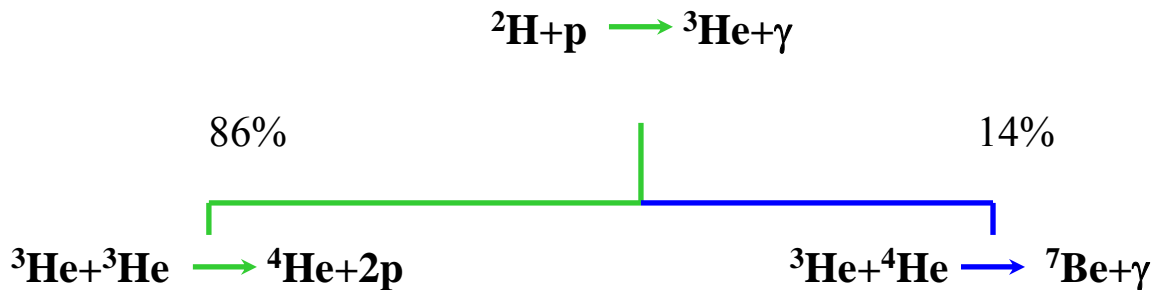
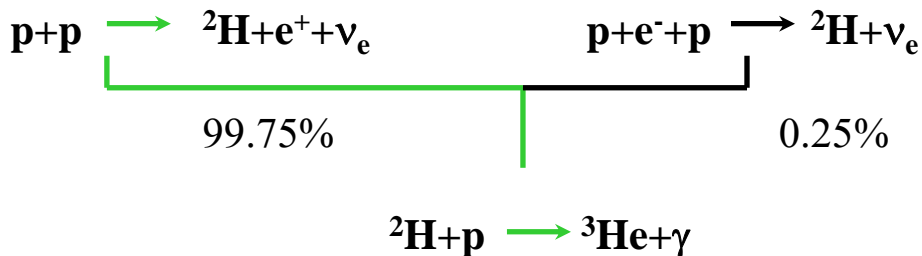
pp chains (H burning)

Weak interaction



→ This combustion takes time

$\sigma \approx 8 \cdot 10^{-27} \text{ b}$
at 5.9 keV
 $\sigma \approx 2 \cdot 10^{-23} \text{ b}$
at 1 MeV



ppI

ppII

ppIII

pp chains (H burning)

Stops because of the lack of stable nuclei with $A = 8$

Central T and ρ too small to enable any 3 body reaction

→ The end of the H burning when H is missing in the core

Time correlation with the mass of the star and its initial composition
(*calculated H burning duration for the Sun : $\sim 10^{10}$ years*)

(*only 10^7 years for a $15 M_{\odot}$ star: 1000 times less*)

→ *a 1st-generation $2-M_{\odot}$ star went out of the H burning stage a long time ago*

As H burning stage is longer than all the others together

→ *such a star also died a long time ago*

→ *only lightest 1st generation stars are still alive*

New gravitational contraction

The star is leaving the main sequence

(cold) CNO cycle (H combustion)

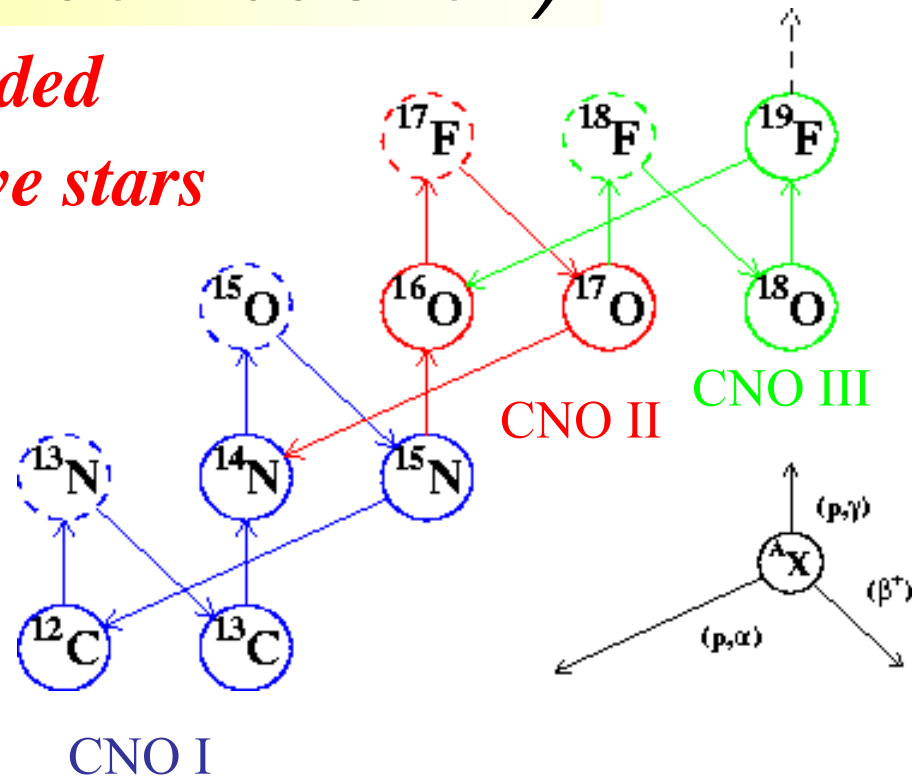
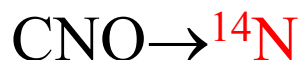
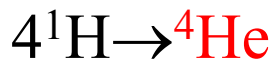
CNO nuclei are needed

2nd generation or massive stars

- $T \gtrsim 20 \cdot 10^6 \text{ K}$
- Via 3 (p, γ), 1 (p, α) reactions and 2 β^+ disintegration
- on CNO (+F) isotopes as catalysts
- $^{14}\text{N}(p,\gamma)^{15}\text{O}$, slowest reaction, it limits the energy production

it induces initial CNO

change in ^{14}N



Only strong interaction: quicker process

\exists other possible « cycles » for H combustion but negligible for energy production: Ne-Na and Mg-Al cycles

Will dominate in stars heavier than $1.8 M_{\odot}$

End of H burning ...

Only a few hydrogen in the core:

- * End of H central burning
- * He core contraction (gravitation)
- * Star contraction (gravitation)...

Core temperature increase
and around the core ;

→ Shell H burning

→ He core increase

→ Star expansion

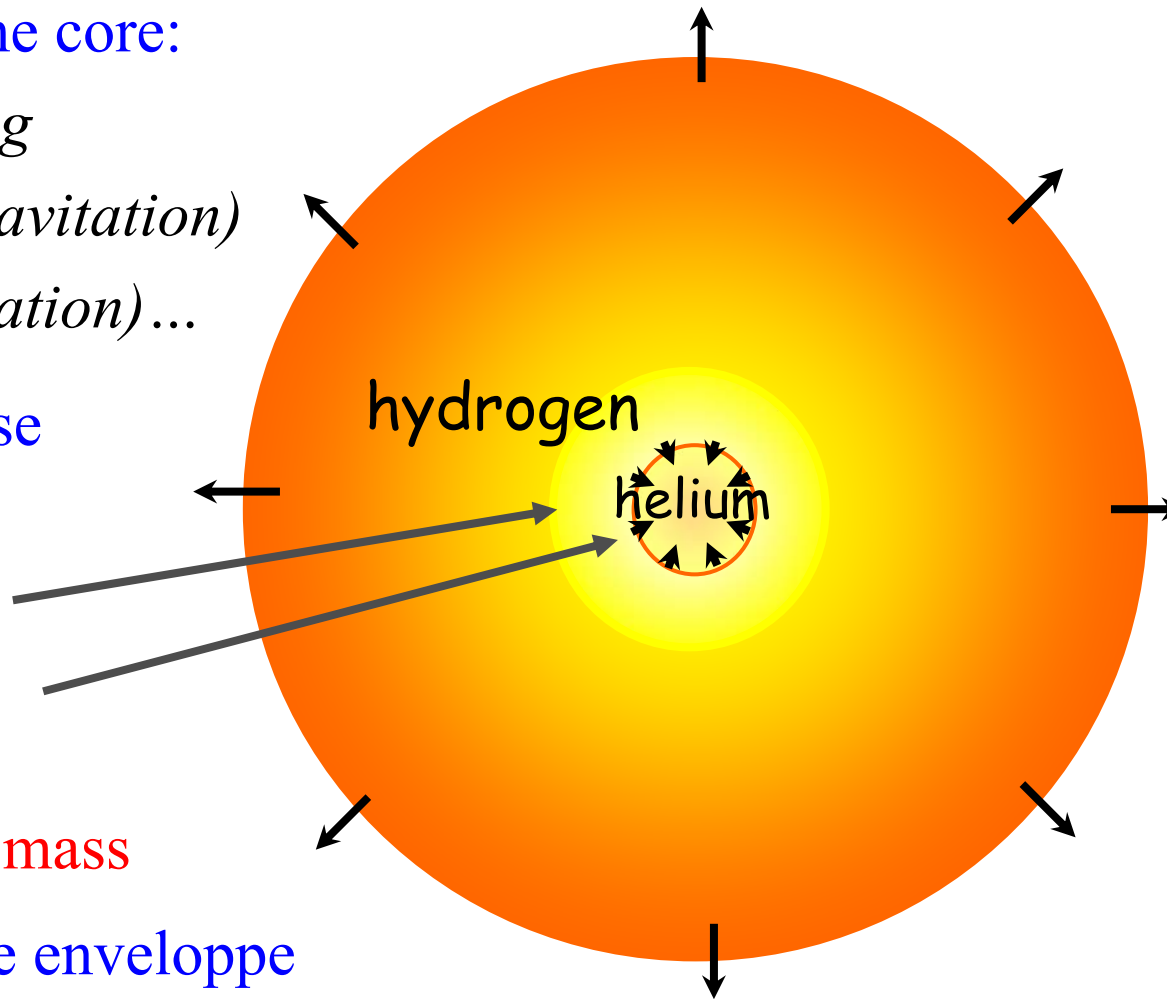
Depends on the star's mass

Convective envelope

Radius increases by a few order of magnitude

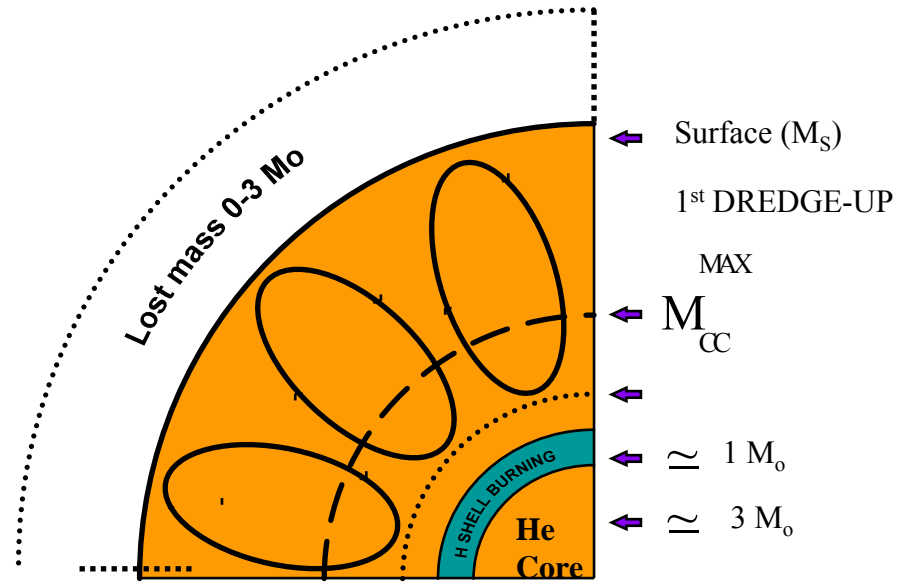
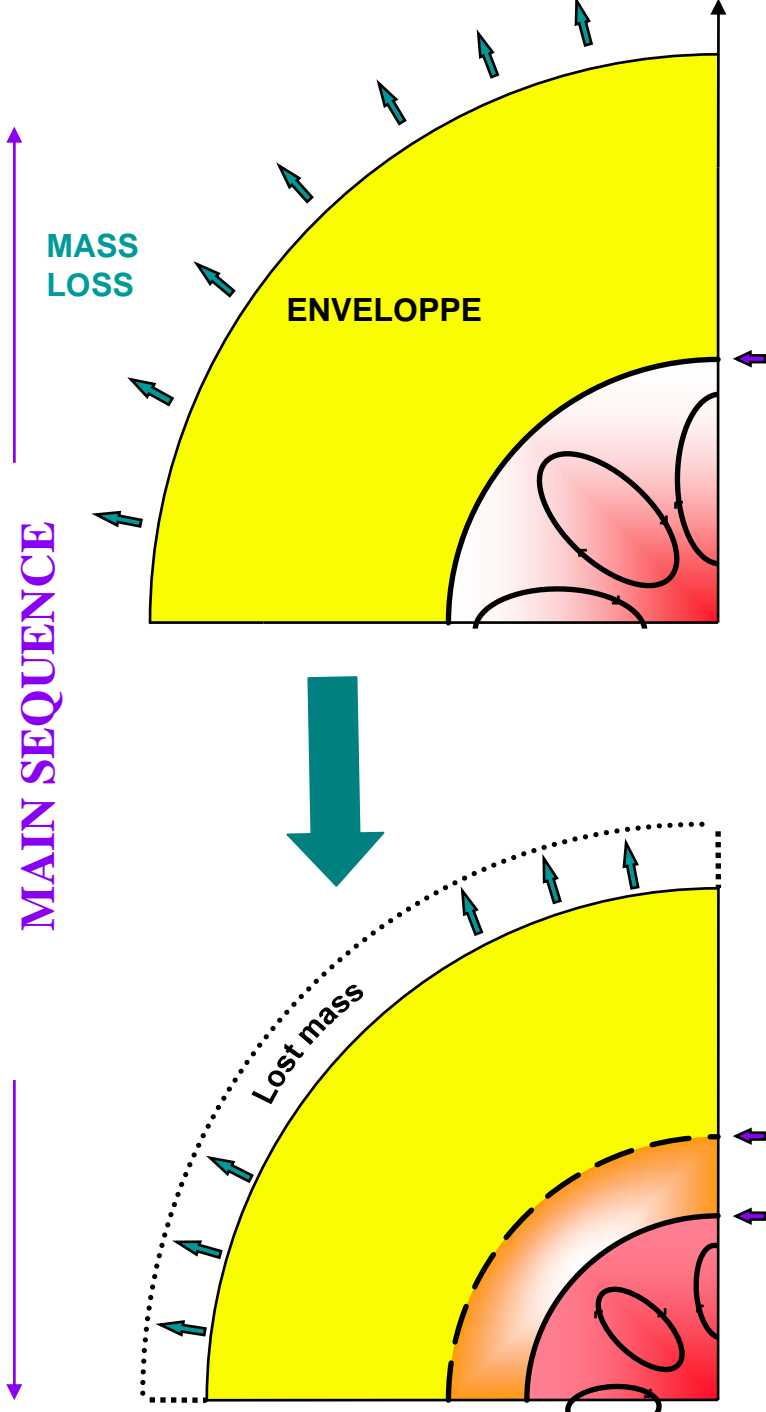
Some H burning ashes are driven to the surface (1st « dredge-up »)

1st generation stars: more He lines may be observed



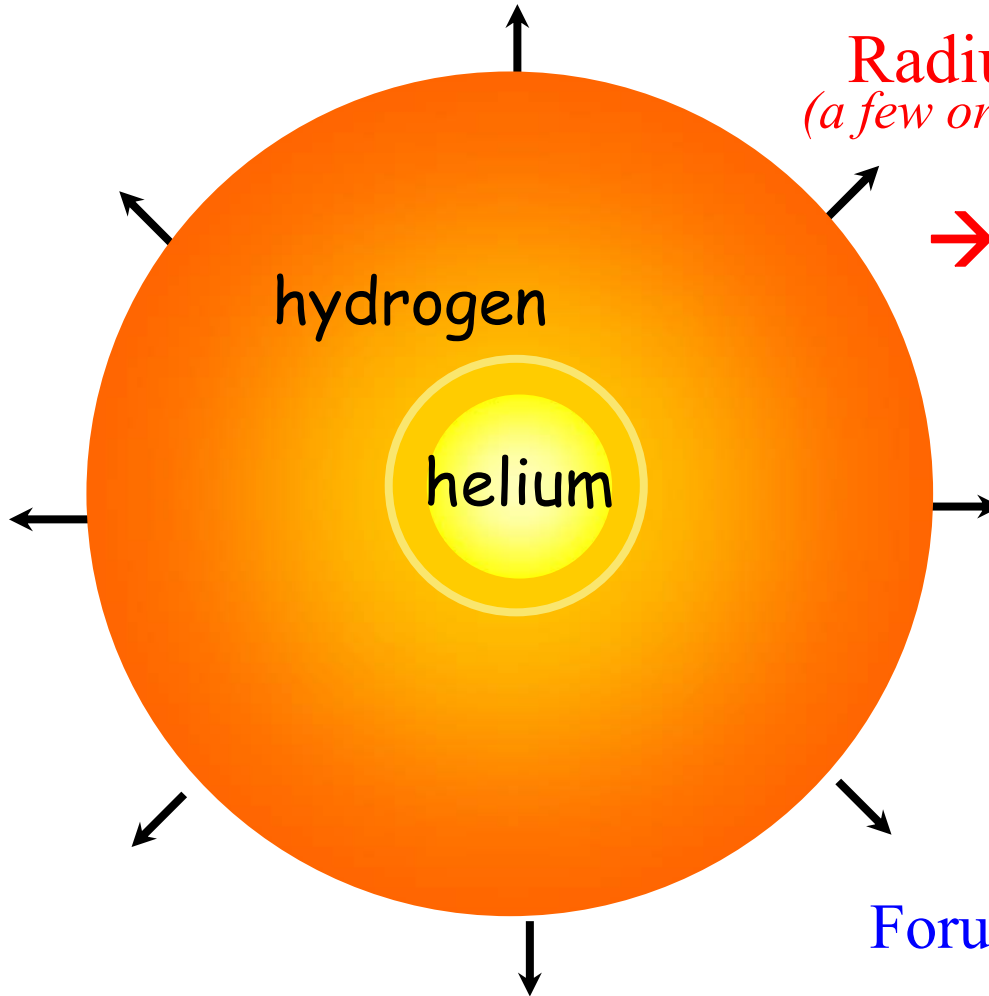
1st dredge-up

M = 15 M_⊙



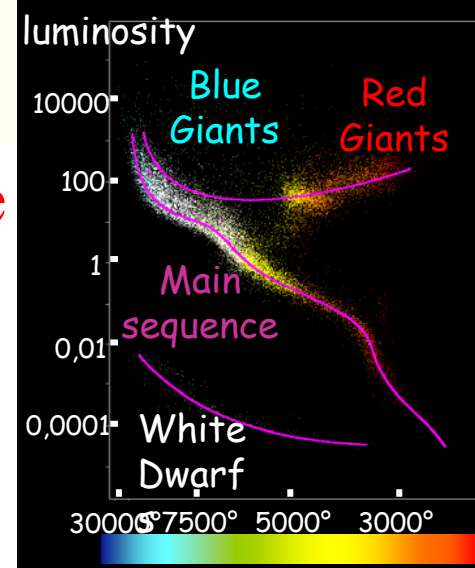
End of H burning \rightarrow He burning

Surface farther from the core: colder temperature



Radius increase
(a few order of magnitude)

\rightarrow **Red Giant**



Central T \nearrow ;

Until nuclear reactions with He
are allowed

\rightarrow He core contraction stops

But no stable $A = 8$ nucleus

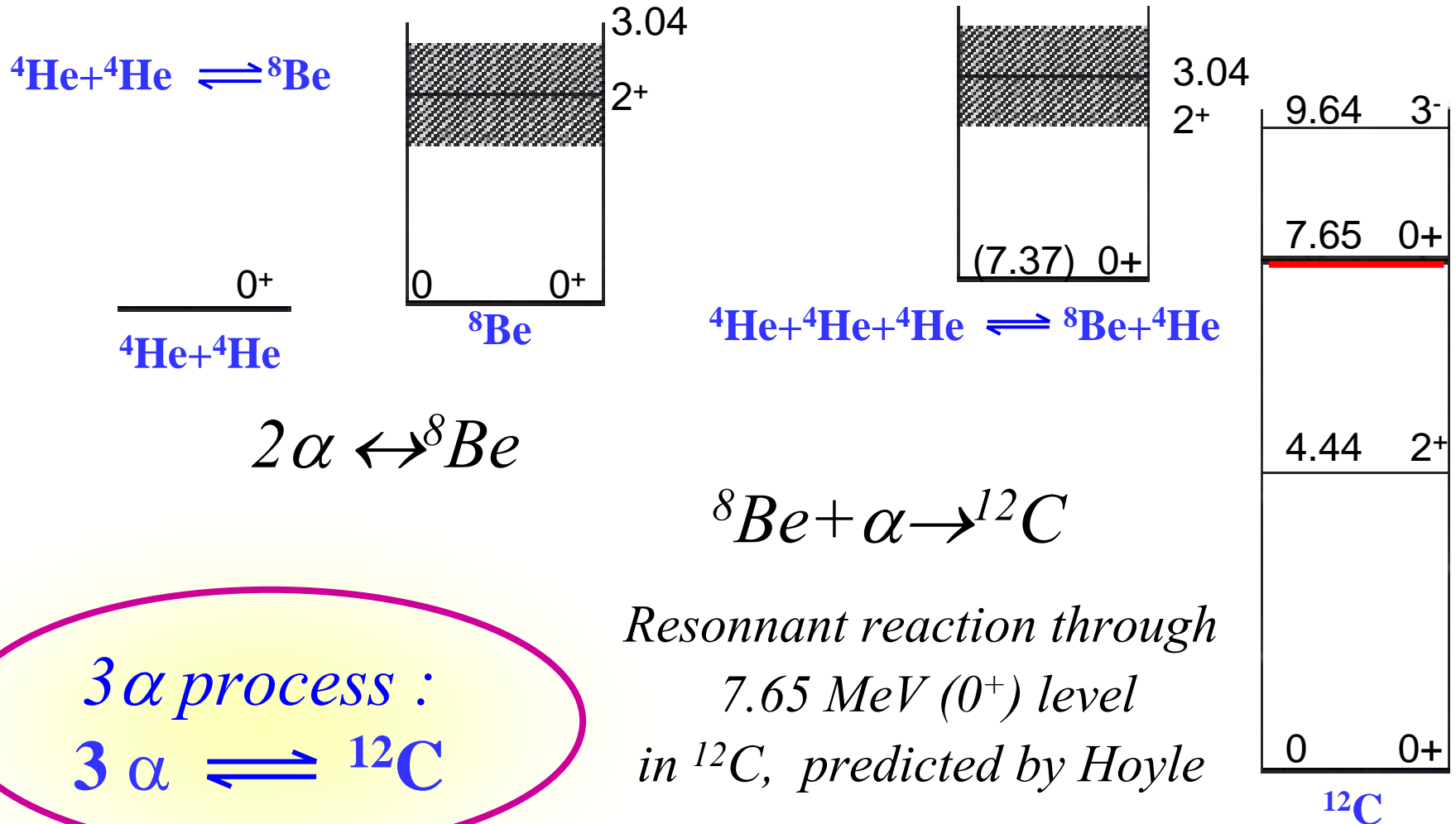
Fortunately: \exists a resonant level in ^{12}C

it allows α capture by ^8Be during the short life of ^8Be ($\tau \approx 2.6 \cdot 10^{-16}$ s):

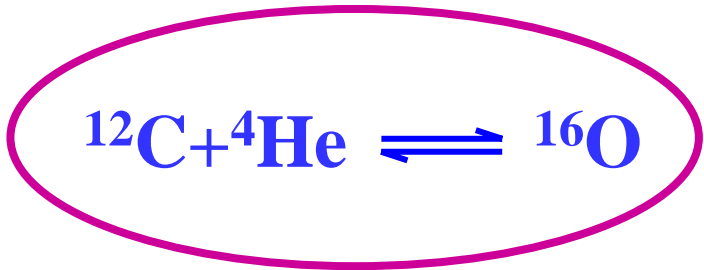
\rightarrow Triple- α process ($T_6 \approx 100$, $\rho \approx 10^5$ g/cm 3)

He burning

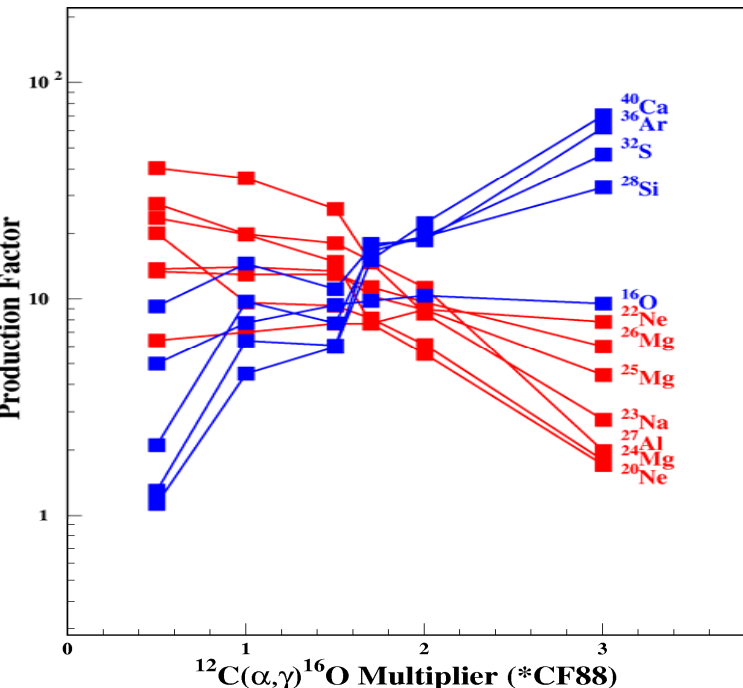
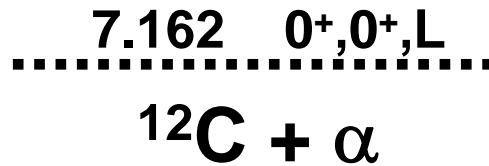
2-steps reaction via unstable ${}^8\text{Be}$ ($\tau \approx 10^{-16}$ s)



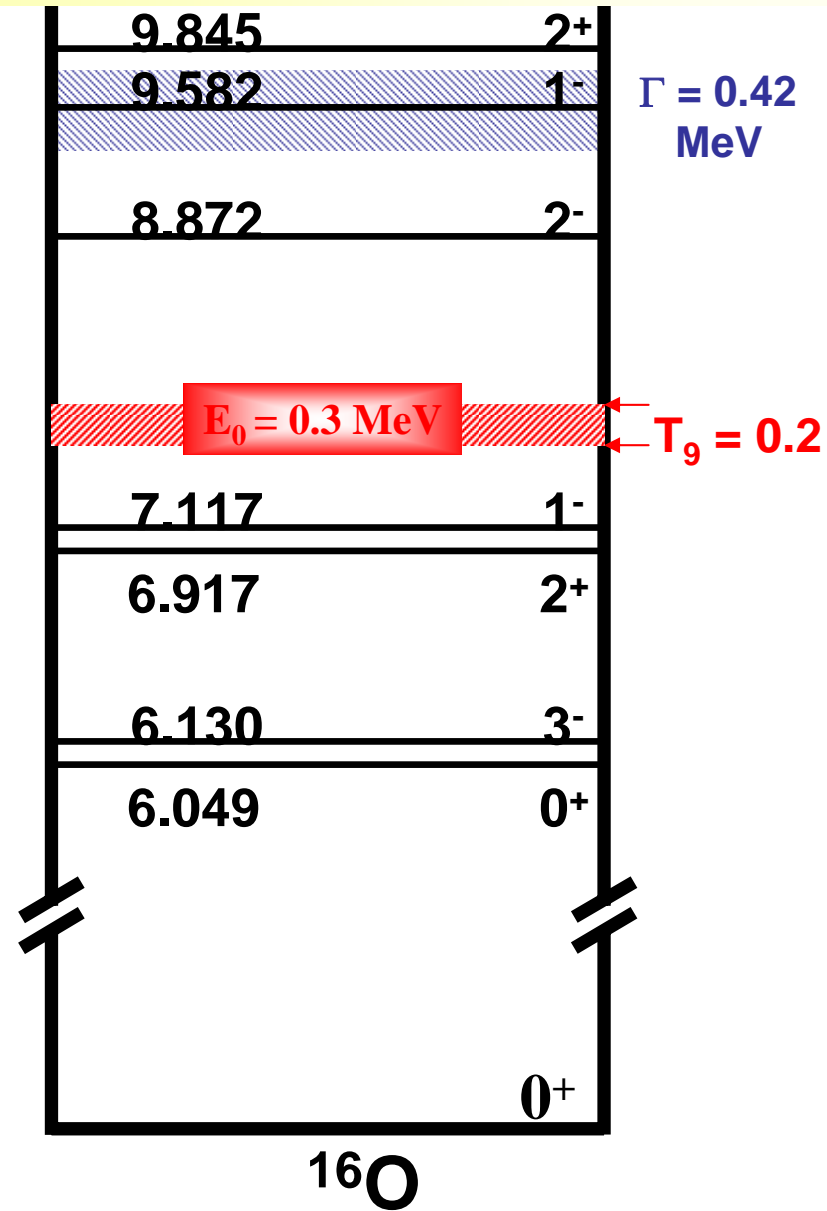
He burning



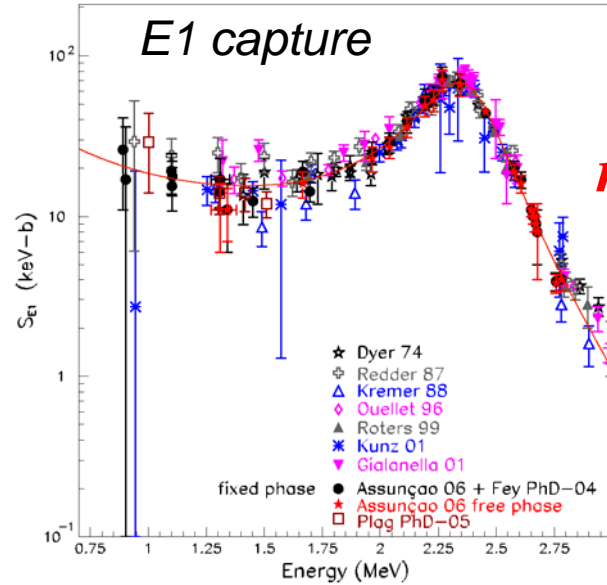
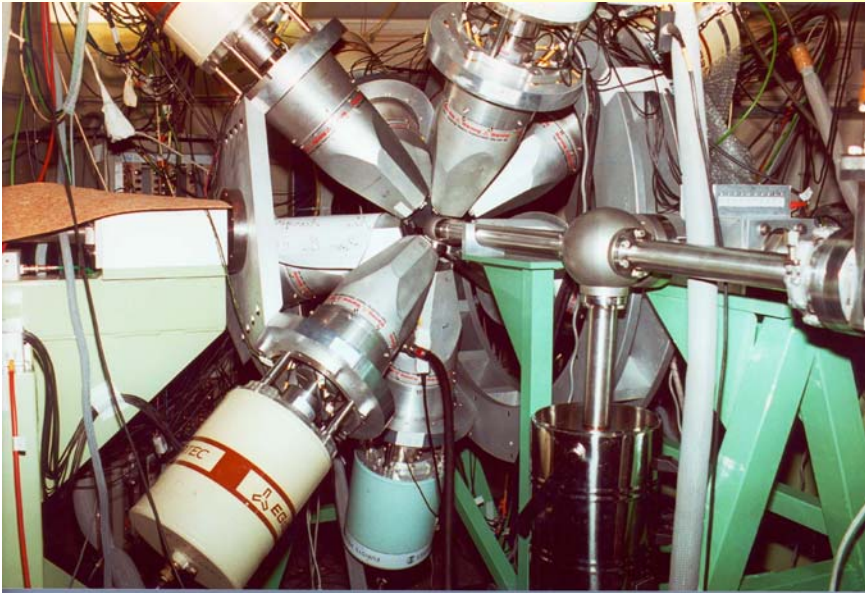
*Massive-stars
model
calculations*



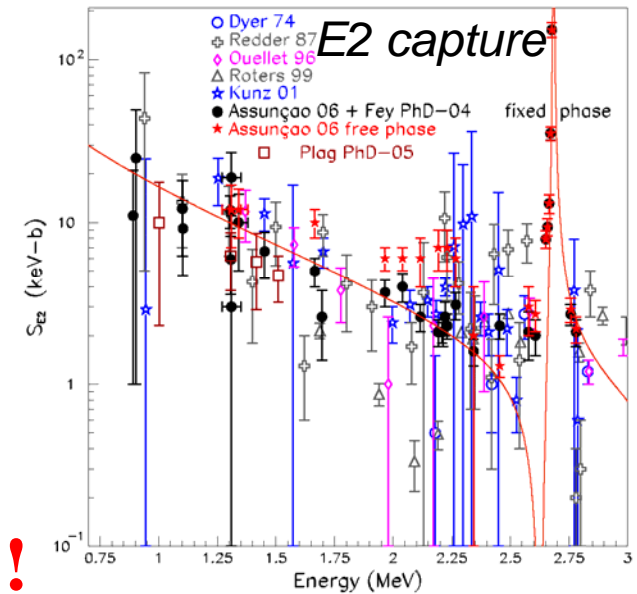
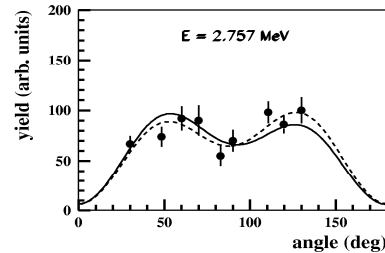
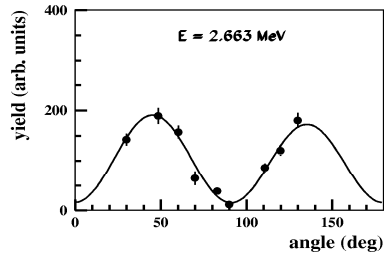
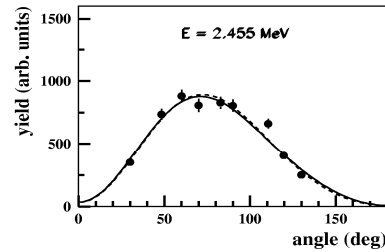
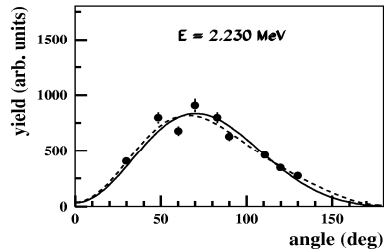
*For
12-13
15-20
25-30
35-40
M_⊙*



Experimental data for $^{12}\text{C}(\alpha,\gamma)^{16}\text{O}$



Low energy measurements down to $\sigma \sim 0.1 \text{ nb}$

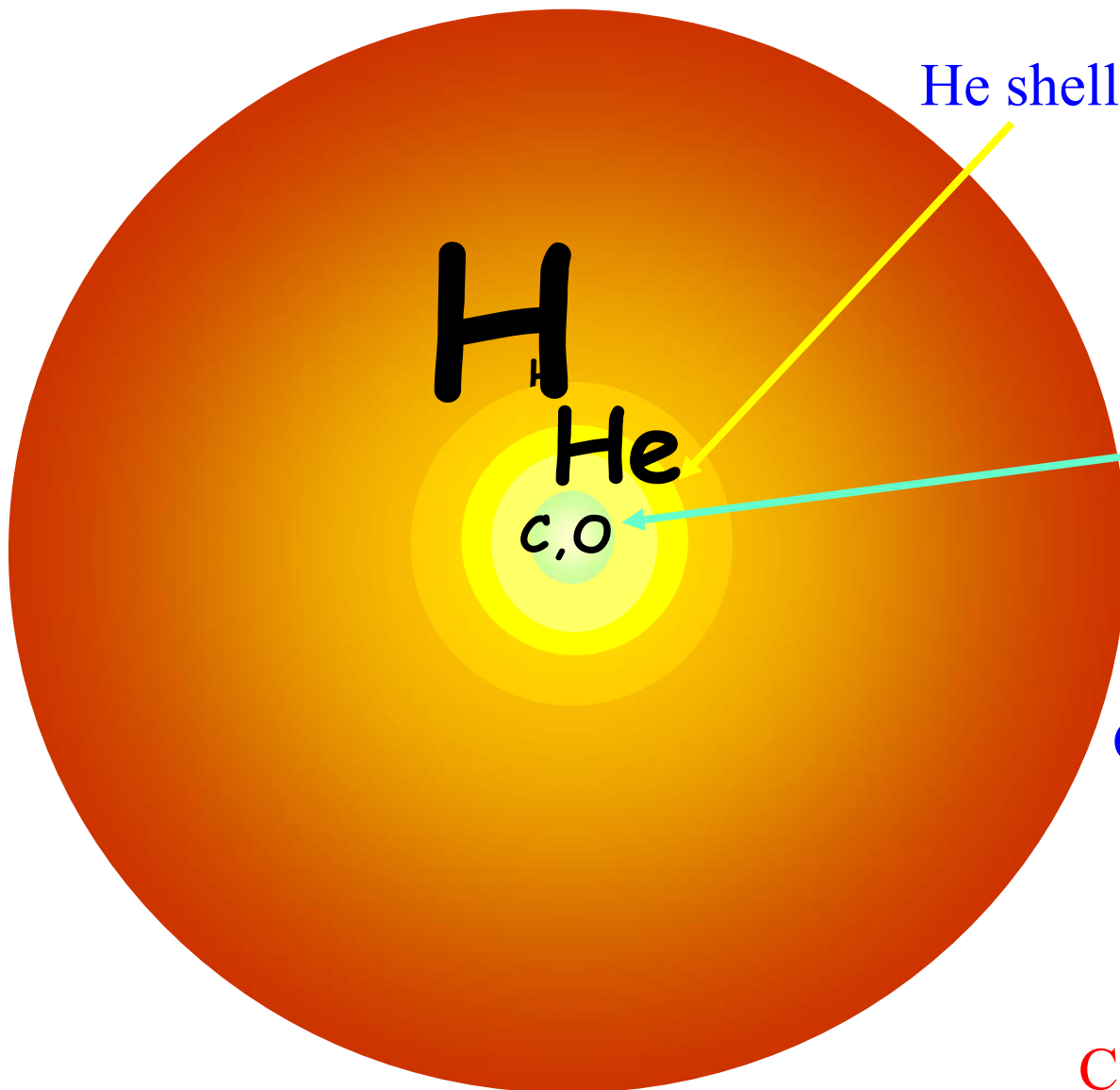


Extrapolations down to 300keV !

Projects: i.e. Triumf; Strasbourg-Orsay-Naples (down to 700keV ?)

He shell burning

(Asymptotic red giants : AGB)



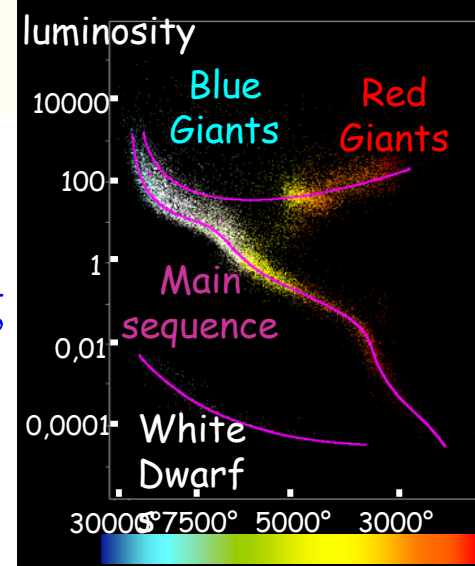
He shell burning

inert ^{12}C and ^{16}O core
(He is exhausted) :
gravitational contraction

Convective envelope (H)
radius ↗

H ashes burning → surface
(2nd « dredge-up »)

C,O observations



H and He shell burnings

(Thermal pulses « TP-AGB »)

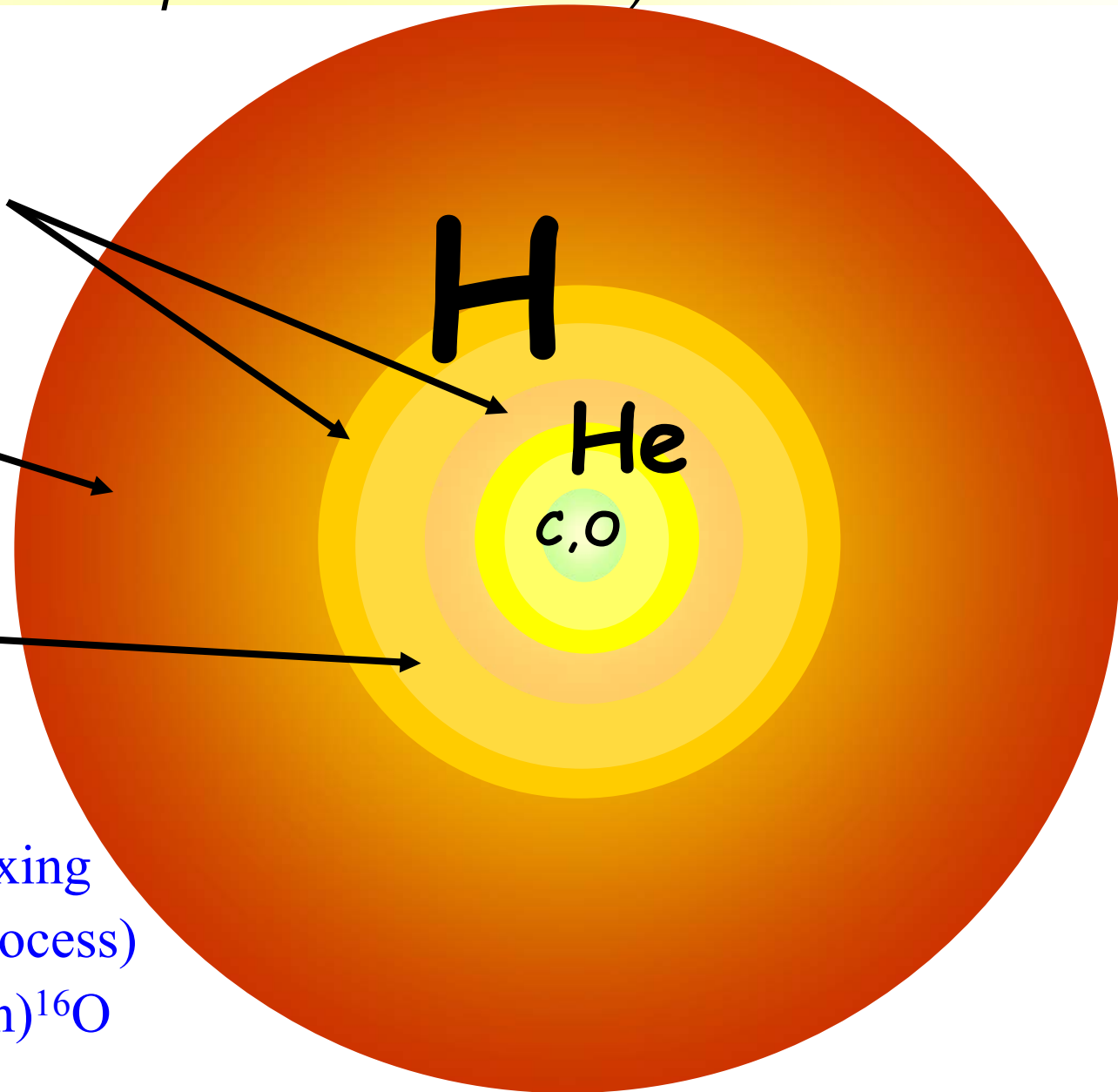
Unstable burning
in H and He shells

Convection in the
enveloppe :

(3rd « dredge-up »)

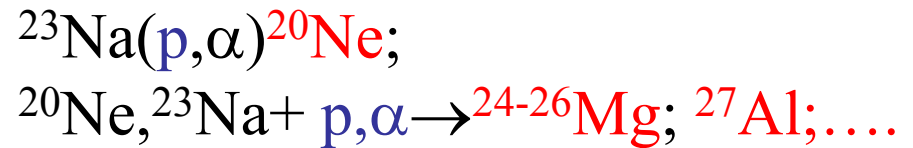
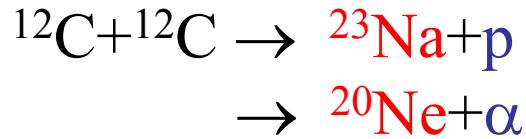
Convective zones
between
« thermal pulses »

H and He and ^{12}C mixing
→ n production (→ s process)

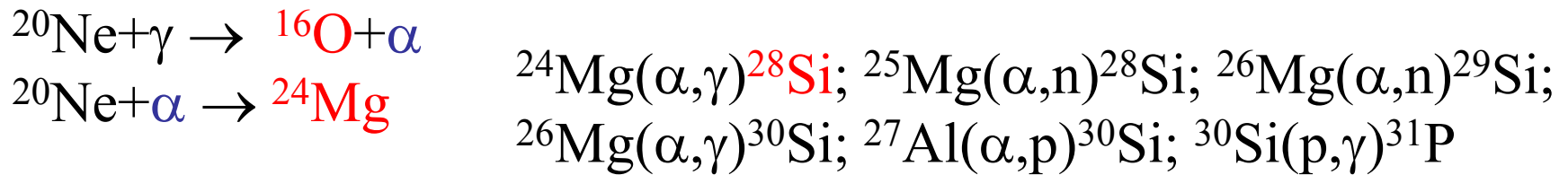


Advanced burnings : C, Ne, O

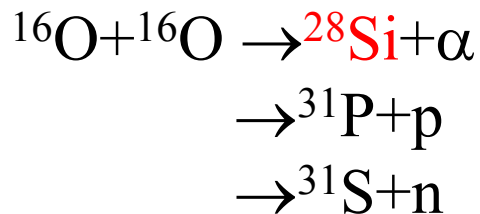
1. Carbon : $T \sim 10^9$ K



2. Neon: $T \sim 2 \times 10^9$ K



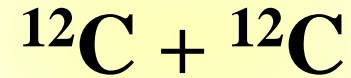
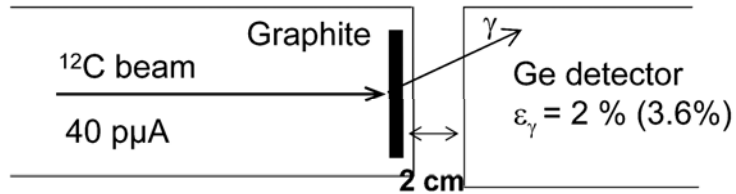
3. Oxygen: $T \sim 3 \times 10^9$ K



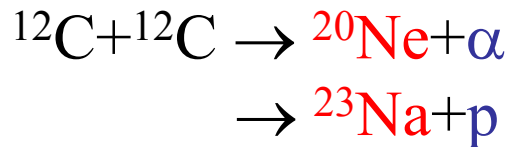
Super Red Giant star

∃ mass loss: function of the star's mass: red → blue

Advanced burnings : C, Ne, O



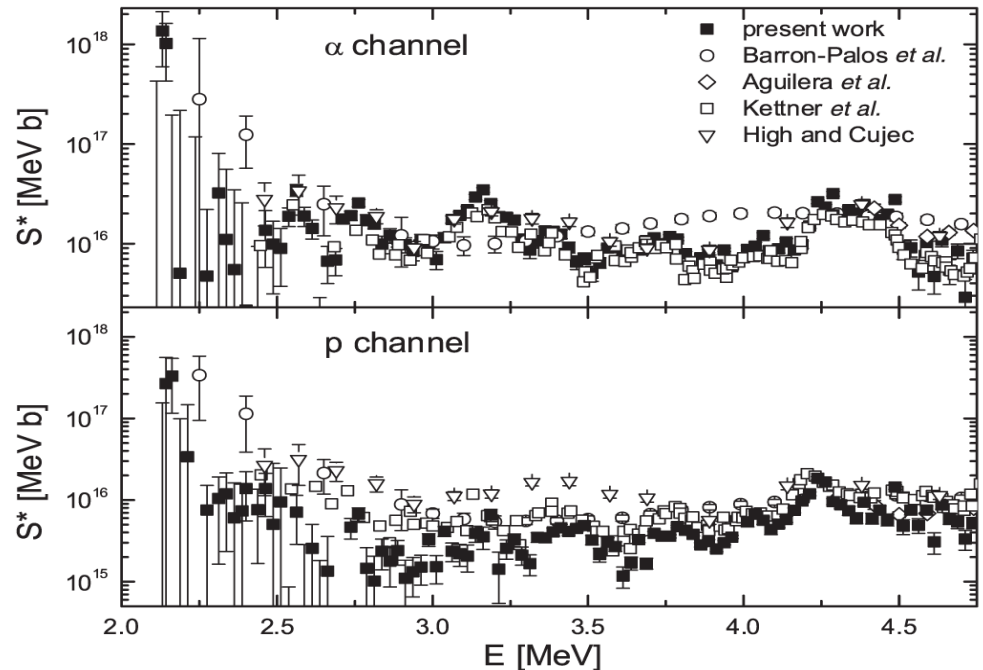
Spillane PRL 122501 (2007)



Resonances

at $E < 3.0$ MeV

**How to extrapolate
down to 1.5 MeV ?**



*Projects: i.e. Strasbourg-Triumf (nuclear structure);
with Naples (down to 1.5 MeV ?)*

Silicon burning

$\sim 4 \cdot 10^4$ K : ^{28}Si photodissociation : $^{28}\text{Si}(\gamma, p)^{27}\text{Al}$, $^{28}\text{Si}(\gamma, \alpha)^{24}\text{Mg}$,

→ fast photodissociation of less bound products : Al, Mg, Ne, O

free p, α , n : react with remaining nuclei.

Reaction rates \implies « *Statistical nuclear equilibrium* » :

$(\gamma, p) \rightleftharpoons (p, \gamma)$; $(\gamma, \alpha) \rightleftharpoons (\alpha, \gamma)$; $(\gamma, n) \rightleftharpoons (n, \gamma)$; $(p, n) \rightleftharpoons (n, p)$, ..

\implies most bound nuclei formation: ^{56}Ni (\rightarrow ^{56}Fe)

^{56}Fe is the most bound nucleus:

→ any reaction between charged particles will dissipate energy
(instead of create in earlier stages)

→ Fe core will increase up to Chandrasekhar mass: instability

→ Shock wave between core instability and free falling layers

→ Supernova explosion

Core \rightarrow n star or black hole

Tidy « onion » star : pre supernova

Burning stages:

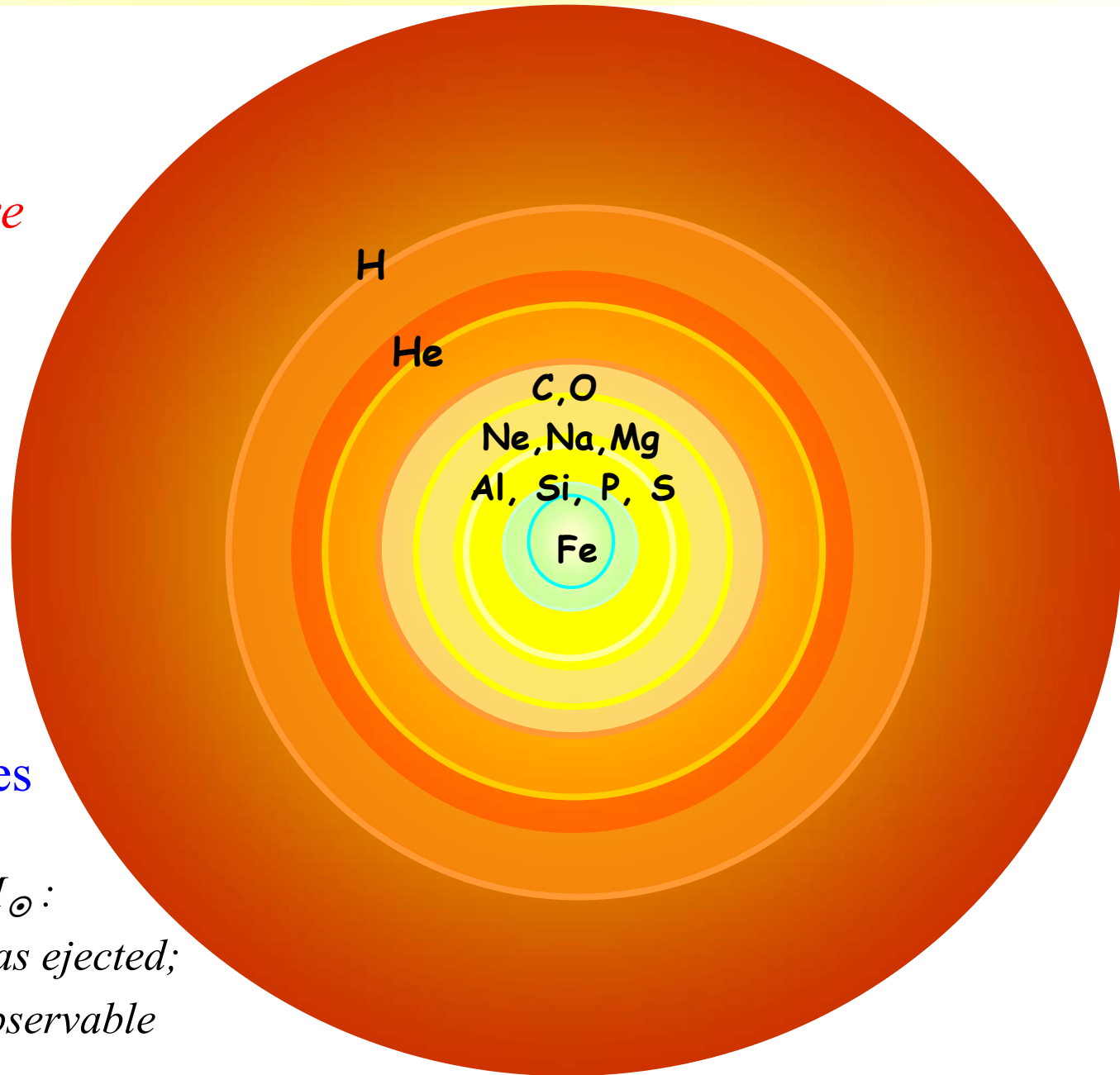
core + shells

« Onion » structure

Only the most massive stars ($M_{\star} \geq 10 M_{\odot}$) will live all these stages

Heaviest stars: important mass losses

*(For $M_{\star} > 40 M_{\odot}$:
the whole H envelope was ejected;
He envelope may be observable*



Wolf-Rayet star (« WR »)

*Heaviest stars
(Wolf-Rayet)*

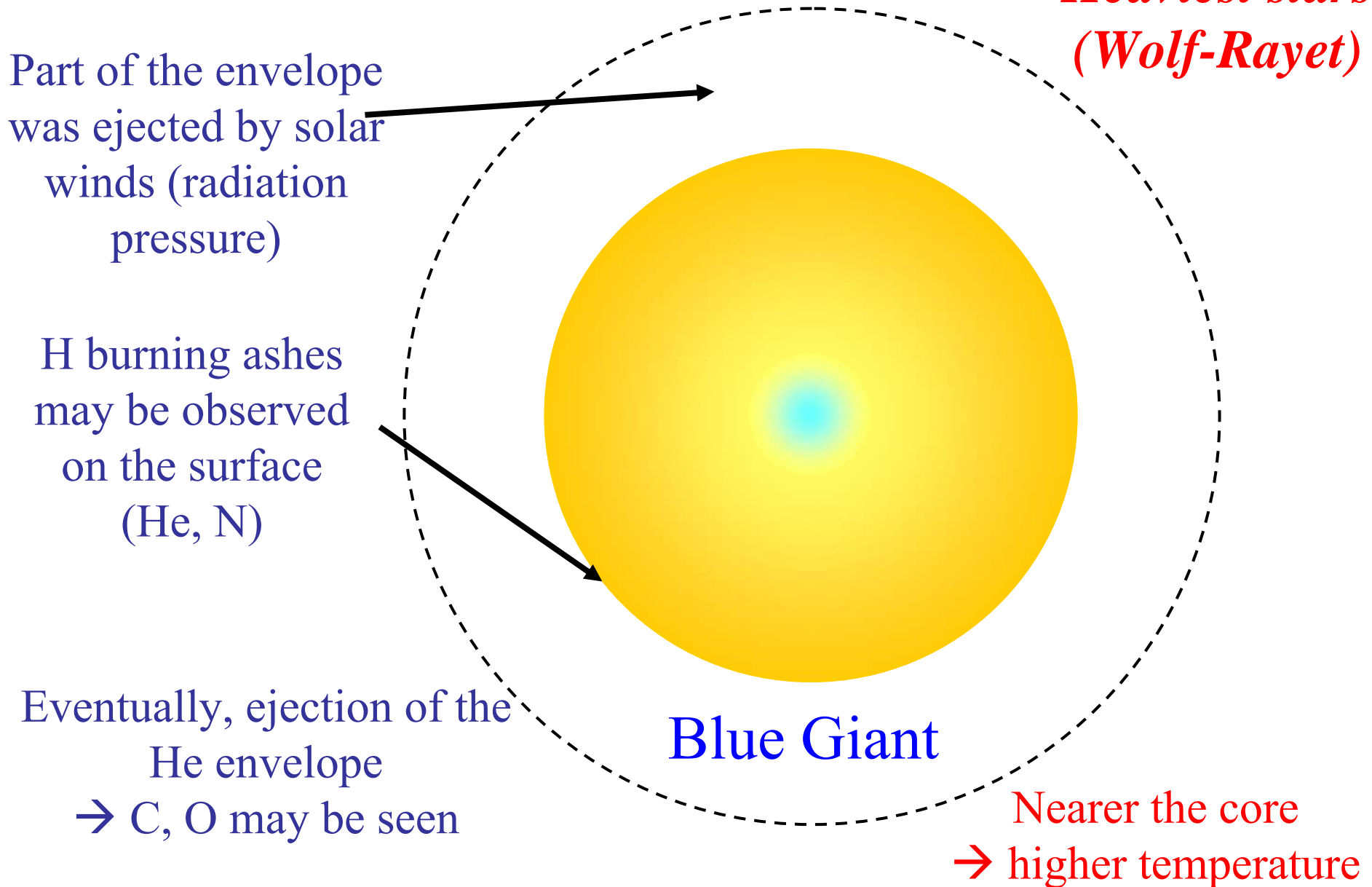
Part of the envelope
was ejected by solar
winds (radiation
pressure)

H burning ashes
may be observed
on the surface
(He, N)

Eventually, ejection of the
He envelope
→ C, O may be seen

Blue Giant

Nearer the core
→ higher temperature



White dwarfs

- Final evolution stage of a star with a mass lighter than 6-9 M_{\odot}
- After He or C burnings for the heaviest
- Stabilized by the degenerate electron pressure
- $M_{\text{WD}} < 1.35 M_{\odot}$ (Chandrasekhar mass)
- Radii of a few thousands of km

From the star formation

Explosion

$\sim 1/4 \text{ He}$
 $+ 3/4 \text{ H}$

Burning sequences C/O

Si

Gravitational contraction

He

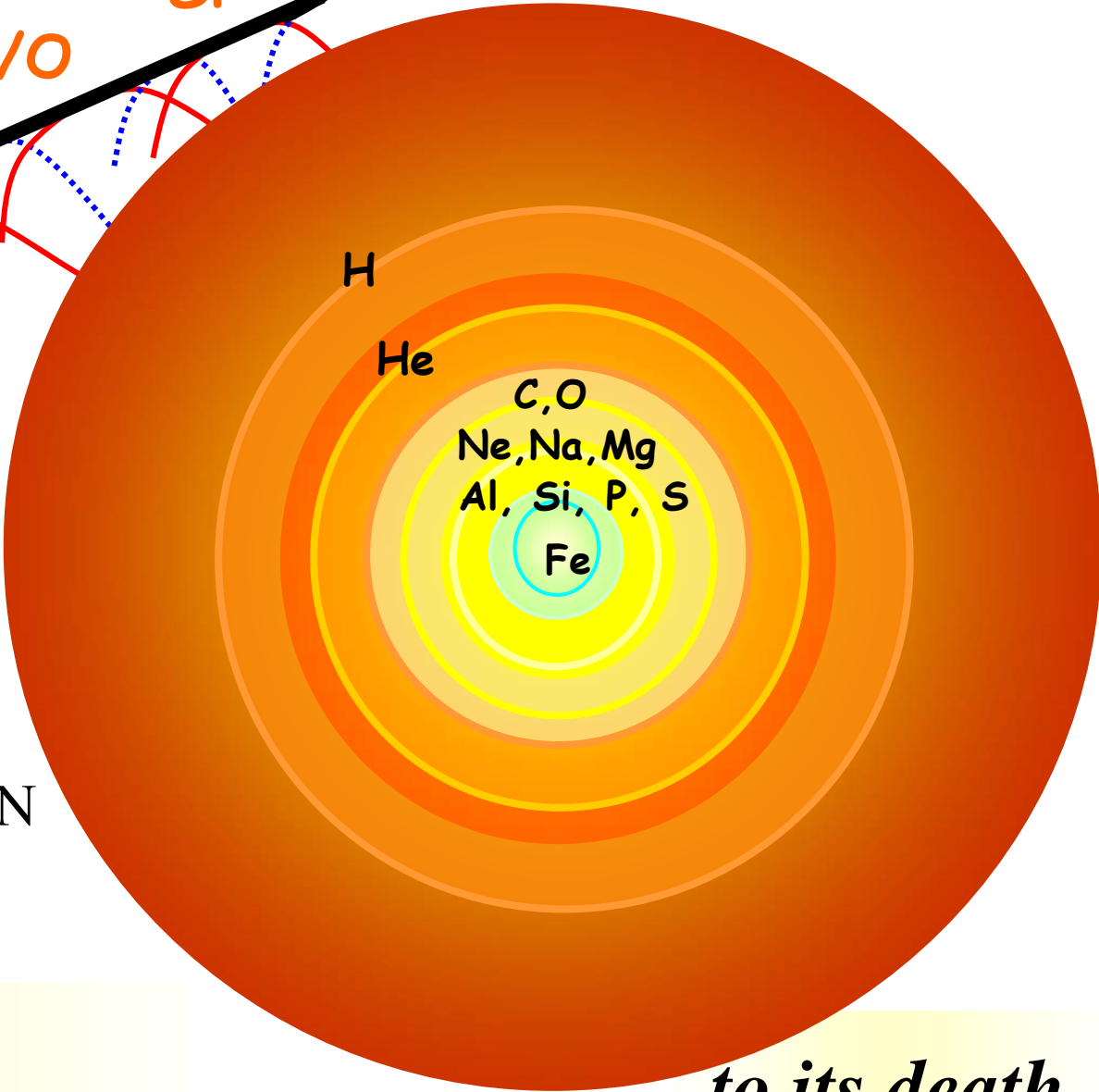
H

white dwarfs

Function of its mass:

Quite death \rightarrow white dwarfs

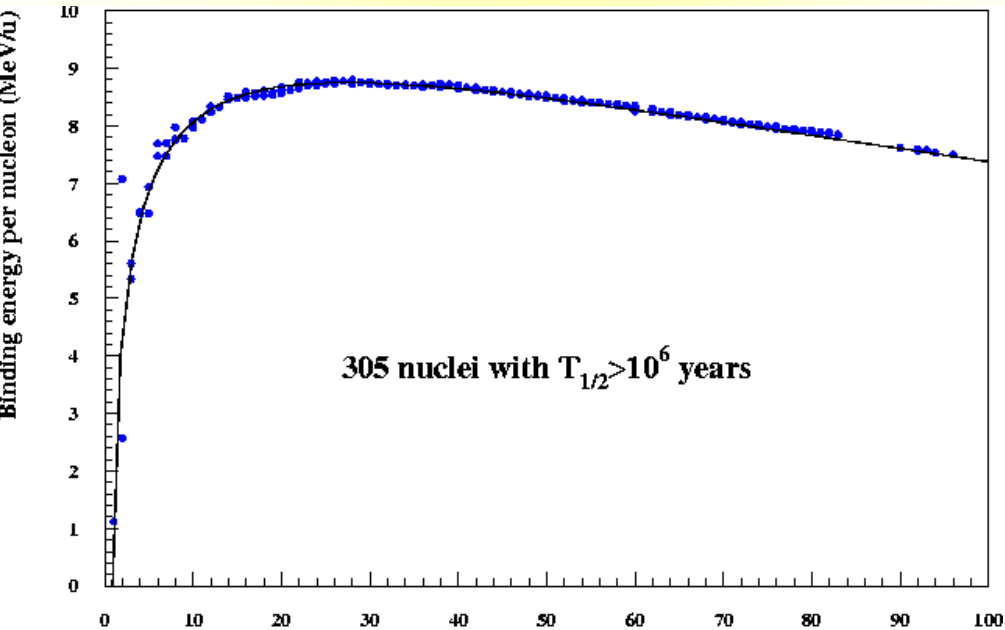
Explosive events: Novae, SN
 \rightarrow Neutron star
 \rightarrow Black hole



Star life

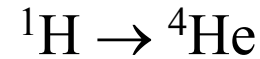
to its death

Nuclear/chemical species synthesis

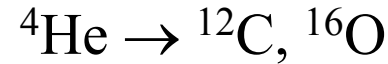


Interstellar gas enrichment:
 stellar winds
 explosions
 → 2nd generations stars
 are metal enriched

Hydrogen burning



Helium burning



Carbon burning



Neon burning



Oxygen burning

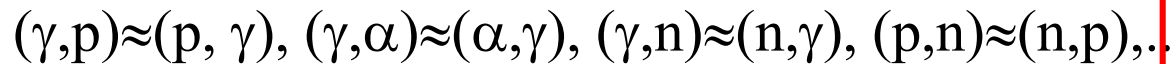


Silicon burning

Photodisintegration



Statistical nuclear equilibrium



→ Most bound nuclei

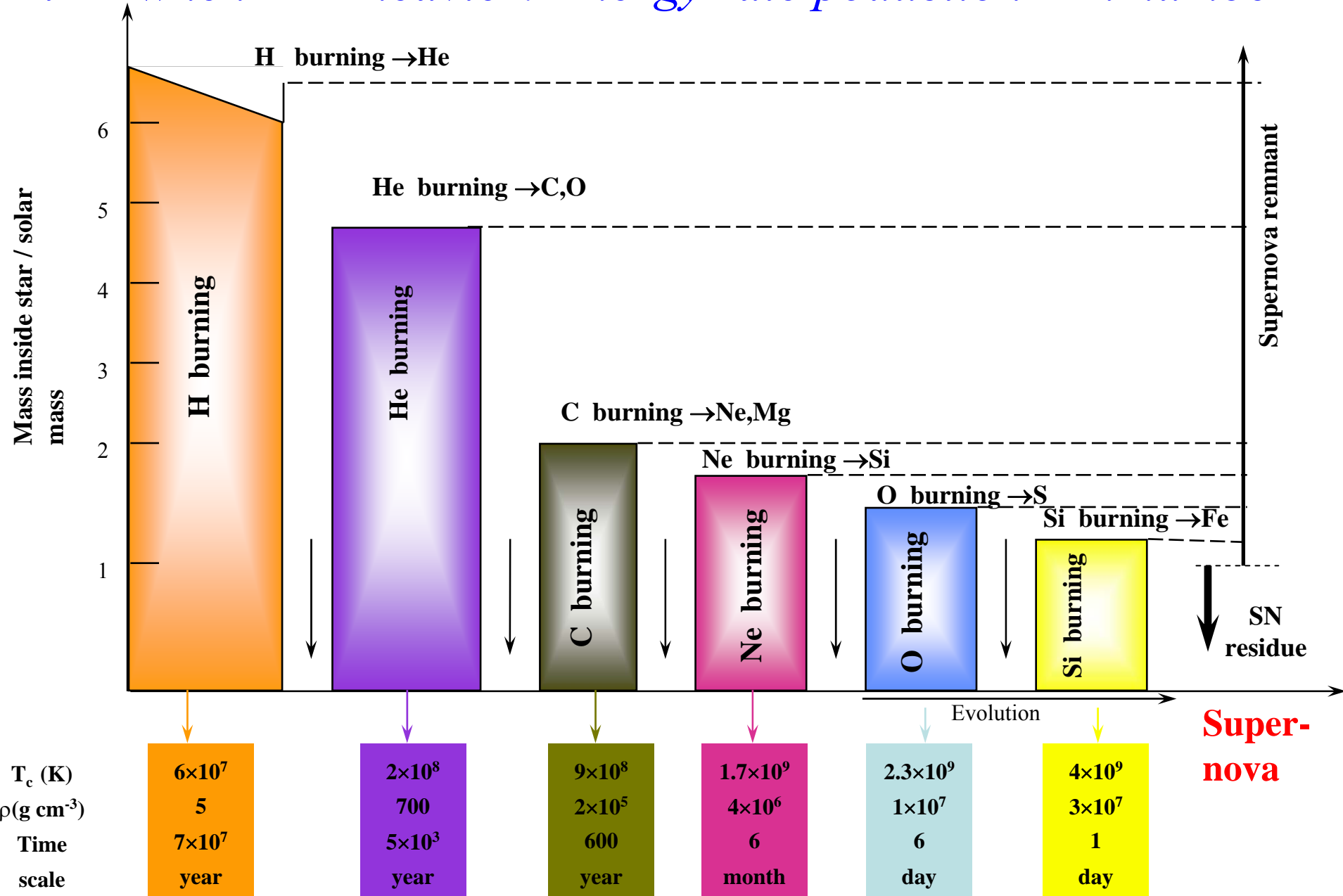


MASS, TIME



Internal structure evolution for a $M = 25M_{\odot}$ star

$\Delta t \searrow$ when $H \rightarrow$ heavier: Energy rate production \searrow n number \nearrow



Crab Nebula: in optic wave lengths

SN 1054 remnant



FORS team, 8.2-meter VLT, ESO

Some star-death examples

☆ *low-mass stars* ($0.08 M_{\odot} \lesssim M \lesssim 0.45 M_{\odot}$):

H core burning then *white dwarfs*

☆ *low-mass stars* ($0.45 M_{\odot} \lesssim M \lesssim 8 M_{\odot}$):

H and He core burnings then *white dwarfs*

...

☆ *massives stars* ($M \geq 10 M_{\odot}$) :

Any stages then *SN II* (H envelope), *SN Ib* (H envelope loss)
or *SN Ic* (H and He envelope loss)

∃ *Numerous binary systems: « binary stars »*

Evolution of one star may influence the other star evolution

Ex. : white dwarf + companion (as a red giant star) → *Nova, SNIa*

neutron star + companion (as a red giant star) → *γ-ray burst*

...

Physical conditions in stars

Core temperatures

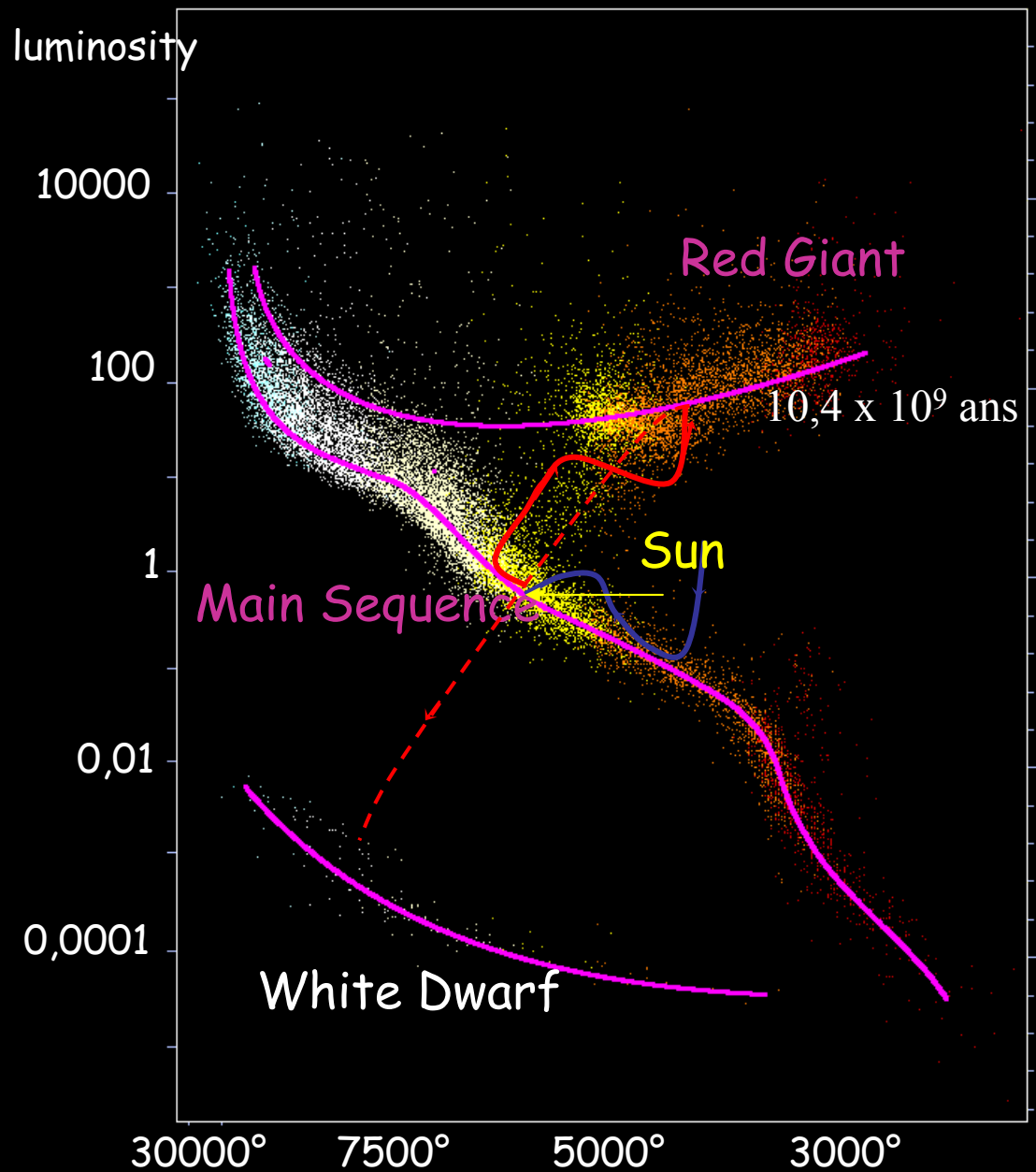
Sun :	$15 T_6$	$E(k_B T) = 1.3 \text{ keV}$	(H burning)
Red Giant :	$\sim 0.2 T_9$	$E(k_B T) = 17 \text{ keV}$	(He burning)
Pre Supernova :	$\sim 4 T_9$	$E(k_B T) = 344 \text{ keV}$	(Si burning)
Nova :	$\sim 0.3 T_9$	$E(k_B T) = 26 \text{ keV}$	(Si burning)
Primordial : Universe (\sim min)	1 to 0.1 T_9	$E(k_B T) \sim 100\text{-}10 \text{ keV}$	(“Big Bang”) (nucleosynthesis)

Core density: 150 g/cm^3 for the Sun

10^{10} g/cm^3 for a white dwarf

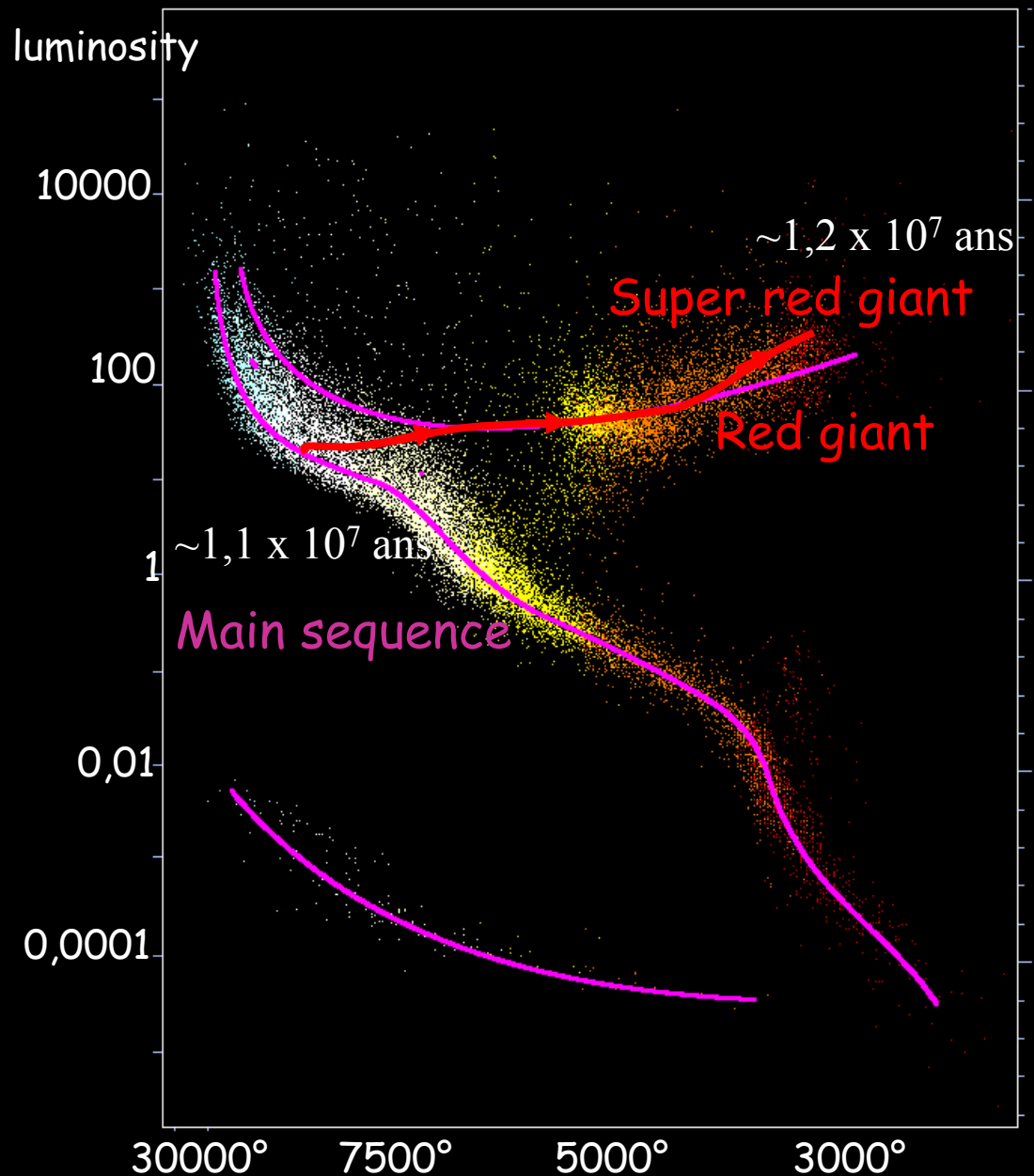
Temperatures: $a T_6 = a \times 10^6 \text{ K}$; $a T_9 = a \times 10^9 \text{ K}$

Life of the Sun

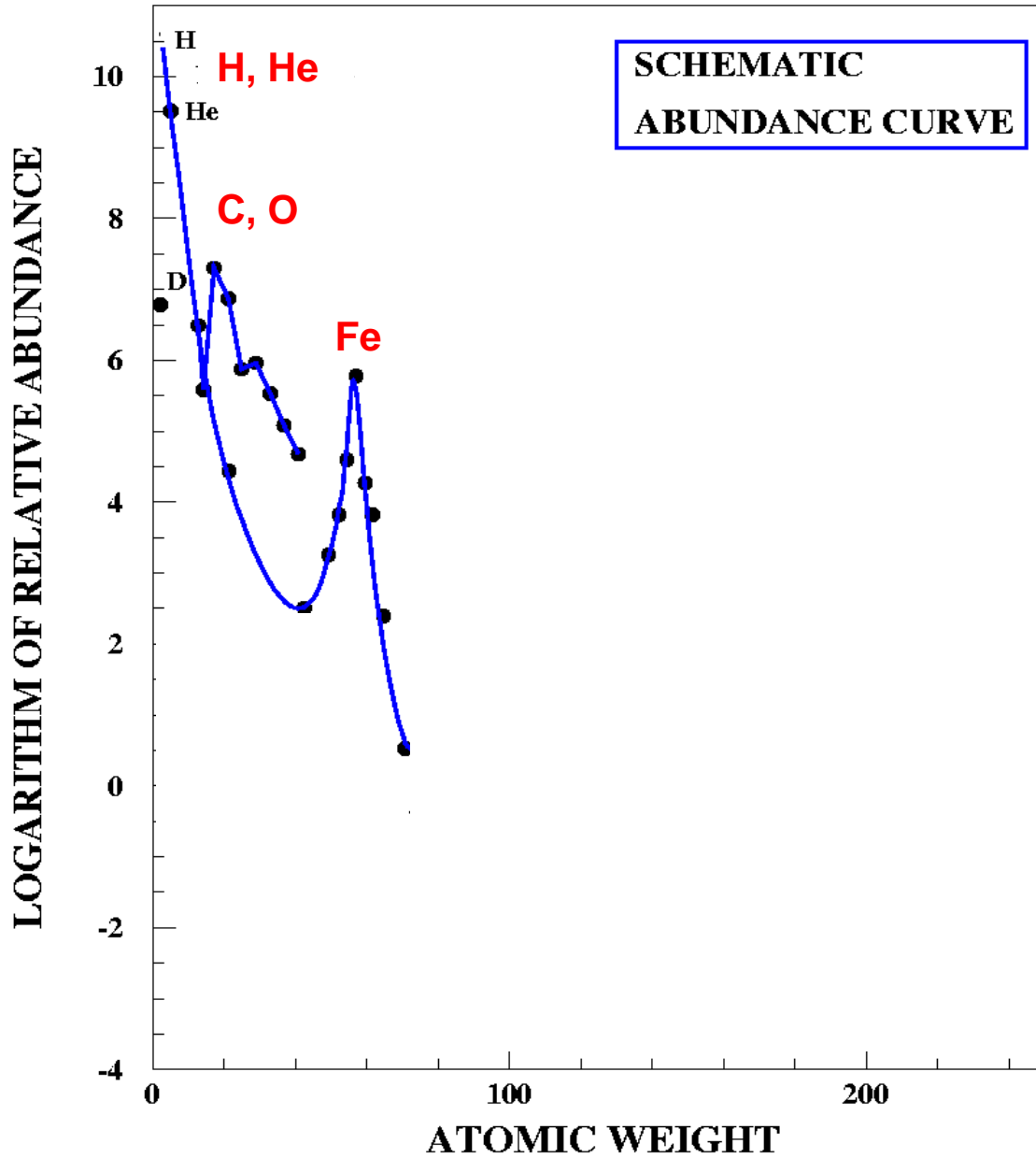


Life of a $15 M_{\odot}$ star

Will die in
Supernova



Abundance understandings



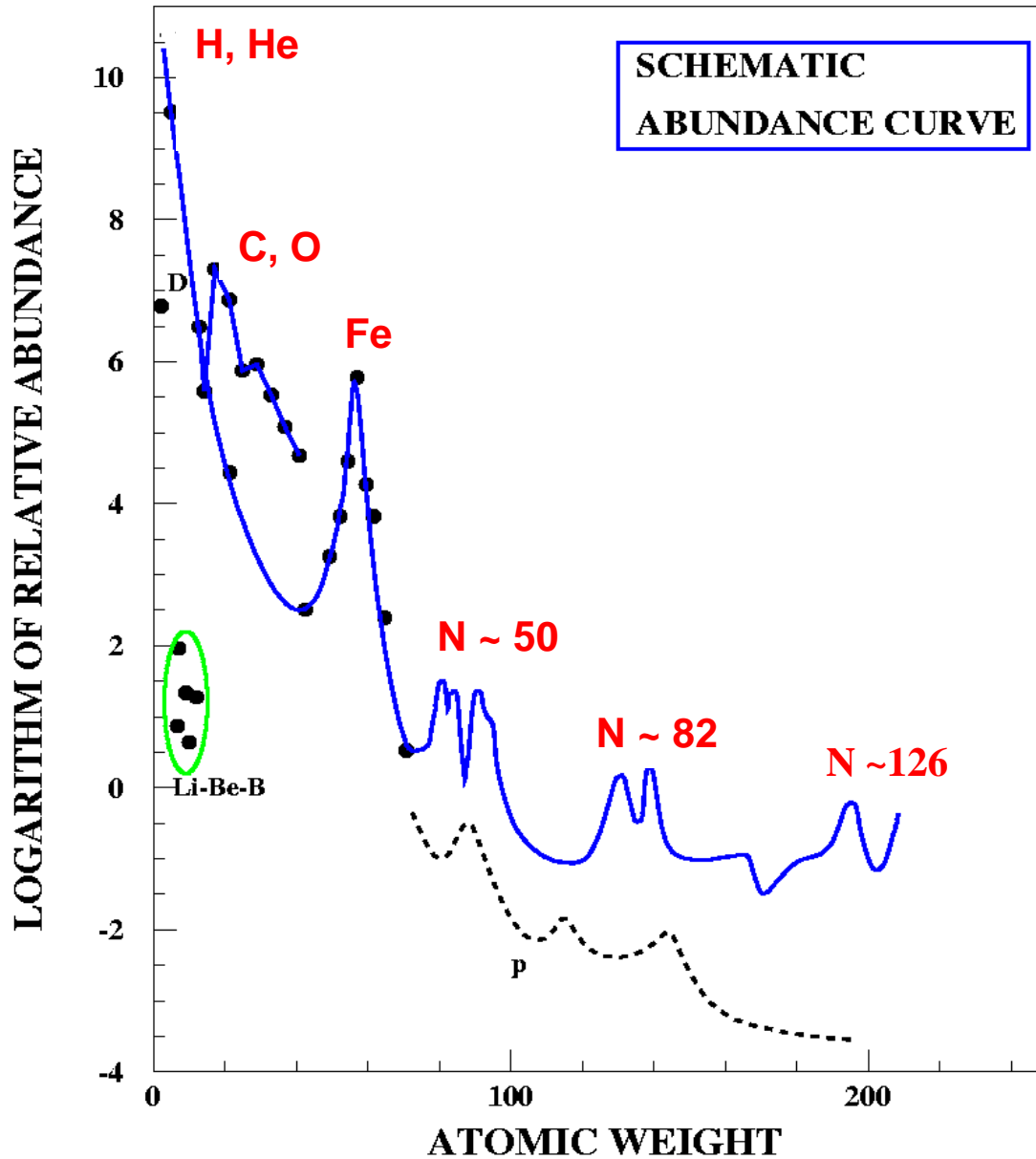
Abundant light elements ;
High abundances
for element species
with $N, Z \sim 2, 8, 28$

→ burning stages
→ magic-numbers
correlation

Primordial
nucleosynthesis:
→ H and He
highest abundances

Stellar nucleosynthesis:
→ C, O, Si, Fe peaks

And the others?



Other processes are
needed

^{56}Fe : most bound nucleus:

Only reactions with
neutrons
may explain
heavier nuclei abundances
no Coulomb barrier

Where?

When?

And LiBeB?

How?

Where?

When?

Production of elements beyond Fe

Charged particle burnings : impossible (Coulomb barrier and binding energy) \Rightarrow neutron capture reactions and photodissociation

- s process

- Slow neutron capture (compared to τ_β)
- During He burning in AGB and massive star cores
- Neutron sources: $^{13}\text{C}(\alpha, n)^{16}\text{O}$ and $^{22}\text{Ne}(\alpha, n)^{25}\text{Mg}$

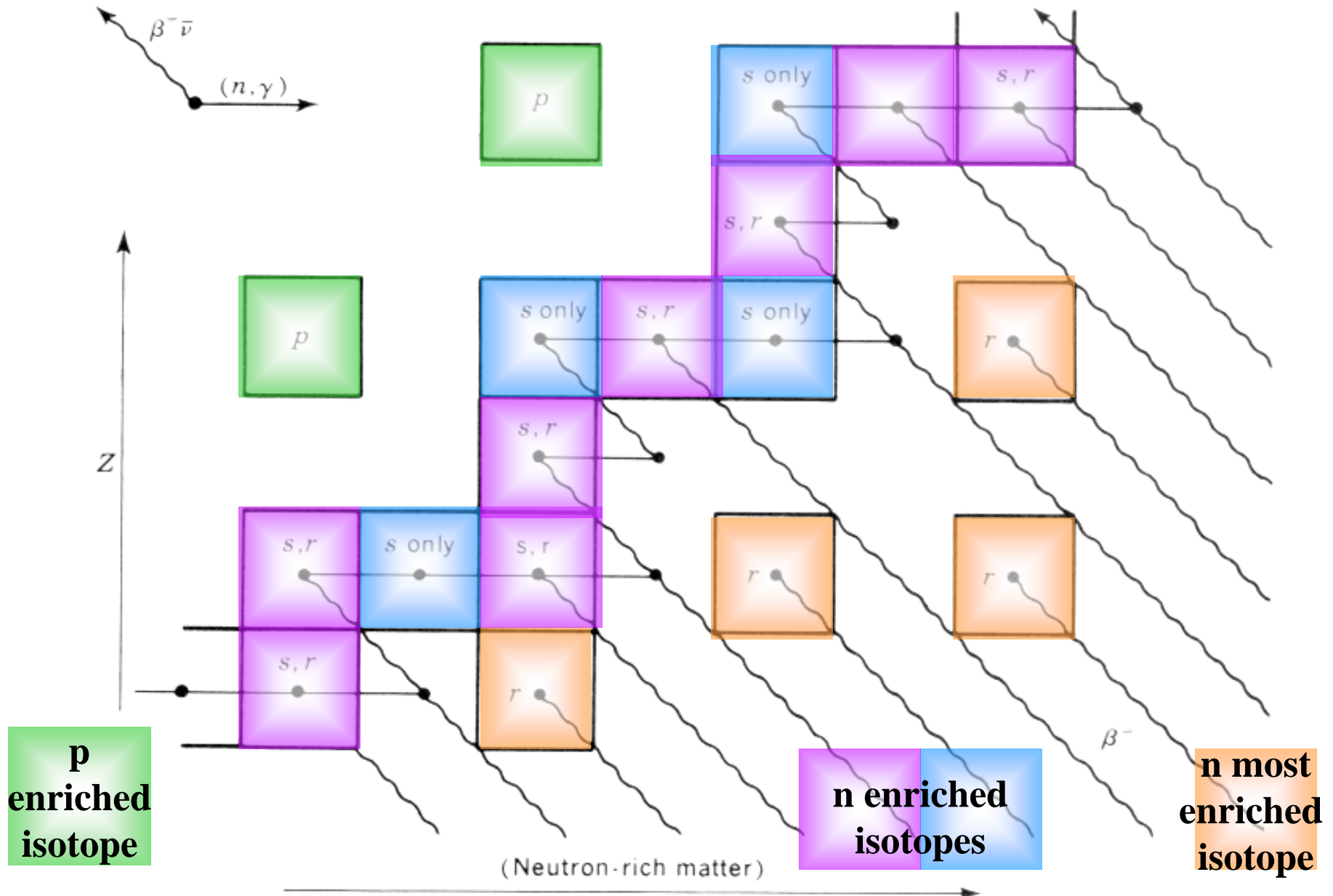
- r process

- Rapid neutron capture (compared to τ_β)
- In deep layers during SNII (?)

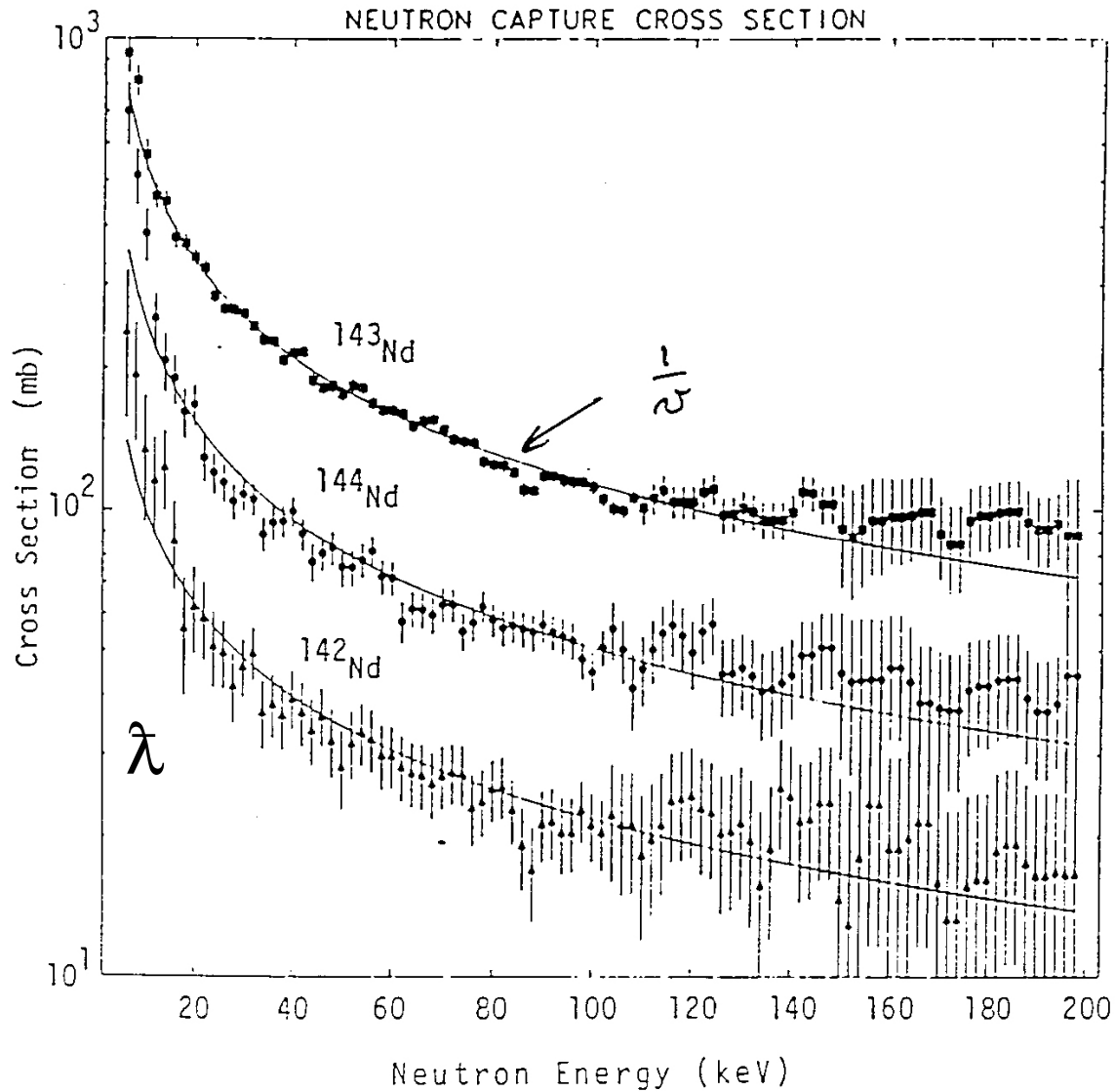
- p process

- Photodissociation of « s » and « r » nuclei
- Massive stars and (pre-)supernovae

s, r and p processes



Neutron captures



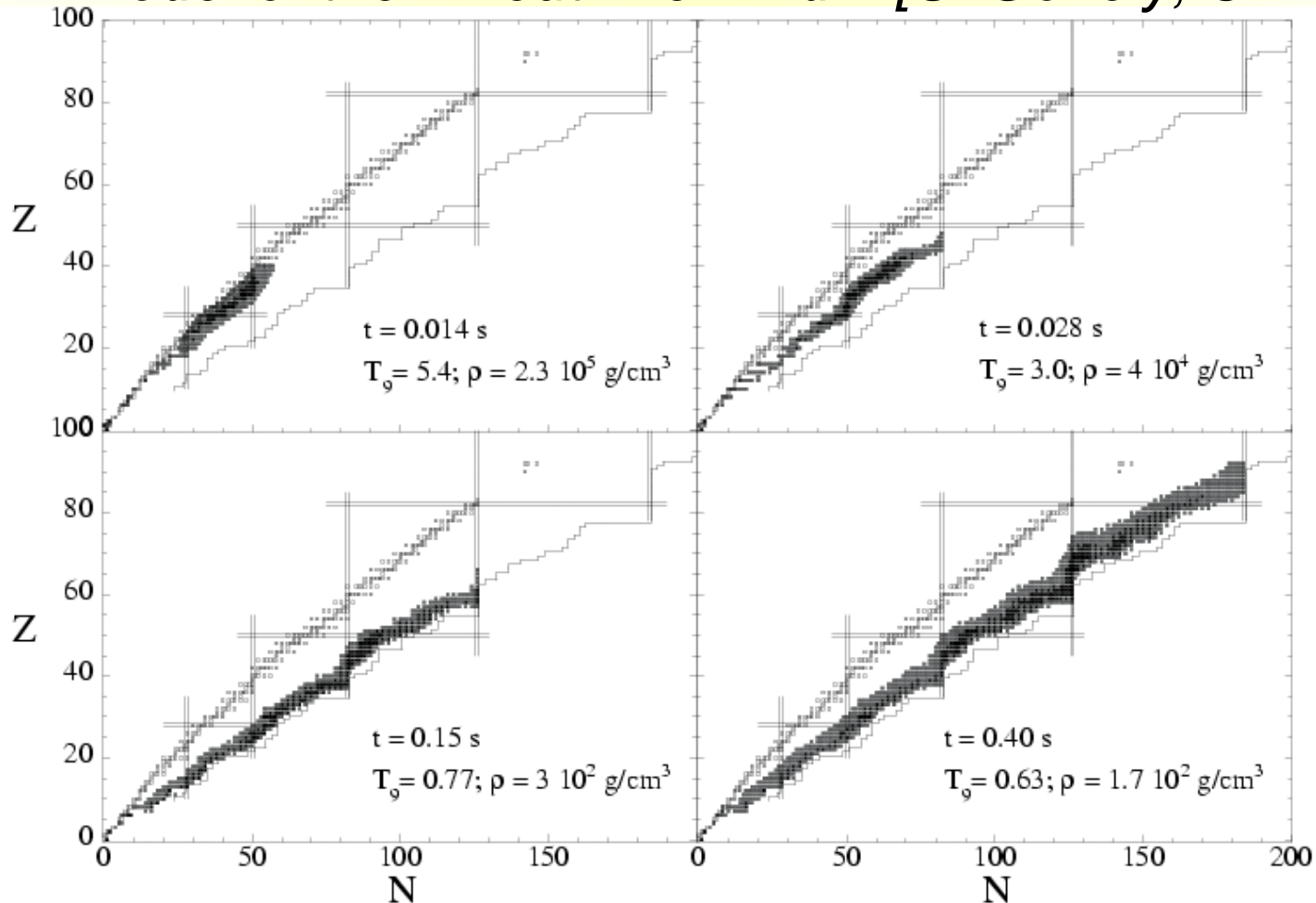
$$\sigma \propto \hat{\lambda} T \propto 1/v$$

In this case :
variation in $E^{1/2}$

$$\langle \sigma v \rangle \approx \text{cste}$$

A dynamical r process calculation

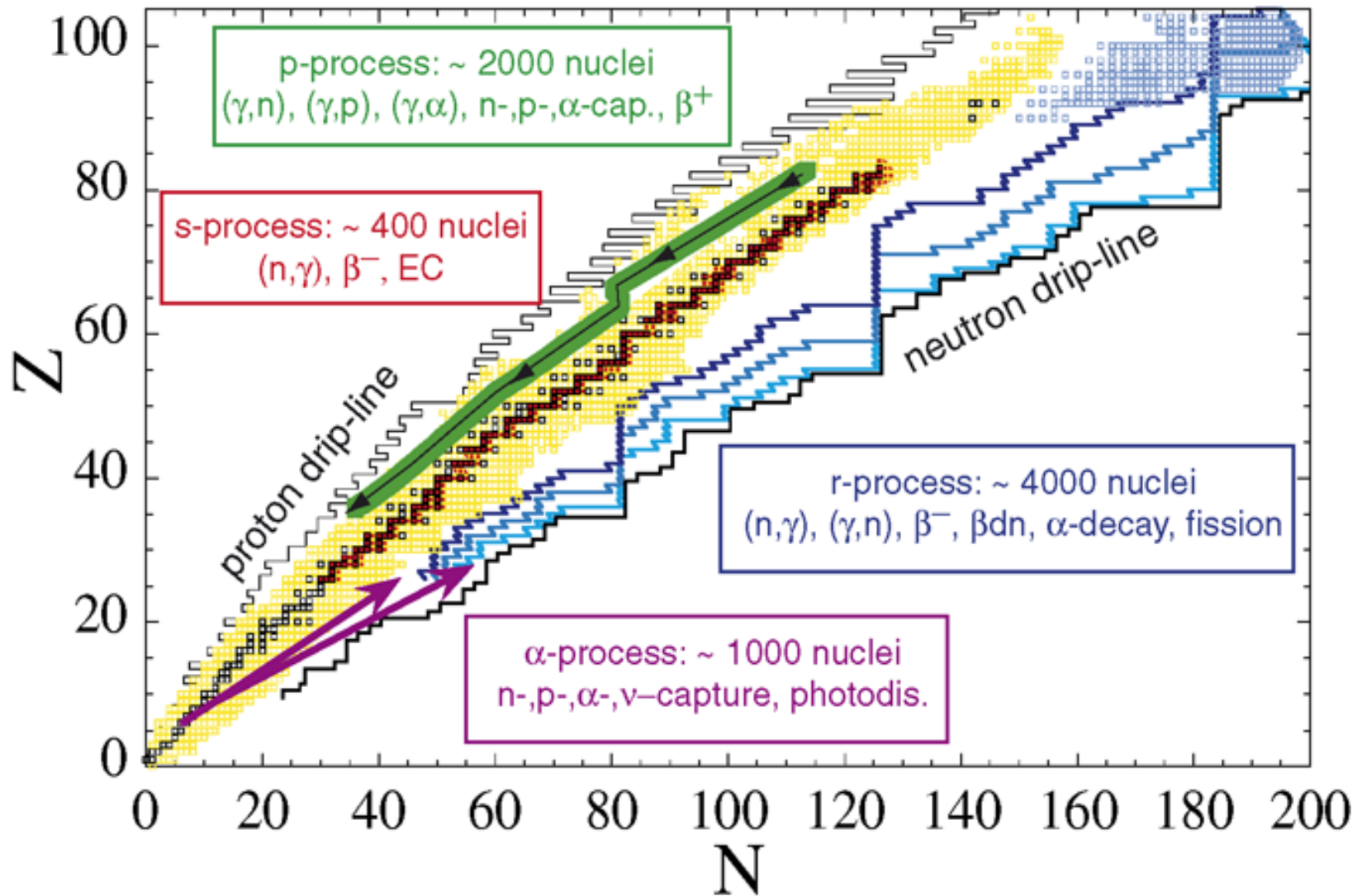
model of the « neutrino wind » [S. Goriely, ULB]



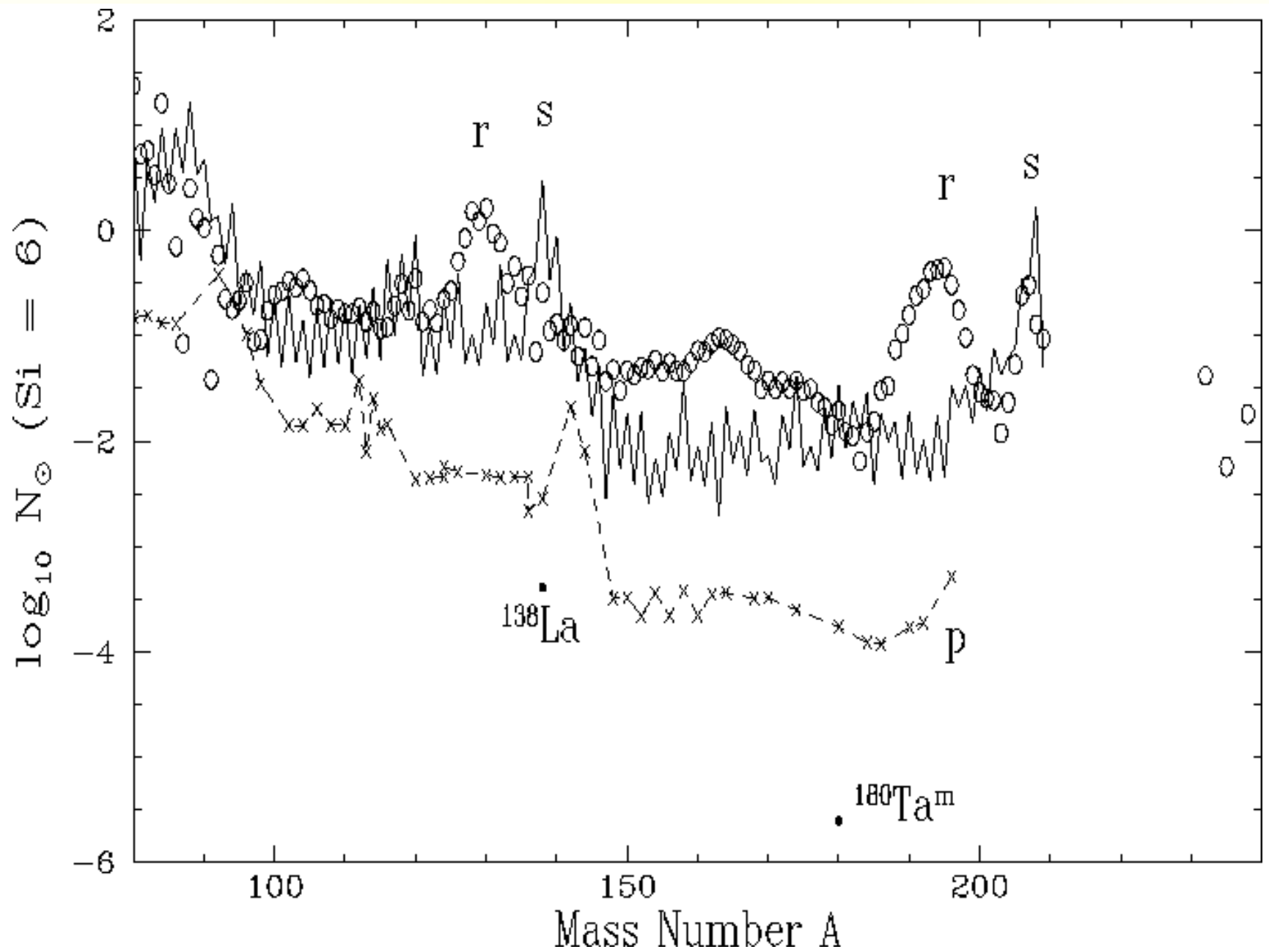
∃ various ways of the r process. Thousands of nuclei, capture reactions, periods, fission and neutrino induced reactions, ...

- ✓ Theory : phenomenological models vs microscopic calculations
- ✓ Chosen measurements (**Rikken, GSI, SPIRAL2, EURISOL**)

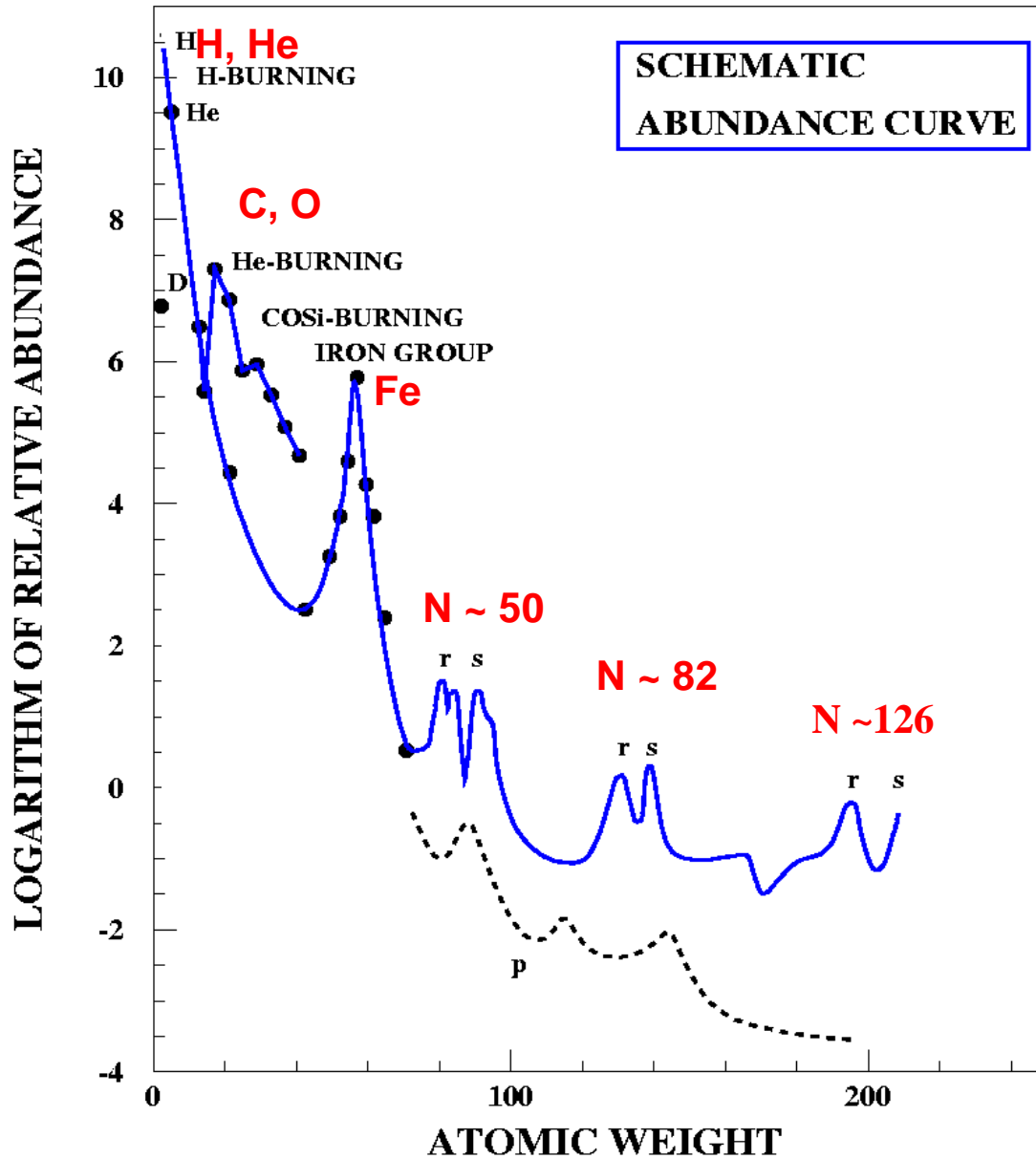
s, r, p processes



s and r nuclei



Abundance interpretations?



Other processes are needed

^{56}Fe : most bound nucleus:

Only reactions with neutrons

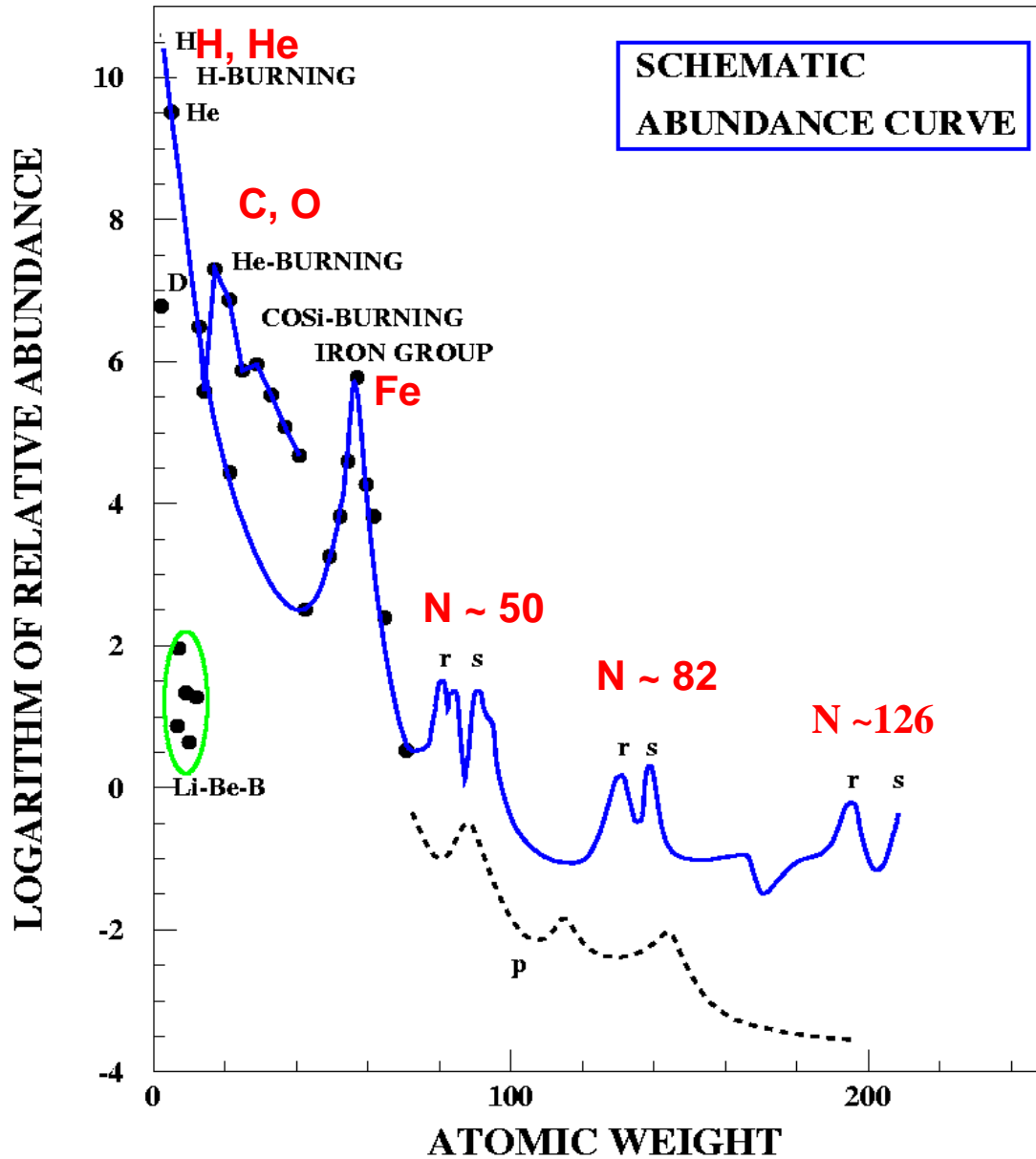
may explain

heavier nuclei abundances

no Coulomb barrier

→ double peaks:
r and s processes

And LiBeB ?



Other processes are
needed

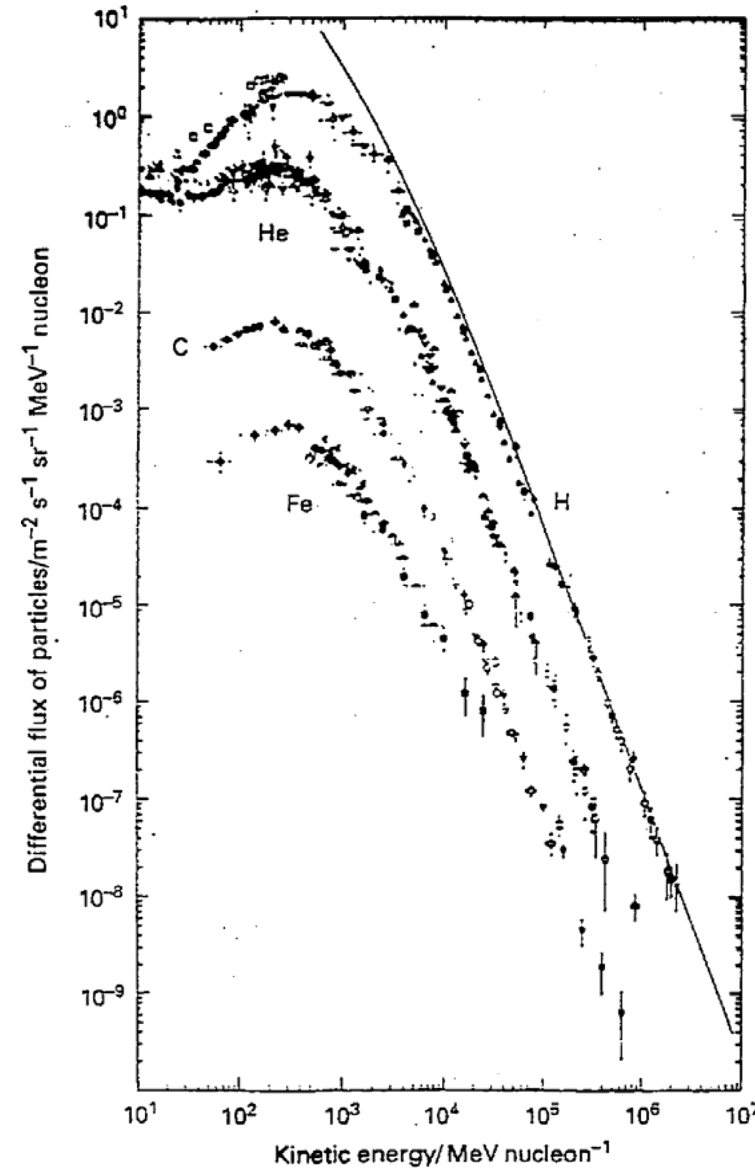
And LiBeB?

How?
Where?
When?

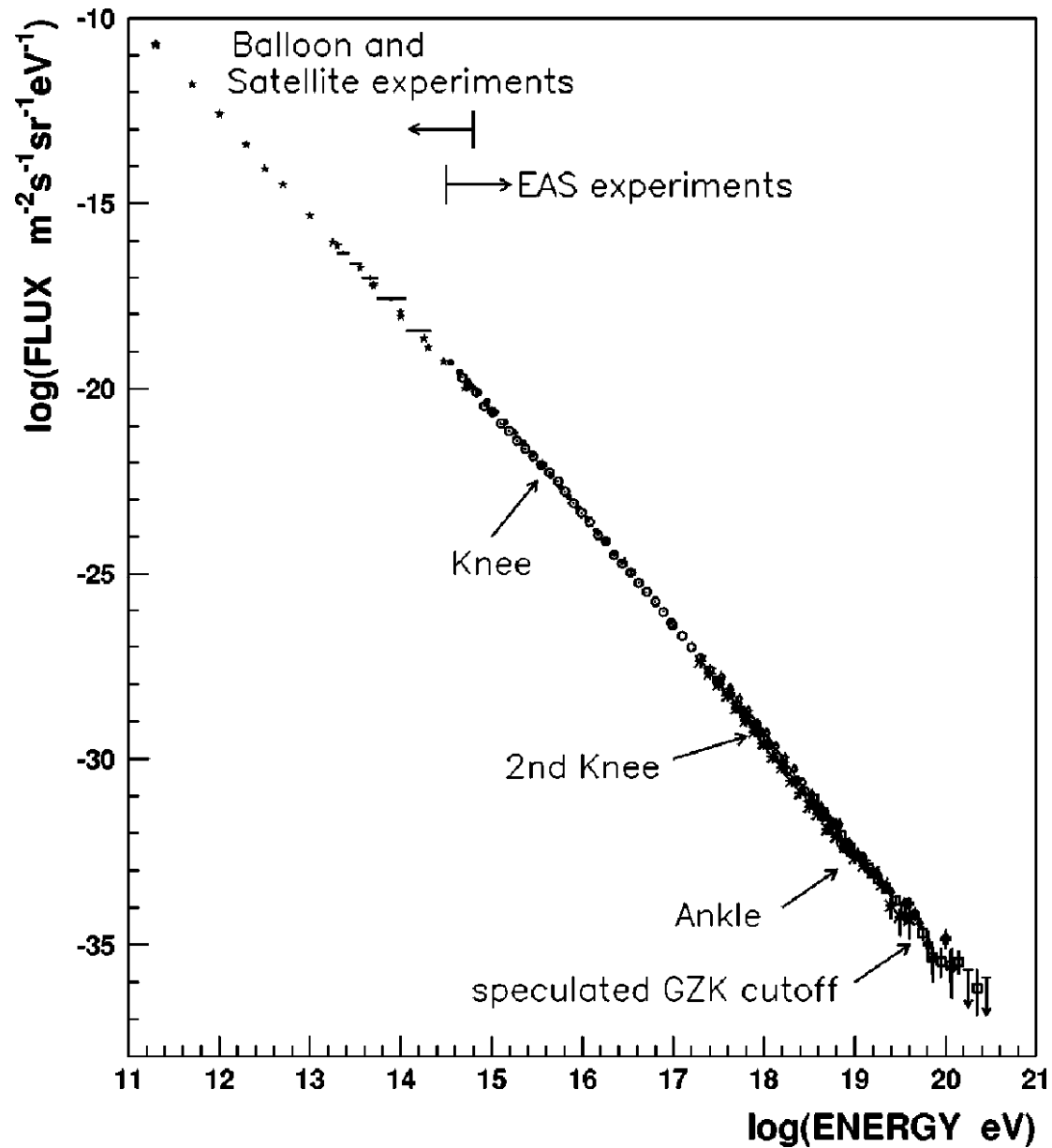
Spallation reactions:
With cosmic rays

LiBeB synthesis

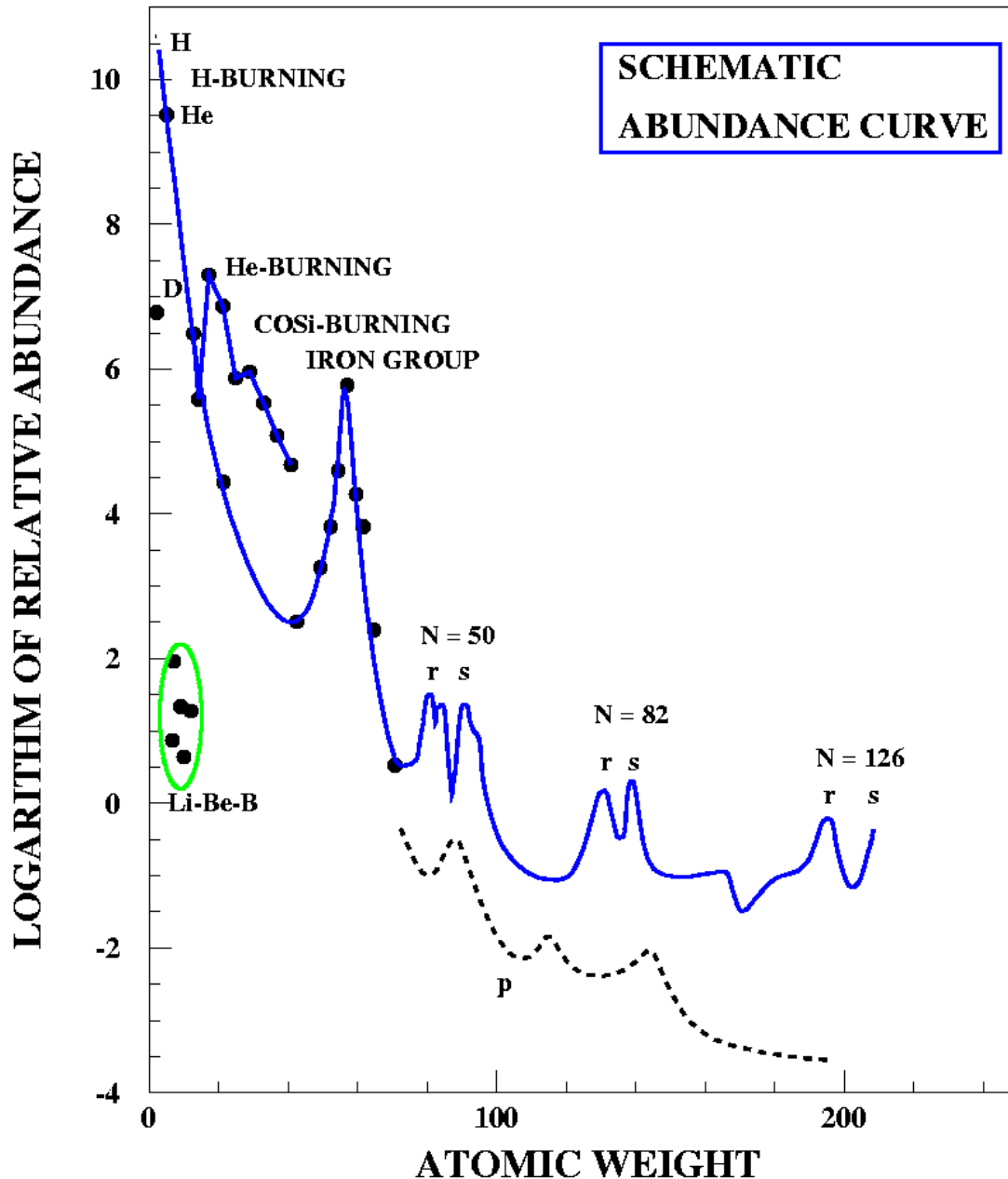
- LiBeB synthesis: spallation reactions in
 - galactic disk : RC (p, α ; GeV/nucleon) & MIS (C, N, O) (Meneguzzi et al, 1971)
 - halo : SN (α , C, N, O ; MeV/nucleon) & MIS (p, α)
- Observations :
 - visible-UV in galactic halo (CFH, HST, Keck, VLT)
 - gamma (COMPTEL, INTEGRAL)
 - RC (ACE, AMS)
- Neutrino spallation during SN explosions
 - ${}^7\text{Li}$ et ${}^{10}\text{B}$ production
 - cross sections are needed



Cosmic rays



Abundance interpretation



May be explained via
nucleosynthesis
processes:

- primordial
- stellar
- explosive
- spallative

Summary of Origins

Element

Site

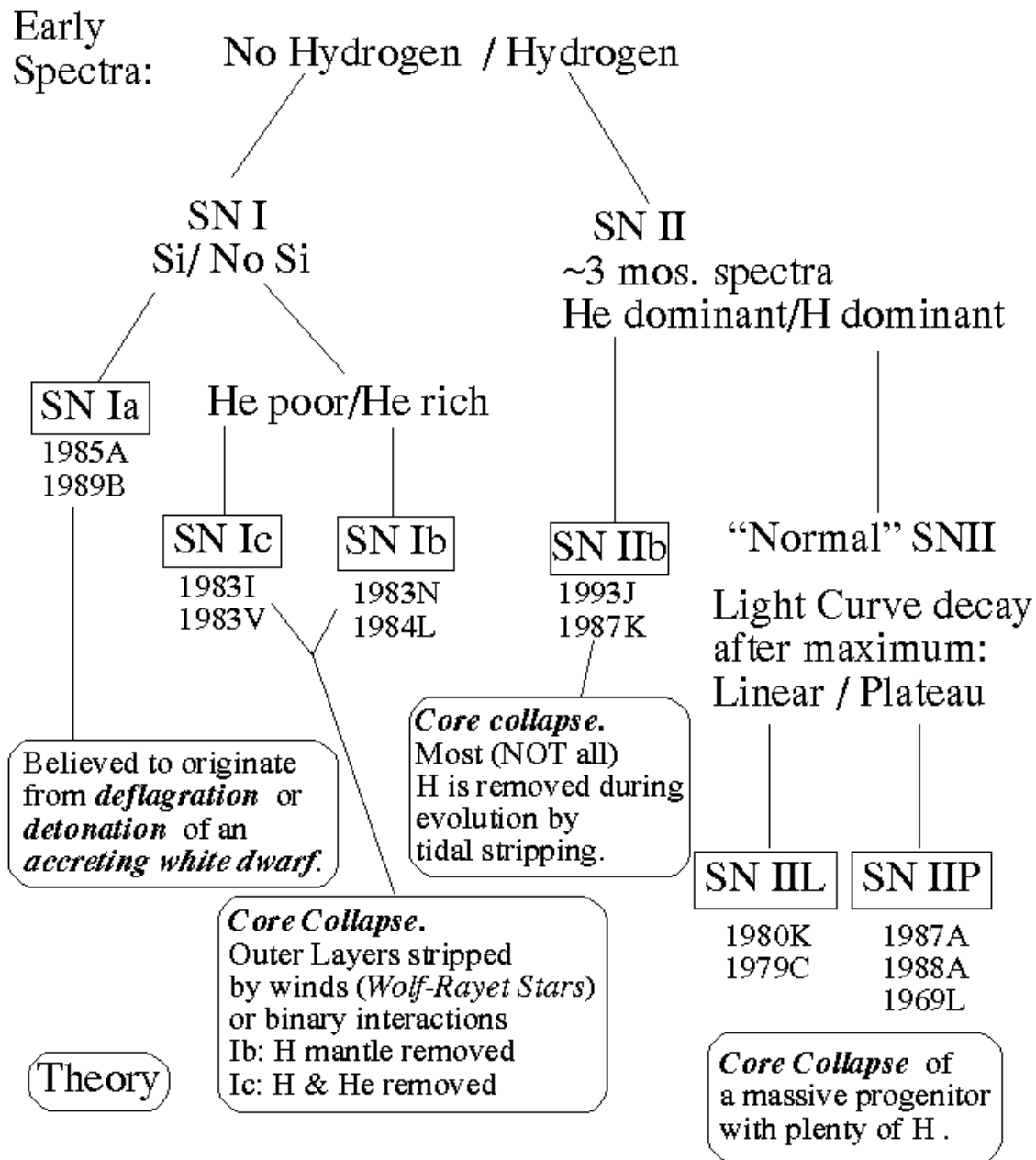
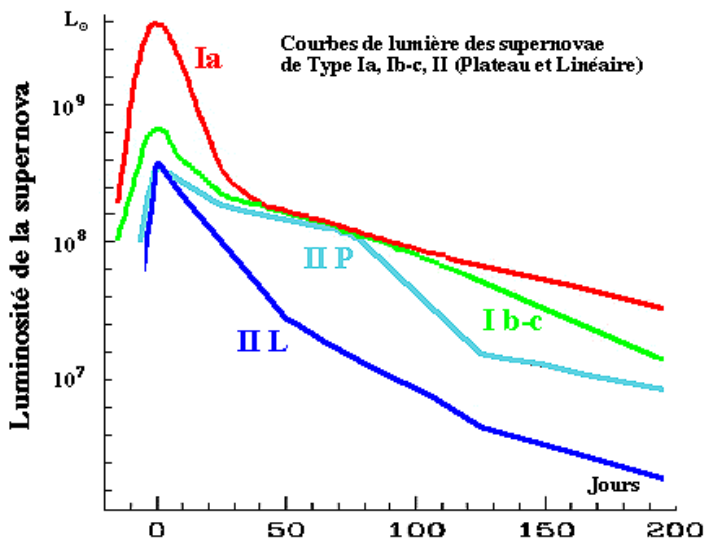
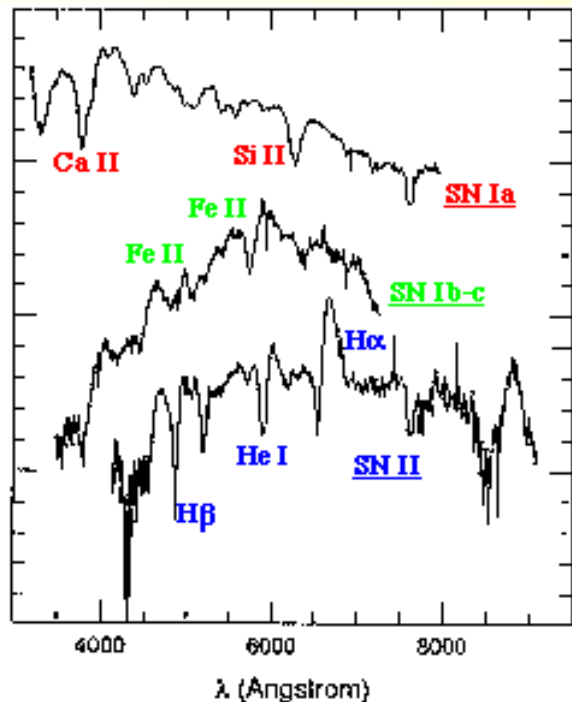
Element

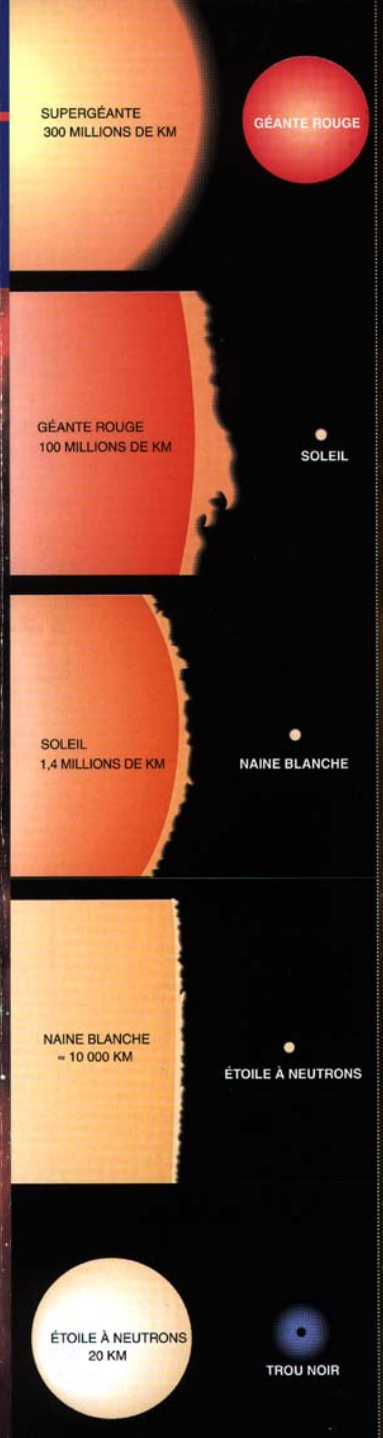
Site

H Big Bang
He Big Bang + stars
Li Big Bang, L* + nu process
Be Cosmic rays
B Nu-process
C Helium burning, L*+M*
N CNO cycle, L*+ VMS
O Helium burning
F Nu-process
Ne Carbon burning
Na Carbon burning
Mg Carbon burning
Al Neon burning
Si Oxygen burning
P Neon Burning
S Oxygen burning
Cl Oxygen burning + s-proc

Ar Oxygen burning
K Oxygen burning + s-process
Ca Oxygen burning
Sc s-process
Ti Expl Si burning
V Expl Si burning
Cr Expl Si burning
Mn Expl Si burning, Ia
Fe Expl Si burning, Ia
Co alpha-rich freeze out
Ni alpha-rich freeze out
Cu alpha-rich freeze out + s-process
Zn Nu-powered wind
p-proc Explosive neon burning, O-burning
s-proc Helium burning, L* and M*
r-proc Nu wind, jets?

Supernovae typology





Relative sizes of stars

Nuclear gamma-ray astronomy *(<10 MeV)*

Interest

- New spectral energy range
- Gamma-ray penetrability
- Constraints on nucleosynthesis (isotopes)
- Cosmic-ray study (acceleration processes, gamma-ray sources)

Production mechanisms

- Annihilation [e^+e^-]

(Galaxy, novae)

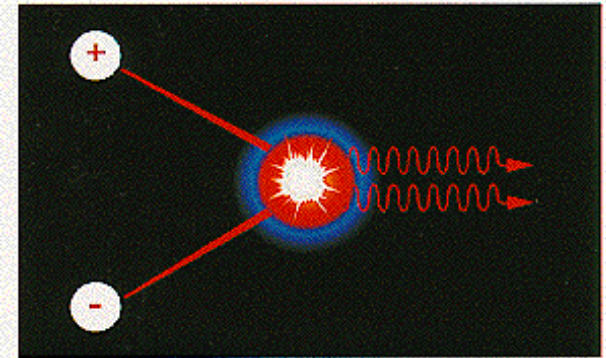
- Nuclear excitations
[(p,p'), (α,α'), ...]

(Solar flares)

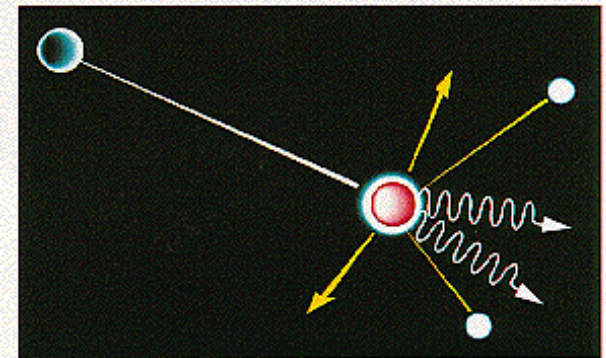
- Radioactivities

[^{26}Al , ^{44}Ti , ^{56}Co , ...]

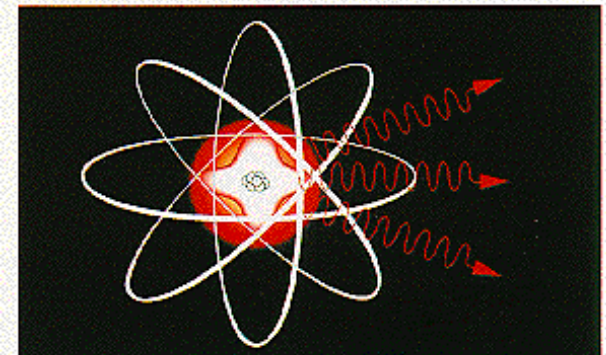
(Supernovae, novae, ...)



MATTER-ANTIMATTER ANNIHILATION

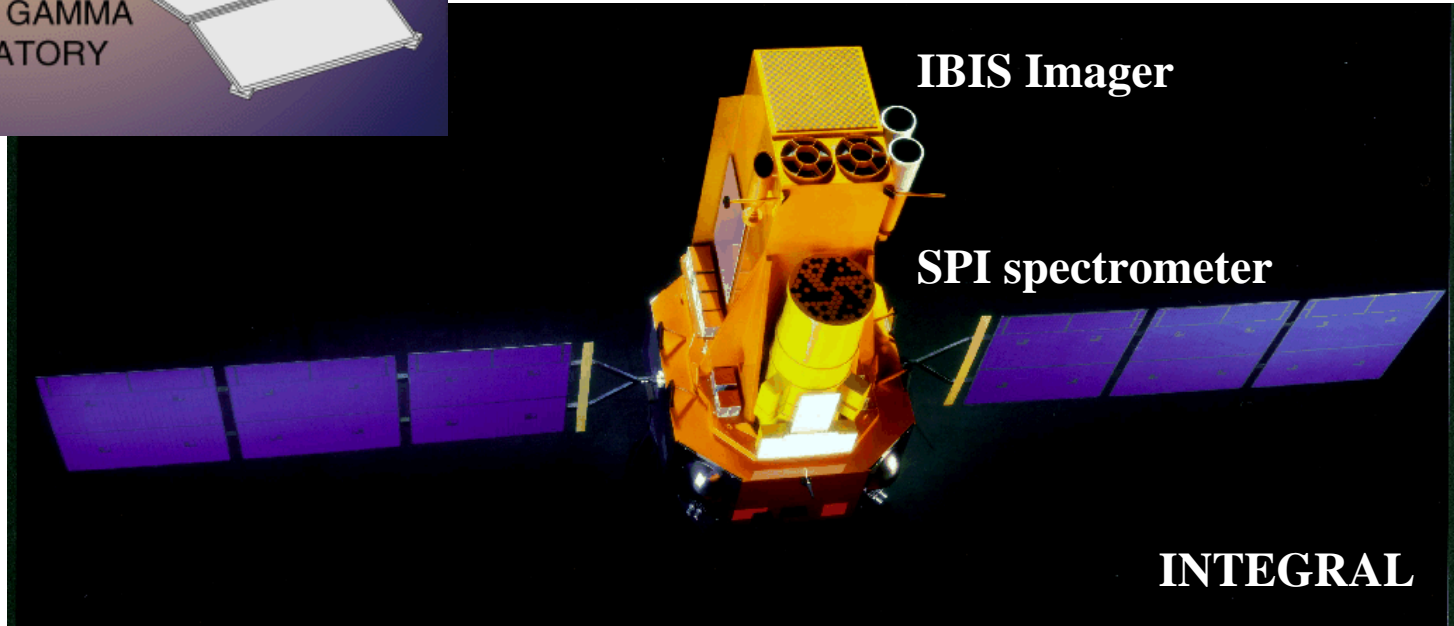
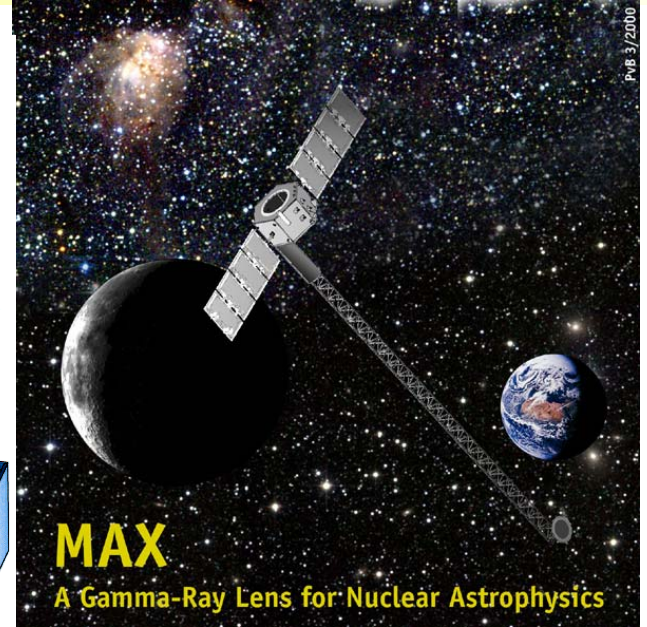
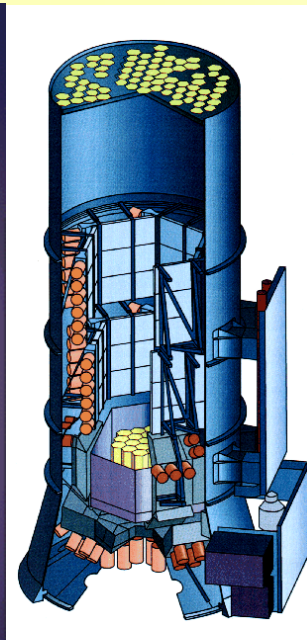
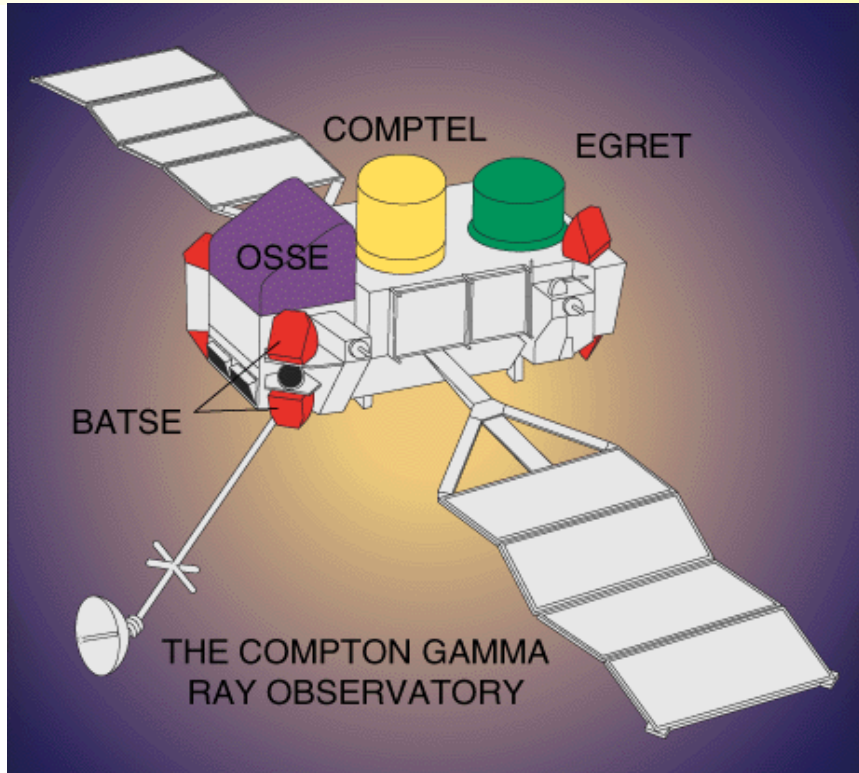


PARTICLE-PARTICLE COLLISIONS

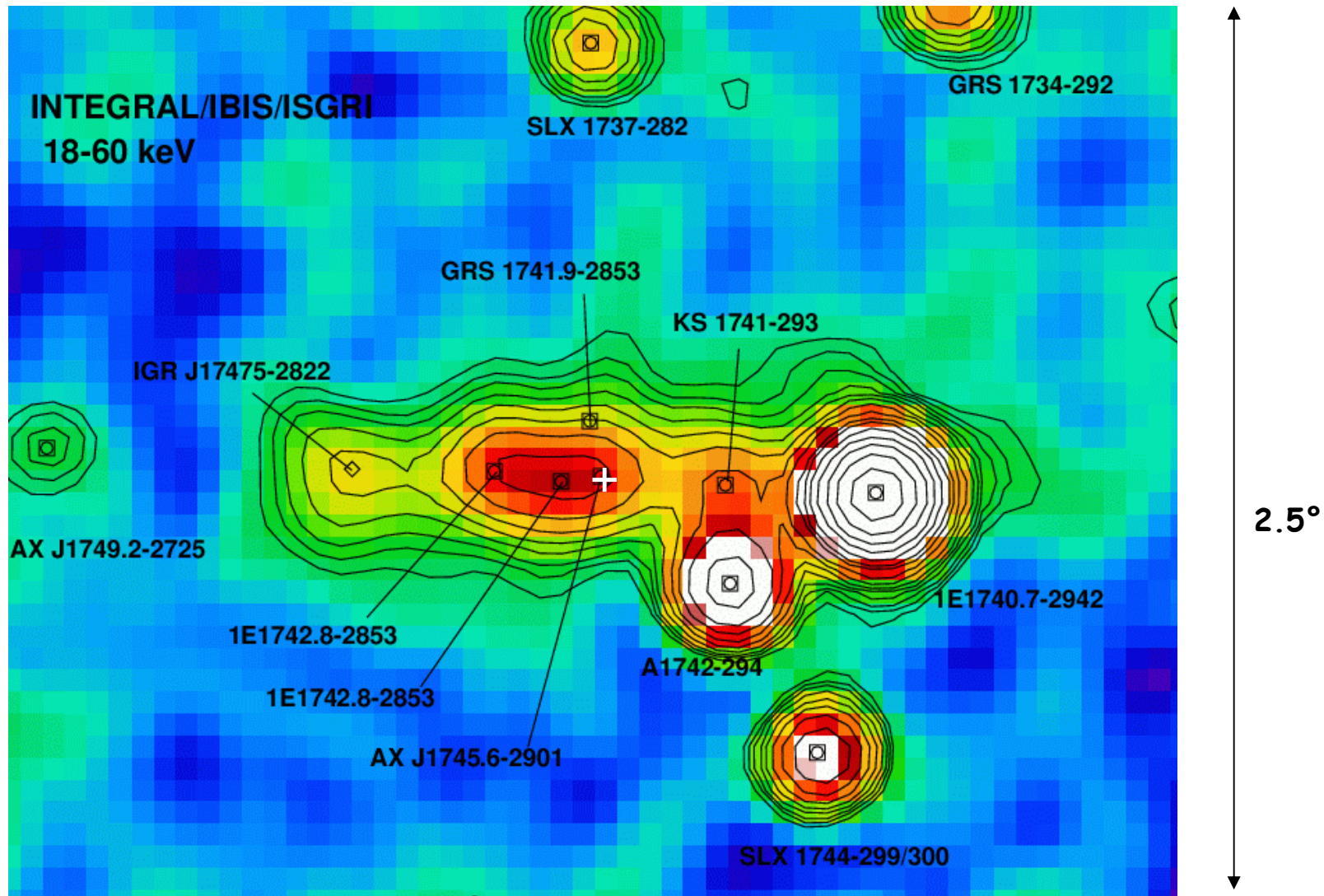


RADIOACTIVE NUCLEI

Some gamma telescopes



INTEGRAL/IBIS imager: galactic center



M. Revnivtsev et al. (IKI Moscow, MPA Garching) et al.

Mostly observable gamma-rays

Isotope	E_γ (MeV)	$T_{1/2}$	Origine	Observé
^7Be	0,478	53,3 j	N, AGB	--
e^+e^-	0,511	--	β^+	Galaxie
$^{18}\text{F}(\beta^+)+e^-$	$\leq 0,511$	110 mn	N	--
^{22}Na	1,275	2,6 an	N	?
^{26}Al	1,809	$7,4 \cdot 10^5$ an	SN,WR,AGB, N	Galaxie
^{44}Ti	1,157	60 an	SN	Cas A, Vela?
$^{56}\text{Ni}/\text{Co}$	0,847; 1,238	77,3 j	SN	SN1987A, SN1991T
$^{57}\text{Ni}/\text{Co}$	0,122;.....	271,8 j	SN	SN1987A
^{60}Fe	1,173; 1,333;..	$1,5 \cdot 10^6$ an	SN	--

AGB : Asymptotic Giant Branch stars

WR : Wolf-Rayet stars

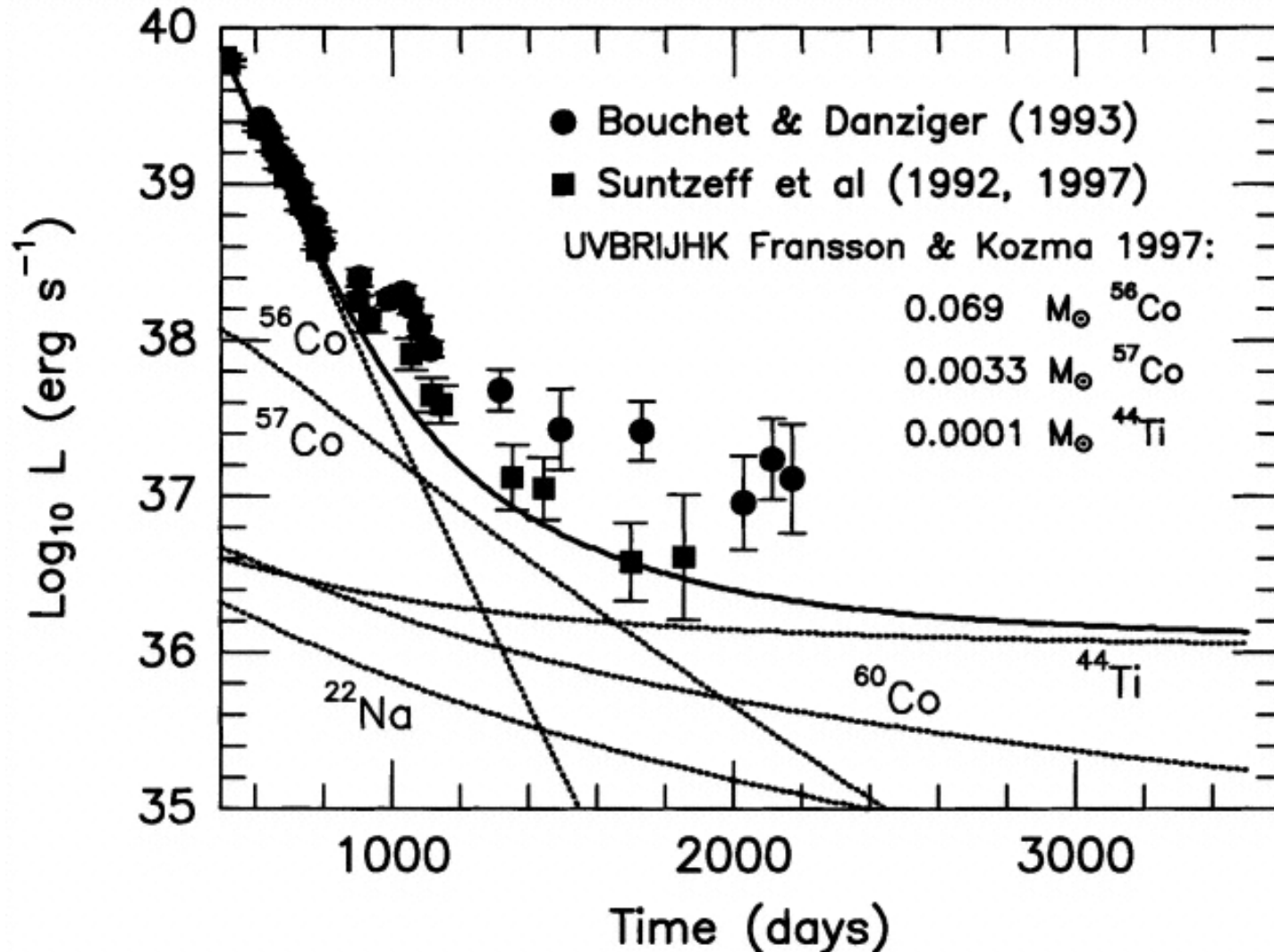
N : Novae

SN : Supernovae

SN1987A light curves

(Diehl & Timmes 1998)

Indirect observation of Co radioactivity:



Some references

General :

Clayton "Principles of Stellar Evolution and Nucléosynthesis
University of Chicago Press

Rolfs &
Rodney "Cauldrons in the Cosmos"
University of Chicago Press

Lang "Astrophysical Formulae" (2 Volumes)
Springer

Ecole Joliot 1990 et 2000 (IN2P3)
Curie

Burbridge "Synthesis of Elements in Stars" (historique)
et al. (B²FH) Reviews of Modern Physics 29 (1957) 547

Wallerstein "Synthesis of Elements in stars: forty years..."
et al. Reviews of Modern Physics 69 (1997) 995

Slezak &
Thévenin "Nucléosynthèse et Abondance dans l'Univers"
CEPADUES-Editions

s process:

Kaeppler "S-process nucleosynthesis - nuclear physics ..."
Rep. Prog. Phys. 52 (1989) 945

Some references

Primordial nucleosynthesis:

Sarkar (Rep. Prog. Phys. 59 (1996) 1493)

Schramm & Turner "Big-Bang enters the precision era"
Reviews of Modern Physics 70 (1998) 303

J. Rich "Fundamentals of Cosmology"
Springer

rp process

Schatz et al. "rp-process nucleosynthesis at extreme ..."
Physics Reports 294 (1998) 167

Novae:

Bode & Evans (Ed.) "Classical Novae"
John Wiley & Sons

Gamma ray astronomy:

Schoenfelder "The Universe in Gamma Rays"
(Editeur) Springer

Supernovae:

Arnett "Supernovae and Nucleosynthesis"
Princeton University Press

Cosmic rays:

Schlickeiser "Cosmic Ray Astrophysics"
Springer